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## Study of Universe Transparency in an LIV Framework

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#### 1. Effective Field Theory

- Lagrangian
- Modified dispersion relations
- 2. Kinematical and Dynamical Modifications
- 3. Mean Free Path
  - Optical depth
  - Cross section
- 4. Plots
- 5. Outlook

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# Backstory



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### Motivation

#### [Cao Z., Aharonian F.A. et al., 2021]

#### Article

# Ultrahigh-energy photons up to 1.4 petaelectronvolts from 12 γ-ray Galactic sources

https://doi.org/10.1038/s41586-021-03498-z	A list of authors and affiliations appears at the end of the paper.
Received: 21 October 2020	The extension of the cosmic-ray spectrum beyond 1 petaelectronvolt (PeV; 10 <sup>15</sup> electronvolts) indicates the existence of the so-called PeVatrons–cosmic-ray factories that accelerate particles to PeV energies. We need to locate and identify such
Accepted: 26 March 2021	
Published online: 17 May 2021	
Check for updates	objects to find the origin of Galactic cosmic rays <sup>1</sup> . The principal signature of both electron and proton PeVatrons is ultrahigh-energy (exceeding 100 TeV) yradiation. Evidence of the presence of a proton PeVatron has been found in the Galactic Centre, according to the detection of a hard-spectrum radiation extending to 0.04 PeV (ref. <sup>2</sup> ). Although y-rays with energies algobily higher than 0.1 PeV have been reported from a few objects in the Galactic plane <sup>3</sup> , unbiased identification and in depth exploration of PeVatrons requires detection of y-rays with energies allowe 0.1 PeV. Here we report the detection of more than 530 photons at energy esaysoure 000 teraelectronvolts and up to 1.4 PeV from 12 ultrahigh-energy y-ray sources with a statistical significance greater than seven standard deviations. Despite having several potential counterparts in their proximity, including pulsar wind nebulae, supernova remnants and star-forming regions, the PeVatrons responsible for the ultrahigh-energy y-rays have not yet been firmly localized and identified (except for the Cra) Nebulo. Jeaving one on the origin of these extreme accelerators.

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- Invariance under rotations
- CPT and P invariance
- No dimension 5 operators ([Bolokhov, Pospelov; 2008])
- Only operators quadratic in the fields

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Lagrangian:

$$\mathcal{L} = i\bar{\psi}\gamma^{\mu}D_{\mu}\psi - m\bar{\psi}\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\kappa\bar{\psi}\gamma^{i}D_{i}\psi + \frac{ig}{M^{2}}D_{j}\bar{\psi}\gamma^{i}D_{i}D^{j}\psi + \frac{\xi}{4M^{2}}F_{kj}\partial_{i}^{2}F^{kj}, \qquad (1)$$

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Photon:

$$E^{2} = k^{2} + \frac{\xi k^{4}}{M^{2}}$$
(2)

(3)

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Fermions:

$$\begin{split} \mathrm{E}_\mathrm{e}^2 &= \mathrm{m}^2 + \mathrm{p}^2 \left(1 + \kappa + \frac{\mathrm{g}\mathrm{p}^2}{\mathrm{M}^2}\right)^2 \\ &\approx \mathrm{m}^2 + \mathrm{p}^2 (1 + 2\kappa) + \frac{2\mathrm{g}\mathrm{p}^4}{\mathrm{M}^2} \end{split}$$

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Minimal model:  $g = \kappa = 0$  [Jacobson et al; 2003].

Photon:

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Fermions:

$$E_{e}^{2} = m^{2} + p^{2} \left( 1 + \kappa + \frac{gp^{2}}{M^{2}} \right)^{2}$$
$$\approx m^{2} + p^{2} (1 + 2\kappa) + \frac{2gp^{4}}{M^{2}}$$
(3)

Minimal model:  $g = \kappa = 0$  [Jacobson et al; 2003]. Scale of Lorentz Violation:

$$\Lambda_{\rm LV} \equiv \frac{\rm M}{\sqrt{\xi}} \tag{4}$$

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Kinematical:

$$E^{2} = k^{2} + S \frac{k^{4}}{\Lambda_{LV}^{2}}, \ S = \pm 1$$
 (5)

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Kinematical:

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Dynamical:

$$\sum_{\lambda=1,2} \varepsilon_{\mu}^{\lambda} \varepsilon_{\nu}^{\lambda} = \operatorname{diag}\left(-\frac{\mathrm{E}_{\gamma}^{2}}{\mathrm{k}^{2}}, 1, 1, 1\right)$$
(6)

