





GAMMA-RAY DATA COLLECTION, CALIBRATION AND ANALYSIS (FOR LIV STUDIES)

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ABOUT ME

- Gamma-ray astronomy since 2002, LIV studies since 2003
- Associate Professor since 2010 Sorbonne Université Paris
- Member of HESS and CTA
- Research activities:
 - Search for LIV-induced energy-dependent time-delays with high-energy gamma-ray sources
 - Study of source-intrinsic effects in flaring blazars
- Past technical activities
 - On-call expert for HESS camera operation
 - Maintenance of HESS cameras

ABOUTTHIS COURSE

- Preliminary comment:
 - This course will be entirely focused on ground-based high-energy gamma-ray astronomy
- Four parts:
 - From image collection to raw data
 - Calibration
 - Reconstruction
 - Analysis + A brief look at Lorentz Invariance Violation analyses
- Before that: some basics, to get familiar with some of the vocabulary

NOTES/CAVEATS

- This lecture is mostly based on my experience with HESS
- Other IACT Arrays can use different implementation of the same basic methods/ functionalities
- A huge thanks to HESS colleagues who helped me gathering the material for this lecture: M. De Naurois, J.-P. Lenain, F. Toussenel, S. Ohm, a lot of PhD students, and many other colleagues...

- Three major experiments/Cherenkov telescope arrays are currently operating
- The goal is to study the VHE sky
- For this, we **take pictures** of the Cherenkov emission from cascades developing in the atmosphere
 - Atmosphere is a part of the detector
- The signal is **faint** and very **short**
- So, to build a Cherenkov telescope, you need
 - A **large mirror** to collect Cherenkov light
 - A **fast** and **sensitive** camera in the focal plane
- Additional constraints exist!



- We are interested in **photons** produced by an astrophysical source
- Dominant background: hadrons



From Aharonian & Akerlof, 1997

γ-ray spectra 1 and 2 correspond to spectral indices of 2 and 3 respectively.

- **Detecting** a particle means
 - **Identifying** it → gamma-ray or hadron?
 - **Characterizing** it → energy, time, direction, impact location
- Notion of event
 - A gamma-ray (or a hadron) interacts in the atmosphere, the Cherenkov light is reflected by the mirror and collected by a camera of HESS
 - Necessary condition: the camera is **triggered**
- **Raw data** is recorded only when the camera triggers
 - For each channel (pixel), amplified integrated pulses converted into ADC counts
 + other parameters
- Calibration is necessary to convert ADC counts to a charge, and assess
 optical efficiency → multiple ingredients required!

- Reconstruction is used for
 - Discrimination/Identification → **discriminating variables**
 - Selection cuts allows to obtain a sample of gamma-like events
 - Different cuts will give different amounts of purity for that sample
 - Characterization
 - Several techniques can be used: Hillas, Model,...
- Analysis is used on gamma-like events for
 - Precise **background estimation** and subtraction
 - Evaluation of **statistical significance**
 - **Signal extraction**: energy spectrum, light curve, morphology
 - Sometimes long observations are needed
 - Reconstructed event list can also be provided at this step

• If the source is suitable, only then can we attempt a LIV analysis

- We are interested in **photons** produced by an astrophysical source
- Dominant background: hadrons
- The signal:
 - Primary cosmic rays hit the atmosphere
 - They produce a shower/cascade of secondary particles
 - Charged secondary particles produce Cherenkov light
- A shower develops within O(10 ns)
 - The camera and electronics need to be **fast** enough!



300 GeV photon (left) and hadron (right)

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300 GeV photons (top) and hadrons (bottom)

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- Let's be more specific!
- We are interested in Cherenkov photons from the cascades which peaks around 350 nm
- But photo-detectors are also sensitive to « normal » light in their range of operation: Night Sky Background (NSB)
- Photodetectors have to be carefully chosen





FROM IMAGE COLLECTION TO RAW DATA

- Characteristics of the pre-upgrade CT5 camera (2012-2019)
 - 2,40 × 2,27 × 1,84 m³
 - Field of view 3.2°
 - 128 drawers with 16 PMTs each
 - 2048 pixels
 - PMT diameter 29 mm
 - PMT spectral range 270-650 nm
 - Peak QE of 30% at 420 nm
 - Gain 2×10^5 at $\sim 800 \vee$
 - 50 ADC counts for I p.e.
 - Integration window width 16 ns
 - Dead time 15 µs
 - Maximum operational trigger rate 5000 Hz



