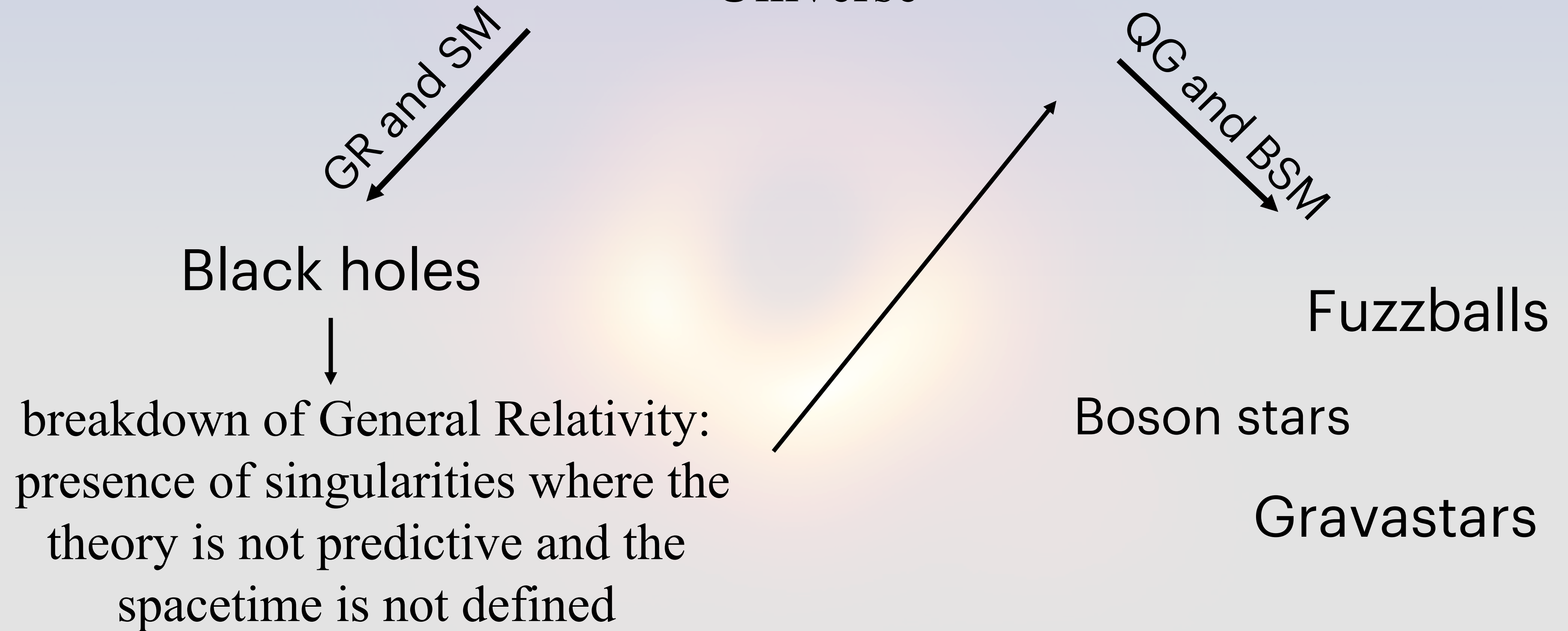


Echoes from backreacting exotic compact objects

Exotic compact objects (ECO)

