



# Finsler $(\alpha, \beta)$ -Gravitational Waves and their Observational Signature

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Based on joint work with Andrea Fuster - [arXiv:2302.08334](https://arxiv.org/abs/2302.08334)

# Outline

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- I. Finsler gravity
- II. Exact solutions
- III. Linearization  $\rightarrow$  gravitational waves (GWs)
- IV. Observational signature of such GWs
- V. Outlook

# Finsler Gravity

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**Finsler gravity** = GR formulated in the language of Finsler geometry

$$ds^2 = g_{\mu\nu}(x)dx^\mu dx^\nu \quad \rightarrow \quad ds = F(x, dx) \quad \text{'Finsler metric'}$$

**Vacuum field equation [1]**

$$\text{Ric} - \frac{F^2}{3} g^{\mu\nu} R_{\mu\nu} - \frac{F^2}{3} g^{\mu\nu} \left( \bar{\partial}_\mu \dot{S}_\nu - S_\mu S_\nu + \nabla_{\delta_\mu} S_\nu \right) = 0$$

# Exact Solutions

## Building blocks

$$\left\{ \begin{array}{l} \alpha = \sqrt{|A|} \quad \text{with } A = a_{\mu\nu} dx^\mu dx^\nu \text{ vacuum solution to EFEs} \\ \beta = b_\mu dx^\mu \quad \text{1-form that is covariantly constant, } \nabla_\mu b_\nu = 0 \end{array} \right.$$

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Prime example:  $\alpha = pp\text{-waves} = \sqrt{|-2du(dv + H(u, x, y)du) + dx^2 + dy^2|}$   
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**Randers metric**

Prime example:

$$ds = \alpha + \beta, \quad \phi(s) = 1 + s$$



# Linearization in 2 Ways

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1. Linearize in **departure from flatness**  $\rightarrow$  pp-wave becomes *standard linearized GW*:

$$\left\{ \begin{array}{l} \alpha = \sqrt{|-dt^2 + (1 + f_+(t-z))dx^2 + (1 - f_+(t-z))dy^2 + 2f_x(t-z)dx dy + dz^2|} \\ \beta = \frac{\lambda}{\sqrt{2}} (dt - dz) \end{array} \right.$$

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2. Linearize in **departure from GR** (i.e. 'size' of 1-form,  $\lambda$ )  $\rightarrow$  Generic  $(\alpha, \beta)$ -metric becomes *Randers metric*

$$ds = \alpha \phi \left( \frac{\beta}{\alpha} \right) \approx \alpha \left( \phi(0) + \phi'(0) \frac{\beta}{\alpha} \right) = \alpha \phi(0) + \phi'(0) \beta = \tilde{\alpha} + \tilde{\beta}$$

# Finsler Gravitational Waves

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**Conclusion:**

To 1<sup>st</sup> order, vacuum solutions are given by Randers metrics  $ds = \alpha + \beta$  with