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DATA ACQUISITION

- Simplified view of front-end electronics for HESS CT5 camera
 - View of I drawer



PMT



• The PMT signal looks like this:



NB: here we see only pulses which triggered the scope, so with an amplitude exceeding a fixed threshold

- Typical width for HESS PMTs is 3.5 ns
- The pulse integral is proportional to the amount of light received
- PMT bases have a protection system so that HV is switched-off when illumination is too high

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AMPLIFIERS

- Linear amplifiers
 - **High gain** (x16) amplification for low amplitude signals (< 200 p.e.)
 - I p.e. \rightarrow 7 mV amplitude
 - **Low gain** (x0.4) amplification for high amplitude signals (< 5000 p.e.)
 - Very high gain (x40) for the **trigger** channel
- Note that
 - Signal is inverted and baseline is shifted towards negative voltages to use the ADC full input range
 - Impedance matching is crucial all along the signal path



DATA ACQUISITION

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TRIGGER

- The array uses **camera triggers** and a **central trigger**
- Different levels of trigger are used
 - L0: pixel level → The charge in a pixel exceeds a given threshold (4 p.e.)
 - LI: camera level → A number (≥ 3) of close-by pixels fulfill the L0 criteria
 - L2: array level
 - L2 trigger makes use of the **central trigger** output
 - ► Only CT5 has triggered → Monoscopic event
 - ► Another telescope has triggered within 80 ns → Stereoscopic event
- Different modes are used for different purposes
 - Internal trigger for normal data taking
 - **External** trigger for specific calibration runs (Single photo-electron runs)
 - Software-generated trigger for testing/warm-up



DATA ACQUISITION

- Simplified view of front-end electronics for HESS CT5 camera
 - View of I drawer



- Analog memory is used to store the signal for some time, until the trigger decision is taken
 - LI trigger decision takes ~200 ns
 - Analog, because ADC conversion takes time (~250 µs for a full event)
- Signal is sampled at I GHz, I ns for each sample
- Circular buffer of 256 cells



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PMT signal sampling



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PMT signal sampling

The read-out is performed (16 ns window)



DATA ACQUISITION

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DATA ACQUISITION

- Once the trigger is received, analog memories are stopped, and the signal is digitized to be sent to the DAQ cluster
- FIFO buffers are used all along the signal path to make the tasks of analog memories and of the rest of the DAQ independent
 - It helps reducing the dead time (15 µs) and make it much smaller than the read-out time (250 µs)
- So, **raw data** consists in integrated charge as ADC counts
 - For each pixel
 - For each gain (HG and LG)
- Slow control data is intertwined with raw data, at lower rate (1 block for ~1000 events)
 - High voltages, high voltage base currents, temperatures

A SIDE NOTE ON CAMERA DESIGN

- The design presented here is the one used for CT5 « old » « French camera ». It is inherited from CAT, and HESS-I camera designs.
- Other designs exist
- Most recent ones show improved performance
 - Better sensitivity, lower energy threshold, lower dead time
- Some examples:
 - MAGIC: analog signals are sent to the ground through optical fibers
 - HESS-II new camera design (NamCam): signals from the PMTs are digitized directly and sent to an FPGA for trigger, data processing and event transmission
 - NectarCam: evolution from HESS-II old camera readout with improved performance
 - Dead time < 500 ns
PART II

CALIBRATION

CALIBRATION

- Calibration is the step where ADC counts are converted into a charge in (photoelectrons) p.e.
- Cherenkov light intensity is also assessed in order to get the **optical efficiency**
- It is a multi-step process requiring specific types of data
 - What is the electronic signal baseline? → Pedestal
 - How to convert ADC counts to a charge at HG? \rightarrow Single photo-electron (SPE)
 - What about LG channel? \rightarrow Hi/Lo ratio
 - How to deal with pixel efficiency non-uniformities? → **Flat-field**
 - How to convert charge to Cherenkov light intensity?
 - What is the optical transmission efficiency? → Use of muons + atmospheric monitoring
- From then on, we are going to deal with **charge distributions**
 - Charge is obtained integrating PMT pulses

PEDESTAL

- **Pedestal** = baseline of the electronics
- It needs to be calibrated with high accuracy since Cherenkov signals are deduced using this baseline as a reference
- Can be acquired from:
 - Dedicated runs: camera lid closed, random trigger