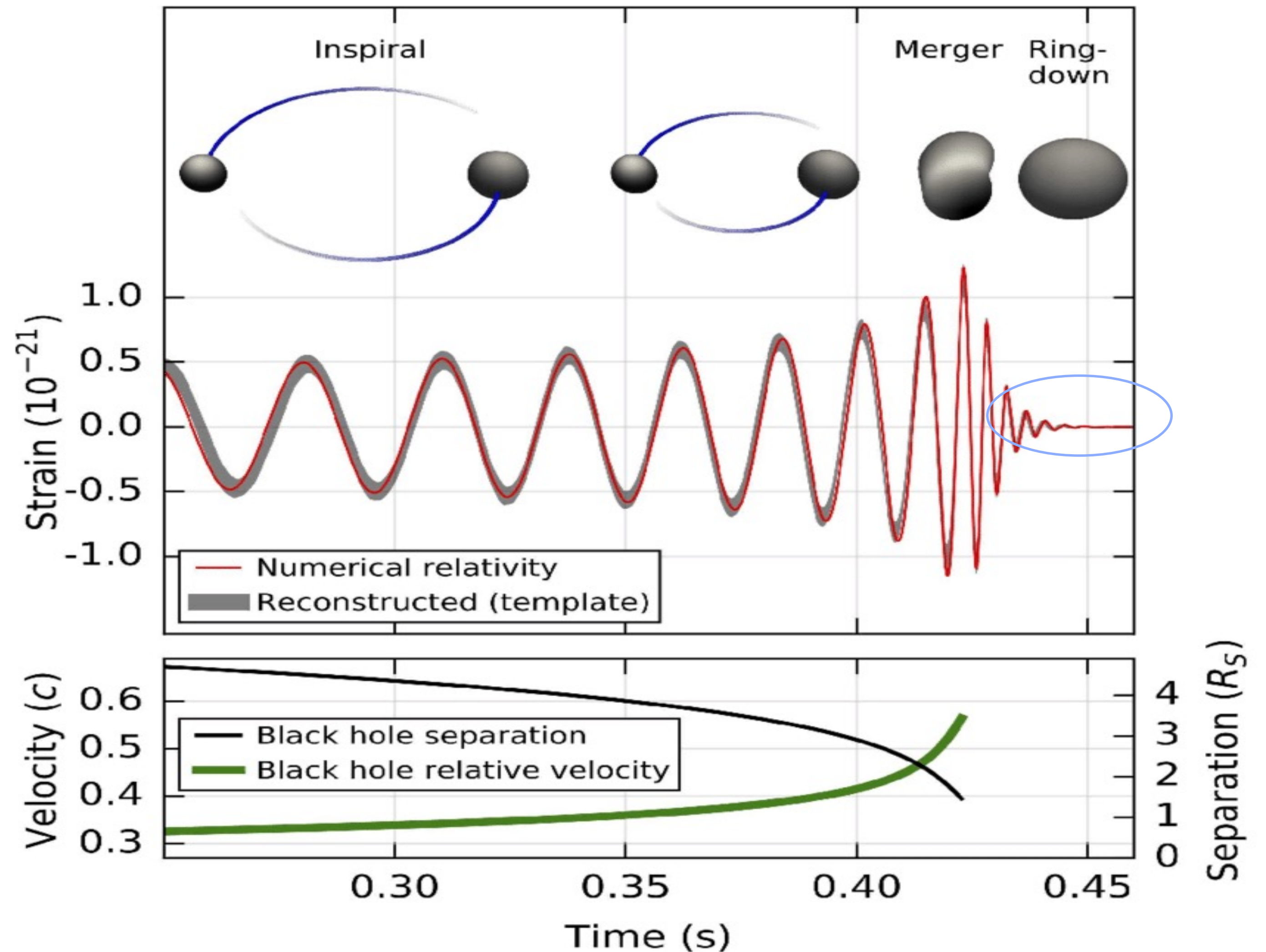
We know that very compact objects exist in our Universe



The source of echoes

The last phase of the GW signal coming from the coalescence of two compact objects is the **ringdown**.

It is caused by the characteristic oscillations of the final object: it can be considered