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$$\Phi_{\rm obs}(E) = e^{(-\tau_{\gamma}(E, z_{\rm s}))} \times \Phi_{\rm int}(E(1 + z_{\rm s}))$$
(7)

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$$\Phi_{\rm obs}(\mathbf{E}) = \mathbf{e}^{(-\tau_{\gamma}(\mathbf{E}, \mathbf{z}_{\rm s}))} \times \Phi_{\rm int}(\mathbf{E}(1 + \mathbf{z}_{\rm s})) \tag{7}$$

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$$\tau_{\gamma}(\mathbf{E}, \mathbf{z}_{s}) = \int_{0}^{\mathbf{z}_{s}} \mathrm{d}\mathbf{z} \frac{\mathrm{d}\mathbf{l}(\mathbf{z})}{\mathrm{d}\mathbf{z}} \int_{-1}^{1} \mathrm{d}\cos\theta \frac{1 - \cos(\theta)}{2} \int_{\varepsilon_{\mathrm{thr}}(\mathbf{E}(\mathbf{z}), \theta)}^{\infty} \mathrm{d}\varepsilon \mathbf{n}_{\gamma}(\varepsilon(\mathbf{z}), \mathbf{z}) \sigma_{\gamma\gamma}(\mathbf{E}(\mathbf{z}), \varepsilon(\mathbf{z}), \theta)$$
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For local observations (Milky Way):

$$\tau_{\gamma}(\mathbf{E}, \mathbf{z}_{\mathbf{s}}) = \underbrace{\int_{0}^{d} dl}_{d} \underbrace{\int_{-1}^{1} d\cos\theta \frac{1 - \cos(\theta)}{2} \int_{\varepsilon_{\mathrm{thr}}(\mathbf{E}, \theta)}^{\infty} d\varepsilon \mathbf{n}_{\gamma}(\varepsilon) \sigma_{\gamma\gamma}(\mathbf{E}, \varepsilon, \theta)}_{\frac{1}{\lambda_{\gamma}}} \tag{9}$$

$$\Phi_{\rm obs}(\mathbf{E}) = e^{(-\tau_{\gamma}(\mathbf{E}, \mathbf{z}_{\rm s}))} \times \Phi_{\rm int}(\mathbf{E}(1 + \mathbf{z}_{\rm s}))$$
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$$\Rightarrow \tau_{\gamma} \approx \frac{\mathrm{d}}{\lambda_{\gamma}} \qquad (10)$$

Pair production:

$$\gamma_{\rm VHE} + \gamma_{\rm soft} \to e^- + e^+$$
 (11)



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Consequences of modifications:

• Threshold condition

$$E\varepsilon(1 - \cos\theta) - \frac{E^4}{2\Lambda_{LV}^2} = 2m^2$$
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• Due to 
$$g_{\mu\nu} \rightarrow g_{\mu\nu} + \text{diag}(E_{\gamma}^2/\Lambda_{LV}^2, 0, 0, 0)$$
 we have

$$\langle |\mathbf{M}_{\mathrm{TOT}}|^2 \rangle = \langle |\mathbf{M}_{\mathrm{QED},\mathsf{A}}|^2 \rangle + \langle |\mathbf{M}_{\mathrm{MOD}}|^2 \rangle \tag{13}$$



Figure 1: Cross section for the subluminal case with  $\omega = 10^{-3}$  eV ,  $\theta = \pi$ .

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# Mean free path



• Both kinematical and dynamical modifications play a role

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- Connect to current experimental observations

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Work in progress:

- Understand uncertainties introduced by choosing different backgrounds
- Connect to current experimental observations
- Study the transparency in a Doubly Special Relativity (DSR) framework

- Vrban F.I.: Photon-photon scattering under Lorentz invariance violation, Master Thesis, University of Rijeka, 2022
- Cao Z., Aharonian F.A. et al.: Ultrahigh-energy photons up to 1.4 petaelectronvolts from 12  $\gamma$ -ray Galactic sources, Nature, 2021
- Rubtsov G., Satunin P., Sibiryakov S.: Calculation of cross sections in Lorentz-violating theories, Phys. Rev. D, 2012
- Jacobson T., Liberati S., Mattingly D.: Lorentz violation and Crab synchrotron emission: a new constraint far beyond the Planck scale, Nature, 2003
- Bolokhov A.P., Pospelov M.: Classification of dimension-5 Lorentz-violating interactions in the standard model, Phys. Rev. D, 2008

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