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 - Data taking runs: identifying and filtering out pixels hit by Cherenkov light
 - More NSB → wider pedestal



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GAIN CALIBRATION

- Gain is used to convert ADC counts into a charge expressed in photo-electrons (p.e.)
- Overall gain is from
 - PMT, and
 - Amplifier
 - Obviously, it is different for HG and LG channels
- Dedicated runs are necessary:
 - With camera protected from ambient light (dark or using a filter)
 - Low intensity flashes so that I p.e. is produced at the photocathode on average
 - External trigger required
- (Simulated) Signal and charge distribution in the HG channel:



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 I p.e. (SPE, single photo-electron) is produced at
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- Luminosity on the camera:



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 I p.e. (SPE, single photo-electron) is produced at
 the photocathode on average
 - External trigger required
- Charge distribution and fit:
 - Nominal gain is 50 ADC/p.e.
- NB: SPE runs are also used to set PMT HV...



GAIN AND AGING

- PMTs are subject to aging \rightarrow gain decrease with time \rightarrow need for regular monitoring
- HV adjustments are made from time to time to re-adjust the gains to 2×10^5



HI/LO GAIN CROSS-CALIBRATION

- SPE cannot be resolved in the LG channel
- Real data is used to cross-calibrate
 High and Low gain channels in
 the intensity range where they overlap



FLAT-FIELD

- **Flat-fielding** is used to correct for non-uniformities of collection efficiency (cones + PMT) and PMT Quantum Efficiency
 - Camera is illuminated by a very uniform light source located at the center of the dish
 - Cross-calibration coefficients are obtained for each pixel



WHAT WE HAVE SO FAR

• Up to now, we have a conversion from ADC to a number of photo-electrons, averaged over the spatial efficiency of the camera:

$$S_{HG} = \frac{ADC_{HG} - P_{HG}}{\gamma_{HG}} \times FF$$

$$S_{LG} = \frac{ADC_{LG} - P_{LG}}{\gamma_{HG}} \times (HG/LG) \times FF$$

with:

- **S**: signal in p.e.
- **ADC**: signal in ADCc, **γ**: gain
- **P**: pedestal
- **FF**: flat-filed coefficient
- HG/LG: Hi/Lo gain ratio
- Charge in p.e. will be the main ingredient of the next steps of data analysis
- But we still want to evaluate optical efficiency

- To get the number of Cherenkov photons hitting the telescope, we need to know:
 - Reflectivity/transmissivity of optical elements as a function of wavelength
 - Telescope collection efficiency
 - PMT QExCE as a function of wavelength.
- All this is difficult to measure...
- Atmospheric muons are used to get the overall optical efficiency

- Muons offer a natural calibration beam:
 - Radiation spectrum is very similar to Cherenkov emission from gamma-ray showers
 - Muons reaching the ground are ultra-relativisitic, making the light yield independent of the muon energy → it depends only on the track length
 - Muon signal is very short in time
- HESS small cameras can see the last ~400 m of the muon track (~700 m for CT5)
- The overall optical efficiency is given by:

Eff. opt. =
$$\text{Re} \times (1-\text{Sh}) \times \text{Co} \times \text{QE}$$

where:

- **Re** is the mirror reflectivity integrated over wavelength (~80%)
- Sh is the shadow of camera and masts on the mirror (~10%)
- **Co** is the collection efficiency (~70%)
- **QE** is the PMT quantum efficiency (~30%)



• Some simulated and real examples:





• The number of photons reaching the camera per unit wavelength, azimuth angle, track length is given by:



- Atmospheric absorption can be obtained from models or deduced from measurements (LIDAR)
 - Transparency is reduced by dust, aerosols...
- The technique must be proven with Monte Carlo simulations!
- Long-term monitoring of optical efficiency is important
 - PMT cleaning, Winston cone replacement, mirror refurbishment...



WHERE ARE WE NOW?

- For each pixel, we end up with a charge in p.e., and a way to assess Cherenkov light intensity
- This will be the main ingredient for the next step: reconstruction and analysis.
- The calibration step is also very important as a first assessment of instrument performance
- Hardware problems detected during calibration must be taken into account in the reconstruction and analysis steps
 - An inoperative pixel can be removed from the analysis only for a part of a run...
- Calibration software requires great care: a bug can introduce biaises in the reconstruction and analysis steps
 - Each time the calibration software is updated, all data sets need to be reprocessed...
 - Fortunately, effects of these bugs are usually very small!