a **gravitational perturbation in the final object space-time**



Echoes

Consider a (scalar) perturbation around a (spherically symmetric) BH or **horizonless**

ECO

$$\Phi = \sum_{lm} Y_{lm}(\theta, \phi) \Psi_{lm}(r)/r$$

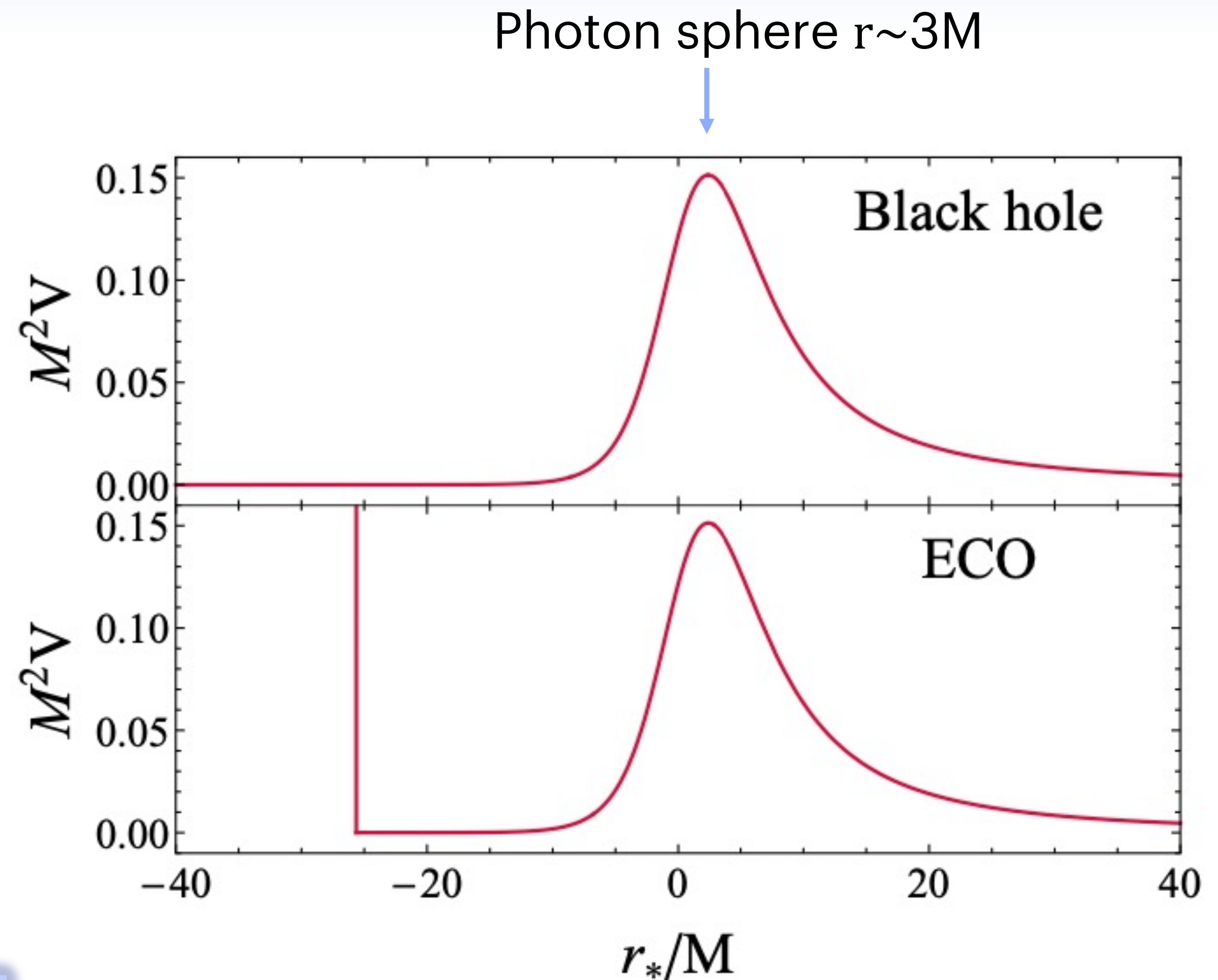
At linear level the field equation is:

$$\left[-\frac{\partial^2}{\partial t^2} + \frac{\partial^2}{\partial r_*^2} - V_l(r) \right] \Psi_{lm}(t, r) = 0$$

But the potential is very different in the two cases

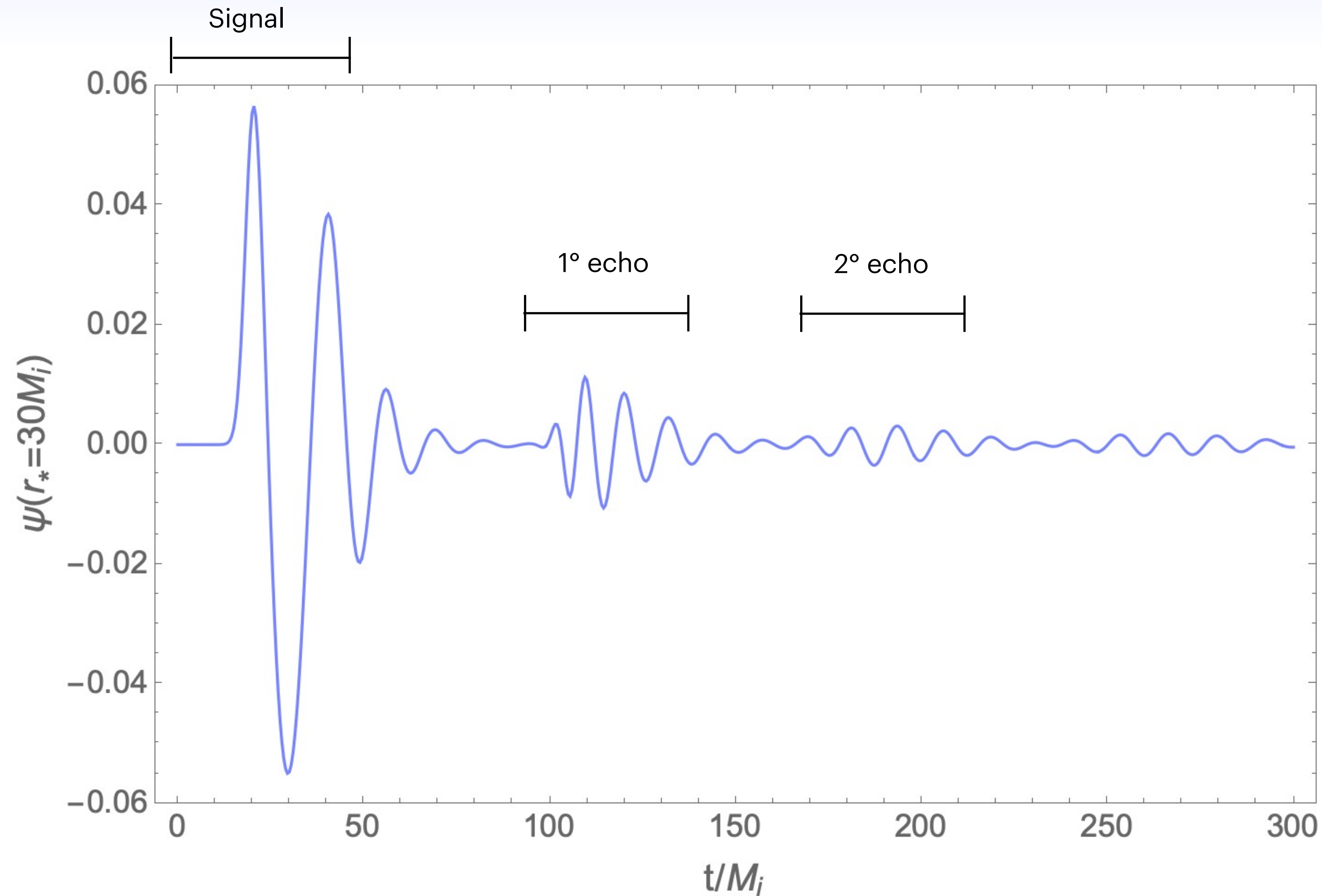
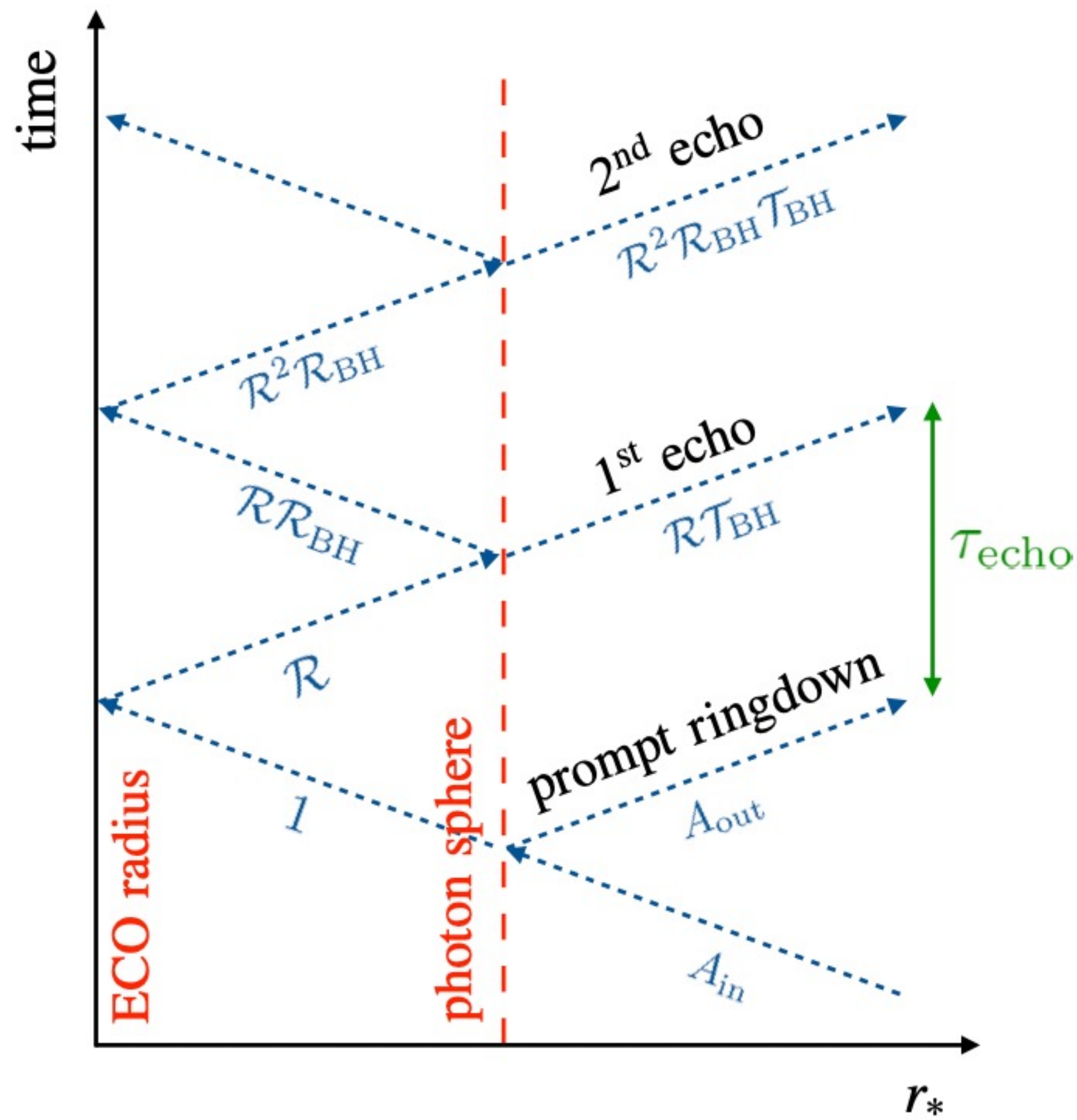


