



Testing CPT with the solar neutrino sector

Pablo Martínez Miravé (IFIC, CSIC - Univ. Valencia) COST CA18108 Fourth Annual Conference Based on 2305.06384 [hep-ph] in collaboration with G. Barenboim, C.A. Ternes and M. Tórtola









CPT violation could manifest as **particles and antiparticles** having different masses and lifetimes.

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- I. With data from solar neutrinos, we set **new bounds**.
- II. Future neutrino observatories, including JUNO, Hyper-Kamiokande and DUNE, will further improve these limits.

About CPT symmetry

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About CPT symmetry

CPT is a keystone of high energy physics which stems from requiring **Lorentz invariance**, **unitarity and locality**.

CPT violation could result from

- Lorentz violation
- Non-local Lorentz-invariant field theories

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Addazi et al, Prog.Part.Nucl.Phys. 125 (2022) 103948

R. Jost, Helv.Phys.Acta 30 (1957) 409-416

About CPT symmetry

From **CPT Theorem**, particles and antiparticles have the **same mass and lifetime**.

In **some CPT breaking models**, a mass splitting between particles and antiparticles is realised.

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Chaichian et al, Phys.Lett.B 699 (2011)

Chaichian et al, Eur.Phys.J.C 73 (2013) 3, 2349

About CPT Symmetry and neutrino oscillations

Neutrino oscillations are sensitive to two mass splittings

$$\Delta m_{21}^2$$
 and Δm_{31}^2 $(\Delta m_{ij}^2 = m_i^2 - m_j^2)$

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three mixing angles

, $\theta_{12},\,\theta_{13}$ and θ_{23} and a phase $\,\delta_{\rm CP}$.

About CPT symmetry and neutrino oscillations

Neutrino oscillations are sensitive to two mass splittings

$$\begin{array}{l} \Delta m_{21}^2 \text{ and } \Delta m_{31}^2 & (\Delta m_{ij}^2 = m_i^2 - m_i^2) \\ \text{ three mixing angles} \\ & \theta_{12}, \ \theta_{13} \ \text{and} \ \theta_{23} \\ \text{ and a phase } \delta_{\text{CP}}. \end{array}$$

SOLAR SECTOR

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About solar neutrinos

Solar neutrino **flux**



Solar neutrino electron survival probability



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About reactor antineutrinos

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Reactor antineutrino electron survival probability



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Reactor antineutrino electron survival probability



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CURRENT STATUS

Limits based on data from 2020

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G. Barenboim et al. JHEP 07 (2020) 155

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$$\begin{aligned} |\Delta(\Delta m_{21}^2)| &= |\Delta m_{21}^2 - \Delta \bar{m}_{21}^2| < 4.7 \times 10^{-5} \text{ eV}^2 \\ |\Delta(\Delta m_{31}^2)| &= |\Delta m_{31}^2 - \Delta \bar{m}_{31}^2| < 2.5 \times 10^{-4} \text{ eV}^2 \\ |\Delta(\sin^2 \theta_{12})| &= |\sin^2 \theta_{12} - \sin^2 \bar{\theta}_{12}| < 0.14 \\ |\Delta(\sin^2 \theta_{13})| &= |\sin^2 \theta_{13} - \sin^2 \bar{\theta}_{13}| < 0.029 \\ |\Delta(\sin^2 \theta_{23})| &= |\sin^2 \theta_{23} - \sin^2 \bar{\theta}_{23}| < 0.19 \end{aligned}$$

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CURRENT STATUS

Updated limits

 $\begin{aligned} |\Delta(\Delta m_{21}^2)| &= |\Delta m_{21}^2 - \Delta \bar{m}_{21}^2| < 3.7 \times 10^{-5} \text{ eV}^2 \\ |\Delta(\Delta m_{31}^2)| &= |\Delta m_{31}^2 - \Delta \bar{m}_{31}^2| < 2.5 \times 10^{-4} \text{ eV}^2 \\ |\Delta(\sin^2 \theta_{12})| &= |\sin^2 \theta_{12} - \sin^2 \bar{\theta}_{12}| < 0.187 \\ |\Delta(\sin^2 \theta_{13})| &= |\sin^2 \theta_{13} - \sin^2 \bar{\theta}_{13}| < 0.029 \\ |\Delta(\sin^2 \theta_{23})| &= |\sin^2 \theta_{23} - \sin^2 \bar{\theta}_{23}| < 0.19 \end{aligned}$

CURRENT STATUS



FUTURE PROSPECTS

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JUNO: medium baseline reactor antineutrinos

Hyper-Kamiokande: high-energy solar neutrinos

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DUNE: high-energy solar **neutrinos**

• • • Extremely **sensitive.**

We consider two sets • of experimental details: **conservative** and **optimal** for each experiment.

FUTURE PROSPECTS

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FUTURE PROSPECTS

From **neutrinos**:

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•	$ \Delta \sin^2 \theta_{12} $	$ \Delta(\Delta m_{21}^2) [10^{-5} eV^2]$
current bound	0.187	3.7
JUNO + HK + DUNE conservative	0.018	2.4
JUNO + HK + DUNE optimal	0.011	0.8

Current limit from the kaon system:

$$|m^2(K^0) - m^2(\overline{K}^0)| < 0.25 \text{ eV}^2$$

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FUTURE PROSPECTS

Testing CPT with the solar neutrino sector

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SPARE SLIDES

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Hyper-Kamiokande

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Factor 2 improvement with respect to Super-Kamiokande IV

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Measuring CPT violation



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Disentangling CPT and NSI

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Solar neutrino detection



 $\nu_x + e^- \longrightarrow \nu_x + e^ \nu_e + {}^{40}\text{Ar} \longrightarrow {}^{40}\text{K}^* + e^-$

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