To be continued...

PART III

RECONSTRUCTION AND ANALYSIS

WHAT DID WE LEARN SO FAR?

- Yesterday, we've seen
 - How images are collected
 - How we can convert from ADC to a charge and to the Cherenkov light intensity
 - Multiple step process
- In practice, reconstruction and analysis use charge in p.e. as main input
- Optical efficiency must be taken into account in Monte Carlo simulations, which provide a charge in p.e. for each pixel
- The next steps are
 - **Reconstruction**: shower images are processed to discriminate gamma-like and hadronlike events and to reconstruct shower parameters

- Analysis:

- if the signal is strong enough (significant), spectral analysis, light curve extraction...
- If the signal is not significant, limits on the flux can still be derived
- Additional criteria apply for LIV analysis
- Before that: run selection, data quality checks, etc. \rightarrow not covered...

RECONSTRUCTION

Goals of reconstruction:

- Reconstruct the parameters of the primary particles: direction, impact location and energy
- Provide discriminating variables to reduce as much as possible the amount of background events. Rejection factor is typically O(10⁴)
- Two methods will be covered in this session
 - Hillas reconstruction
 - Model reconstruction
- They are quite simple to implement, and widely used
- They are also complementary and can be cross-checked or even combined
- Many more exist! Some are based on the use of neural networks, boosted decision trees, etc.
 - More complexity, longer computing times... but better overall performance

- First step: image cleaning
 - Ideally, keeps only the pixels with light from the shower
 - Two thresholds:
 - A pixel with a signal $>s_1$ is kept only if it has at least a neighbor with intensity $>s_2$
 - Different values can be used depending on the amount of NSB (4, 7), (7, 10)...



- Images look like this, for a photon, and for a hadron
- The gamma shower is quite elliptical in shape...



- In 1985, Hillas proposed to model the images by an ellipse
- Hillas parameters:
 - Length **L** and width **w** of the ellipse
 - **Size** (total image amplitude in p.e.)
 - Nominal distance d
 - Azimuthal angle of the image main axis ϕ
 - Orientation angle α
- This parametrization can be expressed from the intensity q_i (in p.e.) and position (x_i, y_i) of each pixel
- From these, several quantities are computed:

$$\langle x \rangle = \frac{\sum_{i} x_{i} q_{i}}{\sum_{i} q_{i}} \quad \langle xy \rangle = \frac{\sum_{i} x_{i} y_{i} q_{i}}{\sum_{i} q_{i}} \quad \sigma_{y^{2}} = \langle y^{2} \rangle - \langle y \rangle^{2}$$

$$\langle y \rangle = \frac{\sum_{i} y_{i} q_{i}}{\sum_{i} q_{i}} \quad \langle x^{2} \rangle = \frac{\sum_{i} x_{i}^{2} q_{i}}{\sum_{i} q_{i}} \quad \langle y^{2} \rangle = \frac{\sum_{i} y_{i}^{2} q_{i}}{\sum_{i} q_{i}} \quad \sigma_{xy} = \langle xy \rangle - \langle x \rangle \langle y \rangle$$



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$$\chi = \sigma_{x^{2}} - \sigma_{y^{2}}$$

$$z = \sqrt{\chi^{2} + 4\sigma_{xy}}$$

$$b = \sqrt{\frac{(1 + \chi/z) \langle x \rangle^{2} + (1 - \chi/z) \langle y \rangle^{2} - 2\sigma_{xy} \langle x \rangle \langle y \rangle}{2}}$$



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- From these, several quantities are computed
- And then we get Hillas parameters from image quantities:



$$d = \sqrt{\langle x \rangle^2 + \langle y \rangle^2}$$

$$L = \sigma_{x^2} + \sigma_{y^2} + z$$

$$w = \sigma_{x^2} + \sigma_{y^2} - z$$

$$\alpha = \arcsin\left(\frac{b}{d}\right)$$

- Single telescope reconstruction
 - Degeneracy for source direction...