We see echoes of the original signal



From arXiv:2105.06410
E. Maggio, P. Pani, G. Raposo

Echoes



Time delay

Defining the compactness parameter as:

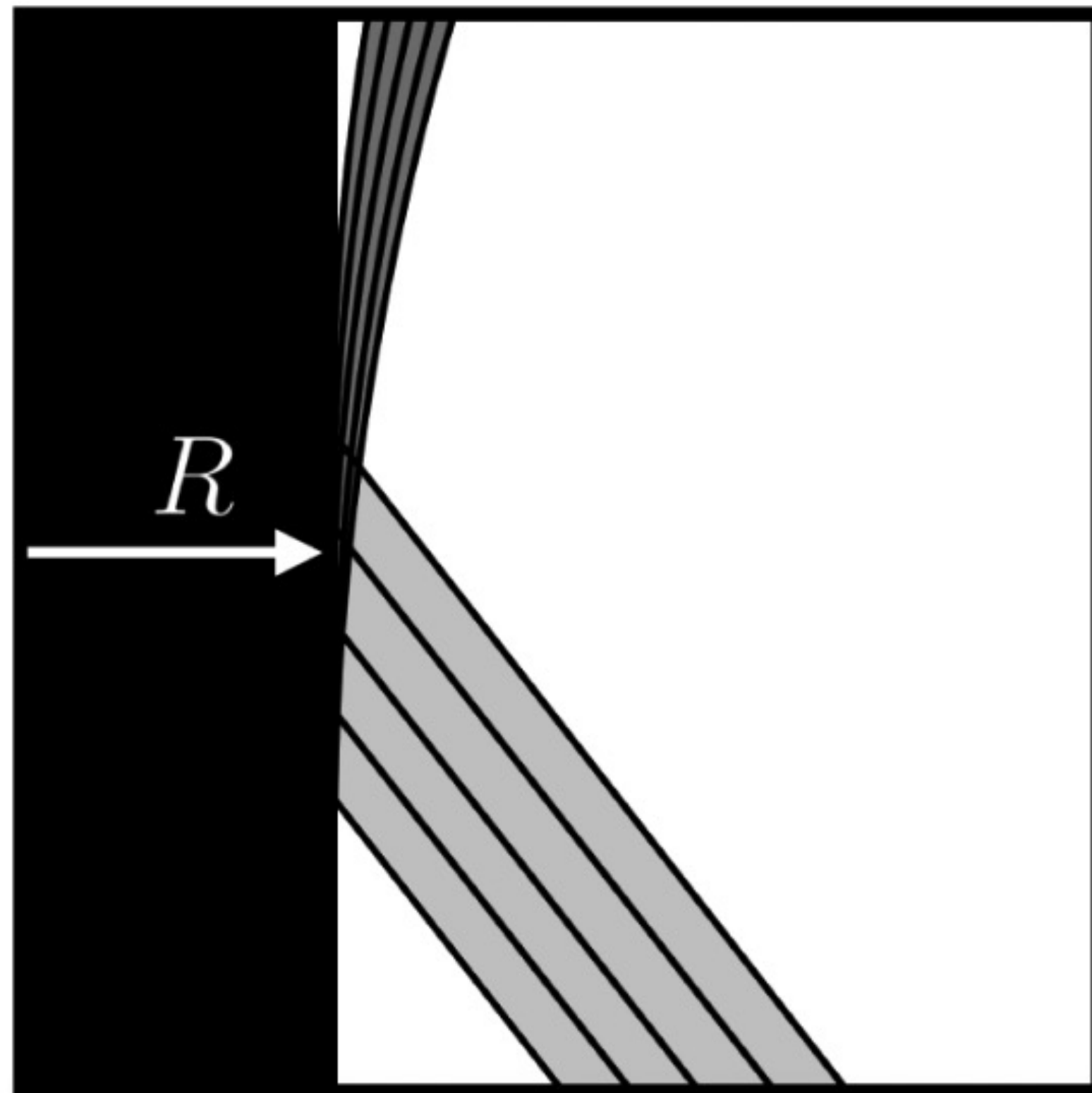
$$\sigma = \frac{r_0}{2M} - 1$$



$$\Delta t_{echo} = 2 \int_{r_0=2M(\sigma+1)}^{r_{peak} \sim 3M} g_{rr} dr = 2 \int_{r_0=2M(\sigma+1)}^{r_{peak} \sim 3M} \frac{dr}{1 - \frac{2M}{r}} \simeq 2M(1 - 2\sigma - 2\ln(2\sigma))$$