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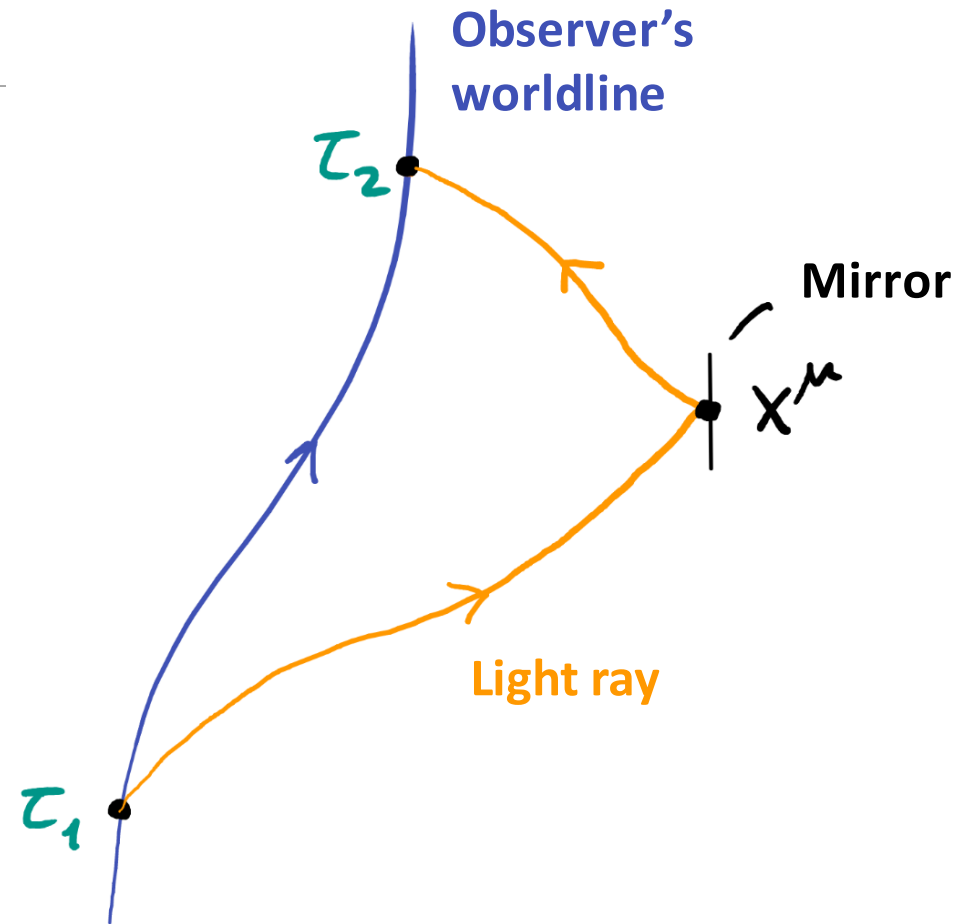
Observational signature...?

# Radar Distance

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**Radar distance** =  $\frac{1}{2}$  time between emission and reception of a reflected light ray

$$R(x) = \frac{\Delta\tau}{2} = \frac{\tau_2 - \tau_1}{2}$$



**GR** radar distance [2]:

$$R = \Delta\ell + \left( \frac{\Delta x^2 - \Delta y^2}{4\Delta\ell} \right) \bar{f}_+ + \left( \frac{\Delta x \Delta y}{2\Delta\ell} \right) \bar{f}_\times$$

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However, define the **measurable quantities**:

$\Delta X$  = radar distance in x-direction - *in the absence of the GW*

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$\Rightarrow$  **Identical to GWs in GR!**

# Conclusion & Outlook

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- The Finslerian GWs are **indistinguishable** from standard GWs in GR **with the same waveform**.
- However, Finsler effects might affect the **generation** of GWs
  - Different waveforms
- Also, different types of Finslerian gravitational waves probably exist.

Thank you!