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 - Additional parameters are added to remove the degeneracy
 - Recent hardware developments can help: addition of timing in image analysis



- Single telescope reconstruction
 - Degeneracy for source direction...
 - Additional parameters are added to remove the degeneracy
 - Recent hardware developments can help: addition of timing in image analysis
 - Possibility to resolve cascade development as a function of time
 - Example with a hadron:





- Stereoscopy is used by all major experiments
 - Intersection of image axes in the camera frame gives the position of the source



- Stereoscopy is used by all major experiments
 - Intersection of image axes in the camera frame gives the position of the source
 - Intersection of planes containing the telescope position and the shower axis gives the _ position of shower impact on the ground (also called shower core)
- NB: pointing accuracy, camera position in the focal plane have to be carefully controlled!





situation!

- Gamma-hadron separation
 - Hillas parameters can be used to discriminate between photons and hadrons
 - Several methods exist. An example is the so-called Scaled Cuts method
 - The idea is to compare image width and length with the expectation values and variance obtained from simulations as a function of image charge **q** and impact distance **ρ**
 - Scaled Width (SW) and Scaled Length (SL)

$$SW = \frac{w(q,\rho) - \langle w(q,\rho) \rangle}{\sigma_w(q,\rho)} \qquad SL = \frac{l(q,\rho) - \langle l(q,\rho) \rangle}{\sigma_l(q,\rho)}$$

- For stereoscopic observations with **ntels** telescopes:
 - Mean Scaled Width (SW) and Mean Scaled Length (SL)

$$MSW = \frac{\sum_{tels} SW}{\sqrt{ntels}} \qquad MSL = \frac{\sum_{tels} SL}{\sqrt{ntels}}$$

 Example of MSW and MSL distributions for the Crab Nebula Simulated gamma-rays in red, Gamma-like events from the Crab in green, Background events in blue



- Main interest of Scaled Cuts:
 - They are scaled by the variance of the parameters, e.g.
 - Small energy \rightarrow big variance \rightarrow small parameters
 - Big energy \rightarrow small variance \rightarrow big parameters
 - This will turn into the fact selection efficiency will not depend on energy
 - → No bias for energy spectrum determination

- Energy reconstruction
 - Use of simulated events to build look-up tables
 - Primary energy is deduced from impact parameter and image size (intensity)
- Tables
 - For each shower:
 - Reconstruct the impact point (the core)
 - Compute the intensity in each telescope
 - Events are binned by reconstructed distance to the telescope and image intensity
 - For each of these bins, average and variance of true energy is computed
 - This is repeated for 20 bins for zenith angle, 7 bins for optical efficiency, 6 bins for off-axis angle. Intermediate values are obtained through interpolation
- This interpolation allows to get the reconstructed energy for each telescope from the size, using the impact point reconstructed with the full array
- Final energy estimation is obtained from a weighted average. Energy resolution ~20%, with biases at low and high energies



MODEL RECONSTRUCTION

- Model reconstruction was developed by De Naurois & Rolland (2009) from early works done in CAT in the 90s (Le Bohec, 1998)
- NB: Cherenkov light distribution of a shower is given by the longitudinal, lateral and angular distributions of charged particles
- The method:
 - Produce these distributions from Monte Carlo simulations
 - Parameterize the distributions \rightarrow shower model
 - The first interaction depth parameter is taken into account
 - Produce the images in the cameras, adding the noise due to NSB in each pixel
 - Reconstruction and discrimination are achieved through the comparison between raw images and modeled raw images
 - Done with a log-likelihood maximization technique
- In HESS, this technique allows a x2 improvement of sensitivity compared to Hillas reconstruction

MODEL RECONSTRUCTION

- Examples of modeled shower images
- I TeV vertical showers
- X,Y units: degrees in the camera frame
- Z scale: intensity in p.e.
- 0.01° pixel size








- It is important that the model is accurate
- Need to compare it with simulations
- Example:



- Average and shower-to-shower fluctuations from simulations is shown in black (I TeV, zenith)
- Model prediction in red, pixel size 0.16° (H.E.S.S.)

- A likelihood fit is used to reconstruct energy, direction, impact, first interaction depth by comparing real raw images and models
 - This comparison is done at the pixel and p.e. level!
 - The main ingredient is the charge distribution in each pixel
 - Pixels are assumed to be independent
- The output of the fit
 - Best guess for the 6 shower coefficients: direction (2), impact (2), depth of first interaction, energy
 - Uncertainties
 - Final log-likelihood value
- Compatibility between recorded events and pure gamma-ray hypothesis is assessed from Goodness-of-fit:

Pixel log-likelihood
$$G = \frac{\sum_{\text{pixel } i} \left[\ln L(s_i | \mu_i) - \langle \ln L \rangle |_{\mu_i} \right]}{\sqrt{2 \times \text{NDoF}}}$$
Pixel expectation valueNormalization

- Discrimination between gamma-ray and hadron induced showers relies on several facts:
 - Hadron induced showers are irregular → image usually show clusters
 - Hadronic part of the shower produce low intensity Cherenkov light which is spread over the camera → « hadronic rain »
 - Primary charged particle can produce Cherenkov light



- From this, we can build discrimination variables
- Shower Goodness (SG)
 - Goodness-of-fit, using pixels above 0.01 p.e. grouped together with three rows of neighbors
- Background Goodness (BG)
 - Goodness-of-fit, using all the other pixels
- As already done before, we can define another parameter using images from several telescopes:

Mean Scaled Shower Goodness (MSSG)



- Where we are now
 - Two methods for reconstructing event parameters
 - These methods can be used also as discrimination tools
- Before reconstruction, we need to chose a number of initial selections such as
 - Minimum charge (p.e.) in the camera
 - Maximum nominal distance (deg.)
 - Minimum number of telescopes required
 - Maximum Shower Goodness
 - Primary depth range
 - Max θ^2 (θ is the angle between the expected position of the source and the reconstructed position) \rightarrow post-reconstruction criteria
- Each set of cuts will lead to
 - Different performance!
 - Different list of reconstructed events!

• Examples:

Name	Min. Charge	Max. Nom.	# Tels	SG_{\max}	t_0	$ heta_{ m max}^2$
	(p.e.)	Distance (deg.)			(X_0)	(deg^2)
Standard	60	2	2	0.6	[-1, 4]	0.01
Faint Source	120	2	2	0.4	[-1, 4]	0.005
Loose Cuts	40	2	2	0.9	N/A	0.0125

- Min. Charge has a direct influence on energy threshold
- Faint source: for low flux, hard spectrum
 - Good angular resolution, so smaller cut on θ^2
- Loose Cuts: for high flux, soft spectrum



- Performance of reconstruction is assessed through
 - Its ability to reconstruct shower parameters with accuracy
 - Its ability to reject hadron background
- Performance of the detector is commonly quantified by effective area, Etrue vs Erec relation i.e. migration matrix, which also contains values for energy resolution and bias → Instrument Response Functions (IRF)
 - Depend on zenith angle, offset angle, observation conditions, reconstruction strategy...



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- Performance of the detector is commonly quantified by effective area, Etrue vs Erec relation i.e. migration matrix, which also contains values for energy resolution and bias → Instrument Response Functions (IRF)
 - Depend on zenith angle, offset angle, observation conditions, reconstruction strategy...



FROM RECONSTRUCTION TO ANALYSIS

- The reconstruction step gives us a list of gamma-like events (« Signal »), including
 - Real gamma from the source under study
 - Real gamma from background
 - Real electrons
 - Mis-identified hadrons
- We now want to focus as much as possible on source gamma events to extract source properties:
 - Spectrum
 - Light curve
 - Morphology

- So, we need to estimate the background as precisely as possible...
 - Goal of the analysis



ANALYSIS

- The detection of a source usually requires multiple runs
- These runs are usually taken in different observation conditions
 - Zenith angle
 - Offset angle
 - Presence of hardware problems
- The goal of the **analysis** is to combine these runs and precisely estimate the background
- Some background estimation techniques require to know the acceptance of the detector → heavy use of simulations
 - Probability that an event of a given energy and type triggers the system, is reconstructed, and is selected according to the considered event class



• It's not the case for the « Multiple-OFF » or « reflected background » technique

• Say we want to investigate a particular position in the sky



- Problem: there may be other sources in the same FOV! \rightarrow exclusion regions
- The ON Region is centered on the expected position of the source



 Acceptance is radially symmetric, so we consider a circle around the observation position



- Background is taken in OFF Regions which are centered on that ring
- There can be several OFF Regions → « Multiple OFF »



• A real example for two different runs



- When a high number of runs is necessary, the whole process needs to be automated...
- A database with exclusion regions is updated regularly