The logarithmic dependence on σ would allow to detect even Planckian corrections ($\sigma \sim l_{Planck}/M$) at the horizon scale

Limits of linear approximation



From arXiv:1809.08238

Raúl Carballo-Rubio, Francesco Di Filippo, Stefano Liberati, Matt Visser

Peeling of outgoing geodesic

The accumulation of geodesics around the gravitational radius produces high densities

Instability

If partial absorption is taken into account, the mass of the object can even increase over the hoop limit $r_0 = 2M$



Non linear interactions should be taken into account

Absorption beyond the test field limit

Partial absorption of the first echo $\longrightarrow M_0 \rightarrow M = M_0 + \Delta E_{1st\ echo}$

$\longrightarrow \sigma_{2nd\ echo} \ll \sigma_{1st\ echo} \longrightarrow \Delta t_{2nd\ echo} > \Delta t_{1st\ echo}$

For high compact object
very small ΔM causes big changes
in the compactness!

For example, if $\Delta M = (5 \cdot 10^{-8})M_0$:

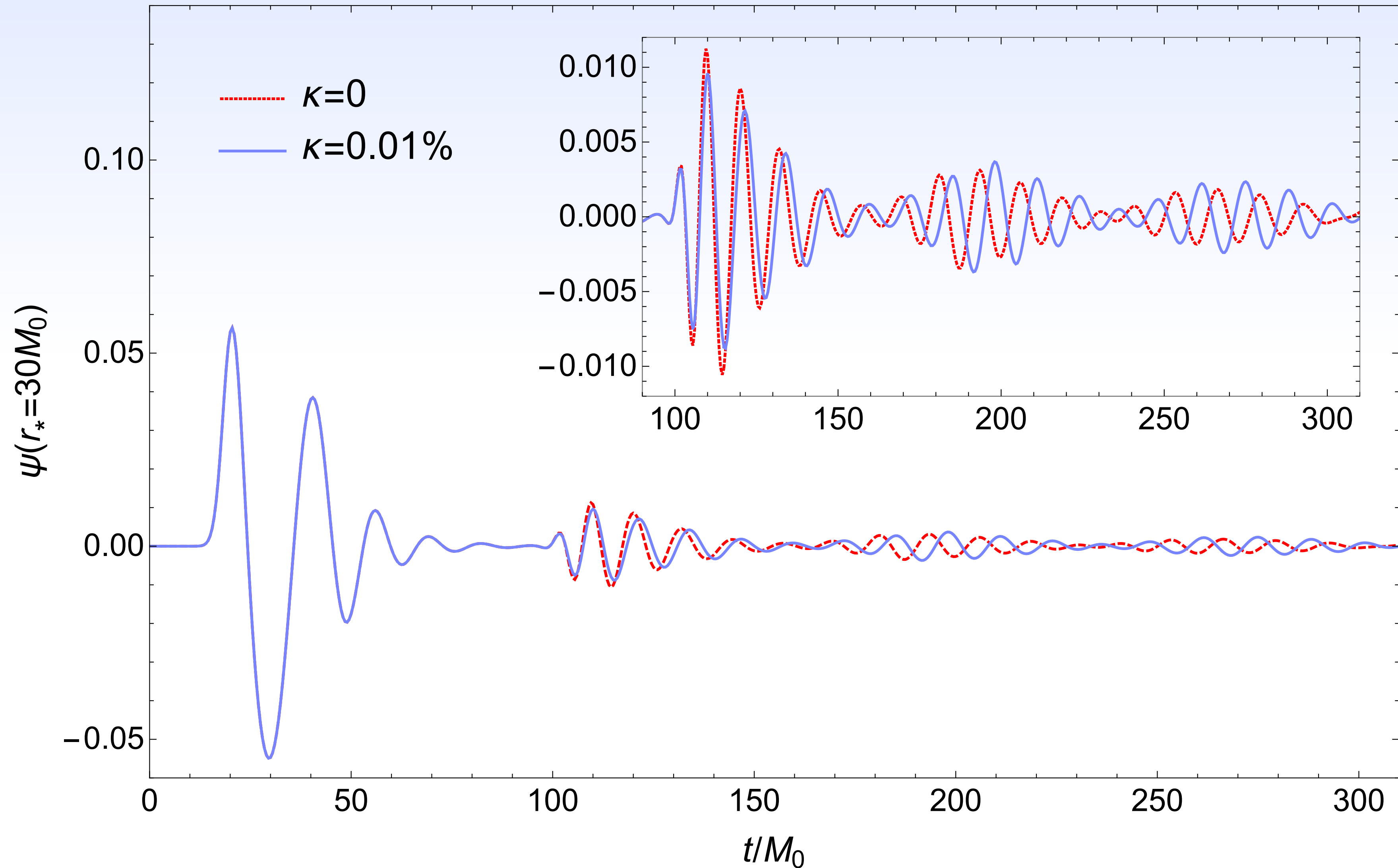
$$\sigma_0 = \frac{r_0}{2M_0} - 1 = \frac{2M_0(1 + 10^{-7})}{2M_0} - 1 = 10^{-7}$$
$$\sigma_f = \frac{r_0}{2M_f} - 1 = \frac{2M_0(1 + 10^{-7})}{2M_0(1 + 5 \cdot 10^{-8})} - 1 = 5 \cdot 10^{-8}$$