~~Randers~~

$$ds = \alpha + \beta$$

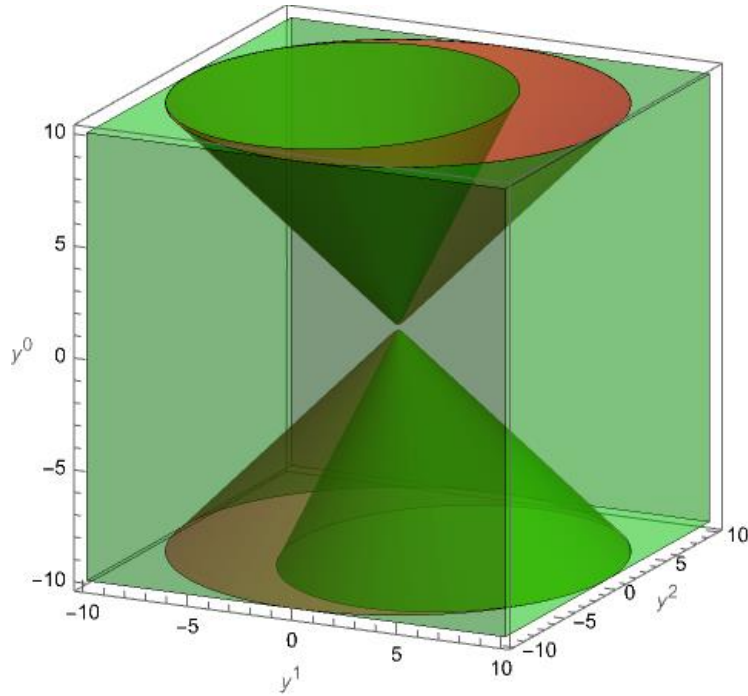


Modified Randers

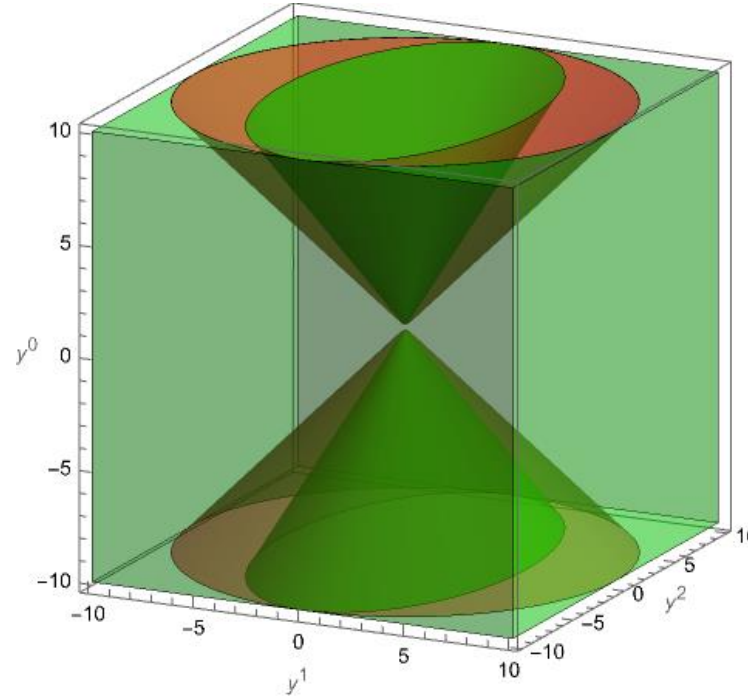
$$ds = \text{sign}(A)\alpha + |\beta|$$

where

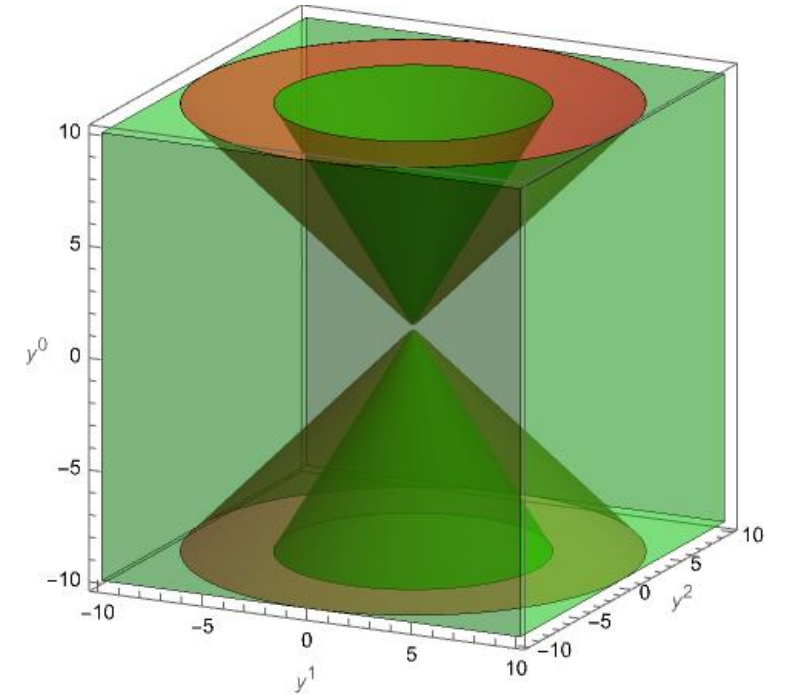
$$\begin{cases} A = a_{\mu\nu} dx^\mu dx^\nu \\ \alpha = \sqrt{|A|} \end{cases}$$



Lightlike  $\beta$



Spacelike  $\beta$



Timelike  $\beta$

Green = Lorentzian signature

Red = positive definite signature

# References

M. Rakhmanov, On the round-trip time for a photon propagating in the field of a plane gravitational wave, *Classical and Quantum Gravity* 26, 155010 (2009).

C. Pfeifer and M. N. R. Wohlfarth, Finsler geometric extension of Einstein gravity, *Phys.Rev. D* 85, 064009 (2012).