SIGNIFICANCE

- We define:
 - **Non** the number of events in the region of interest (ROI)
 - Noff the number of events in control regions (OFF regions) with only background events
 - α the ratio of the expected number of background events in the ROI over the actual number of background events accumulated in OFF regions
 - α can be understood as the ratio of exposures between the ON region and the OFF regions
 - The excess is given by Non α Noff
 - **Significance** is obtained from Li & Ma (1983): $S = \sqrt{-2 \ln \lambda}$

with

$$\lambda = \left[\frac{\alpha}{1+\alpha} \left(\frac{N_{on} + N_{off}}{N_{on}}\right)\right]^{N_{on}} \times \left[\frac{1}{1+\alpha} \left(\frac{N_{on} + N_{off}}{N_{off}}\right)\right]^{N_{off}}$$

- λ is the **likelihood-ratio** between two hypotheses:
 - Non events result from Signal + α Background
 - Non events result from background only (null hypothesis)

SIGNIFICANCE

- For pure Poisson fluctuations for both ON and OFF events, S follows a normal distribution of average 0 and width 1
- Excess events show up as a tail on top of this distribution
- Another control plot: θ^2 distribution



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SENSITIVITY

- For a given analysis chain and a given set of cuts, it is possible to estimate the sensitivity of the array
- Sensitivity: minimum flux which can be detected with Nσ significance during T hours of observations. Other criteria may be added: 10 detected gamma-rays per energy bin, S/B > 1/20...
- Example for point-like sources:



WHAT'S NEXT?

- From now on, we will focus on a simple case: point source, with negligible background
 - This can happen for e.g. a flaring AGN or a GRB!
- We will also consider that the source is detected with a significant signal
 - Usually a detection is claimed when S > 5 σ
- We are usually interested in
 - The energy spectrum of the source
 - Its light curve
 - Ingredients for LIV studies!
- We'll now briefly see how to get the spectrum

SPECTRUM

- A commonly used technique for spectrum extraction is the so-called « forward folding » technique
- It requires a good knowledge of Instruments Response Functions (IRFs): migration matrix (Etrue vs Erec), energy resolution *R*, effective area *A*
 - IRFs are different for each reconstruction method and depend on selection cuts!
 - IRFs vary with observation conditions!
- Steps:
 - Assume a spectral shape Φ (power law, broken power law, etc.)
 - Compute in each bin of energy the **expected** number of gamma events:

$$n_{\gamma} = \int_{E_{rec,1}}^{E_{rec,2}} \mathrm{d}E_{rec} \int_{0}^{\infty} \mathrm{d}E_{true} \mathcal{R}(E_{rec}, E_{true} | \delta, \Psi, \epsilon) \times \mathcal{A}(E_{true} | \delta, \Psi, \epsilon) \times \Phi(E_{true} | \vec{\alpha})$$

with δ the zenith angle, Ψ the off-axis angle, ϵ the optical efficiency

- A log-likelihood is built from the number of ON and OFF events, taking into account a proper live time normalization $\pmb{\beta}$

$$P(N_{ON}, N_{OFF} | n_{\gamma}, n_h) = \frac{n_{\gamma} + \beta n_h^{N_{ON}}}{N_{ON}!} e^{-(n_{\gamma} + \beta n_h)} \times \frac{n_h^{N_{OFF}}}{N_{OFF}!} e^{-n_h}$$

where n_{γ} and n_h are the expected numbers of gamma and background events in the energy bin.

- The log-likelihood is then maximized

SPECTRUM

- The results of the maximization are:
 - The best fit parameters α
 - The expected number of events in each bin
 - The covariance matrix → used to derive flux uncertainty in each bin
 - The final value for log-likelihood →
 ≈ goodness-of-fit
- Flux points are obtained from the model using best fit parameters and residuals obtained in each energy bin
- Main advantage of the forward-folding technique: N_{ON} and N_{OFF} used directly, no need for flux point calculations



PART IV

LIV ANALYSIS

WHAT WE WANT TO TEST

Some theoretical approaches predict a modified dispersion relation for photons in vacuum



E_P: Planck scale ~10¹⁹ GeV

- So, photons of different energies could travel at different speeds in vacuum
- Two photons emitted at the same time at redshift z would arrive with a delay

$$\Delta t_n \simeq s_{\pm} \, \frac{n+1}{2} \, \frac{E_h^n - E_l^n}{E_{QG}^n} \, \int_0^z \frac{(1+z')^n}{H(z')} dz'$$

where the integral deals with Universe expansion

- This effect is very small but cumulate on large distances
 - Example values for $E_{QG} = E_P$
- When no significant lag is measured, a limit is set on E_{QG}

	z = 0.1	z = 1
ΔE = 100 GeV	0.3 s 3x10 ⁻¹⁸ s	4 s 5x10 ⁻¹⁷ s
∆E = 1 TeV	3 s 3x10 ⁻¹⁶ s	40 s 5x10 ⁻¹⁵ s
$\Delta t_n \text{ for } n = 1$ n = 2		