$$\Delta t_0 \sim -4 \ln(2 \sigma_0) = 61.7 M_0$$

$$\Delta t_f \sim -4 \ln(2 \sigma_f) = 64.5 M_0$$

Absorption beyond the test field limit

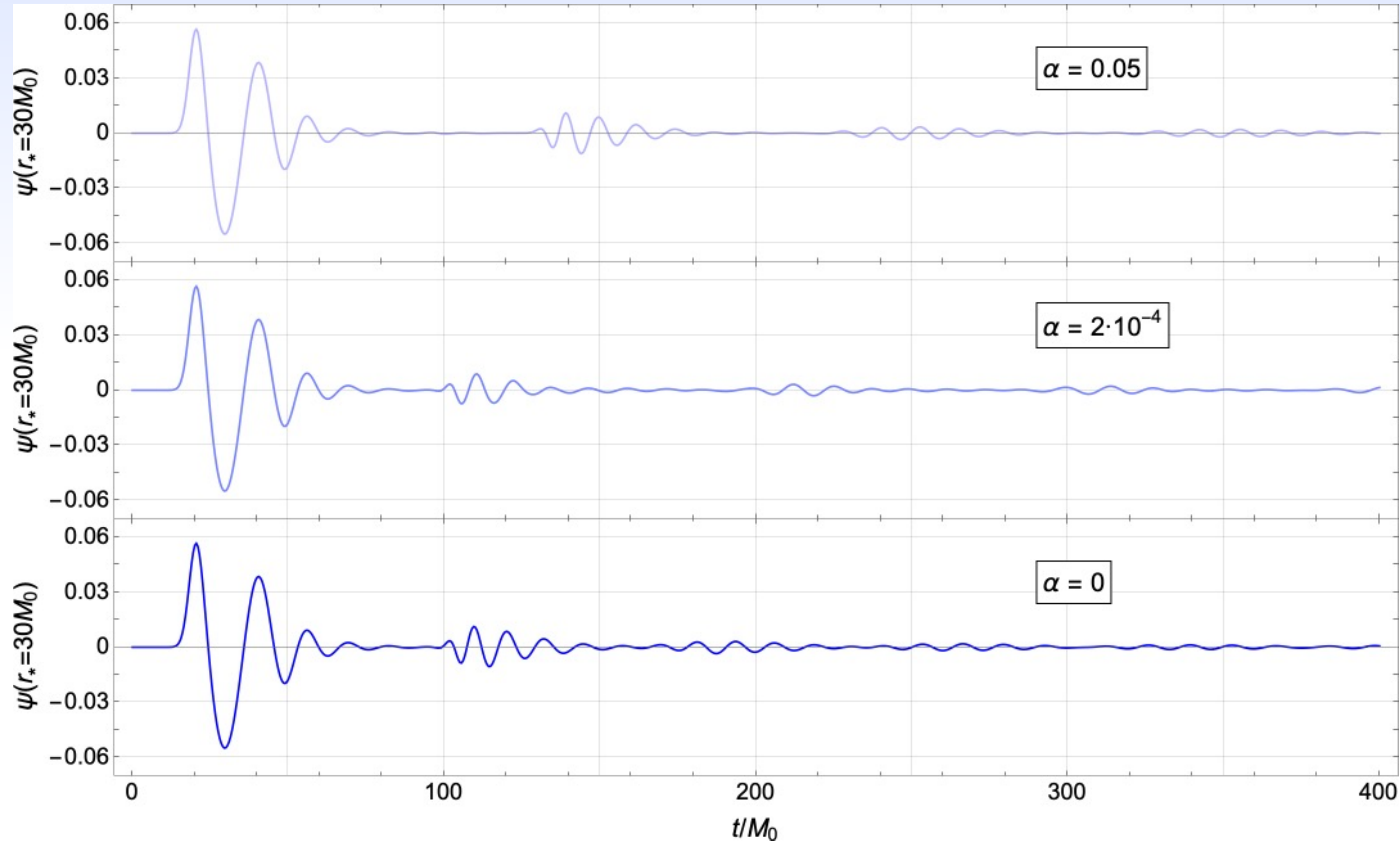
Scalar field in the
Schwarschild space time
with $M \rightarrow M(t)$
and a (partially) reflective
boundary at
 $r = r_0$



arXiv:2205.14170

V.V. , E. Franzin, S. Liberati

Preventing instability I: Asymptotic compactness



Absorption coefficient:

$$\kappa(\sigma) = \alpha \left(1 - \tanh \frac{\beta}{\sigma - \sigma_{Planck}} \right)$$

$$\beta = 10^{-10}$$

Preventing instability II: Expansion

Expansion at constant compactness

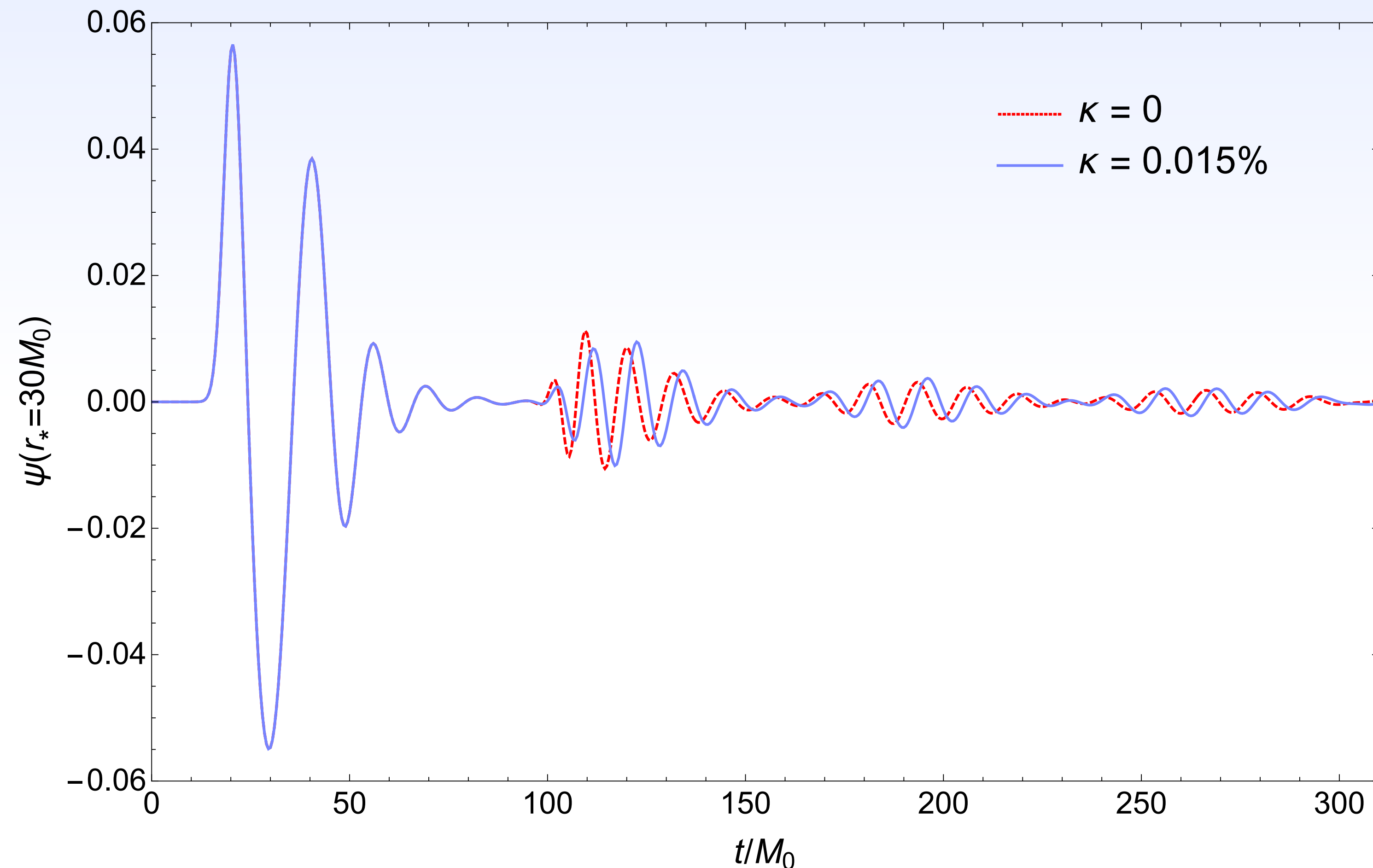
$$\frac{r_0(t)}{2 M(t)} = \sigma_0 + 1$$

$\Delta t_{echo} \sim \text{constant}$

BUT

- There might be a transient phase
- Other, more general model of expansion can be possible

Transient phase $\tau \sim 65M_0 > \Delta t_{echo}$



Conclusions

We saw that, taking into account **possible non linear interaction** between the ECO and the perturbation field, may lead to the partial loss of the main feature of **echoes signal**: the quasi-periodicity

This is important because the searches for echoes in GW ringdown are usually based on this **quasi-periodicity**

It seems that further investigations in the theory and phenomenology of ECOs and more refined searching strategy for such objects are needed.