HOW

- So, we want to measure a time lag between photons of different energies
- Need for sources that are
 - Distant
 - Variable or transient
 - Energetic
- Candidates:
 - Gamma-Ray Bursts (GRBs)
 - Flaring Active Galactic Nuclei (AGNs)
 - Pulsars (PSRs)
- These sources all have advantages and drawbacks
- The « Variable or transient » criteria is essential
 - The more variability we have, the more sensitive we'll be
 - A galactic millisec. PSR can give more stringent limits than a distant flaring AGN...
- Several techniques have been developed to measure energy-dependent time delays

Source	Observed by	Distance	Variability time scale
AGN flare	IACT	z _{max} ~ 0.6	O(1 min)
GRB	Satellites	z _{max} ~ 9	O(0.1 s)
PSR	Satellites & IACTs	d _{max} ~ 50 kpc	O(1 ms)
Source	Random ?	Intrinsic effects	Emission mechanisms
Source AGN flare	Random ? Yes	Intrinsic effects Hints	Emission mechanisms poorly understood
Source AGN flare GRB	Random ? Yes Yes	Intrinsic effects Hints Yes	Emission mechanisms poorly understood poorly understood

BASICTECHNIQUES

- The main ingredients of LIV analyses are:
 - light curves in different energy bands, or
 - a list of photons (gamma-like) with arrival times, reconstructed true energies, etc. together with IRFs
 - As seen previously, each type of analysis, cuts, etc. will lead to different IRFs and different lists of photons
- A simple technique:
 - Fit the peaks in the light curves in different energy bands
 - Compare the fitted positions of the peaks to get Δt and the corresponding uncertainty
 - From the spectral index and energy ranges, get the value of ΔE
- Other methods were used based on signal processing techniques: CCF, wavelets...





A MORE ADVANCED TECHNIQUE

- Likelihood technique proposed by Martinez & Errando (2008)
- The pdf is computed from the time and energy distributions and takes IRFs into account (effective area, energy resolution)

$$P(E_i, t_i | \tau_n) = N \int_0^\infty D(E_i, E_S) \Lambda(E_S) F(t_i - \tau_n E_S^n) dE_S$$

• The likelihood is then obtained from the events with

$$\mathcal{L} = \prod_{i} P(E_i, t_i | \tau_n)$$

- Maximizing the likelihood gives the best estimate for parameter au
- Main advantages:
 - Good sensitivity
 - Easy to combine multiple sources

A MORE ADVANCED TECHNIQUE

- An example with PKS2155-304 2006 flare:
 - Parameterization of the light curve at low energies
 - Parameterization of the spectrum
 - Likelihood maximization and obtained best estimates





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IMPORTANT CONSIDERATIONS

- Key points of analyses are
 - Performance assessment
 - Statistical and systematic error assessment
- Both of these are usually done using dedicated toy-MC simulations
 - Simulate hundreds of light curves with same properties as real data
 - Reconstruct the lag for all of them and study its distribution → statistical error estimate
- Systematic uncertainties have also to be careful assessed.



	Change in estimated т _I (s/TeV)	Change in estimated τ _q (s/TeV²)	
Selection cuts	< 5		
Background contribution	< 1		
Acceptance factors	< 1		
Energy resolution	< 1		
Energy calibration	< 2		
Spectral index	<	1	
Calibration systematics	< 5	< 1	
F _s (t) parameterization	≈ 7	≈3	
Total	< 10.3	< 6.6	

FINAL REMARKS

FINAL REMARKS

- Hardware, reconstruction, analysis is a huge topic to cover
- With only 4 hours available, a lot had to be ignored
 - Physics of the showers (we'll cover that a bit with the tutorial)
 - New generation instruments (sorry CTA, SiPMs, etc.!)
 - New techniques in reconstruction and analysis (neural networks, etc.)
 - Acceptance determination
 - Other background rejection techniques
 - Details on LIV analyses \rightarrow will be covered during the 3rd training school!

- ...

• Lecture notes will cover that a just a little bit more, including references, and a glossary of most important terms

WE MADE IT!



