Effect of the cosmological model on LIV constraints from GRB Time-Delays datasets

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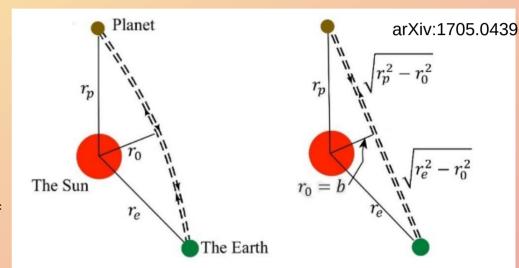




Time delays - overview

- Classical TD from interaction between gravity and particles and fields
 - -- **Observed** in Shapiro delay, graviational and cosmological redshift
 - -- Independent of the energy of the photons
- Quantum Gravity TD from String theory, Loop OG, etc
 - -- modified dispersion relations, i.e. dependent of the energy of the photon
 - -- should affect all messengers (neutrino, GW, photons)
- -- Strict constraints from GRBs (E_{QG}>7xE_{Pl}, Vasileiou et al 2013, E_{QG}>0.5xE_{Pl} Acciari et al 2020)

$$E^2 = p^2 c^2 \left[1 - s_{\pm} \left(\frac{E}{\xi_n E_{OG}} \right)^n \right] ,$$



Time delays in GRBs

Gamma-Ray Bursts

- -- high energies (E_{iso}>10⁵²erg)
- -- high redshifts (z~9)
- -- numerous observations

A perfect probe for TD but a lot of difficulties

- -- Progenitor model vs GRB light curve
- -- Intrinsic vs propagational effects
- -- What about the cosmology?

$$t_{LIV} = \int_0^z [1 + \frac{E}{E_{QG}}(1 + z')] \frac{dz'}{H(z')}$$

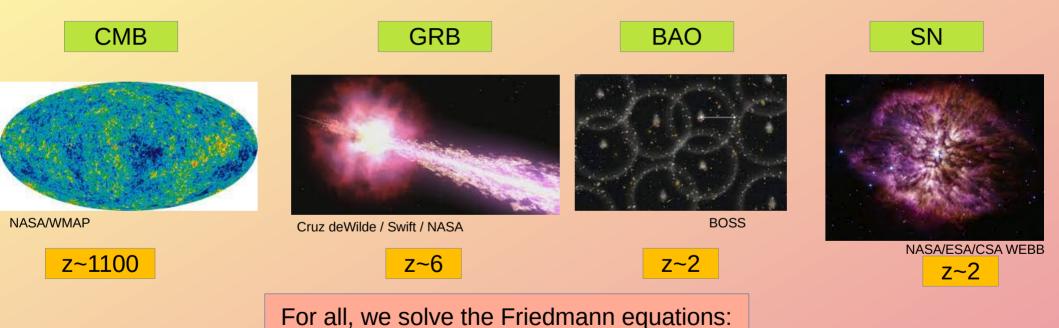
$$\frac{\Delta t_{obs}}{1+z} = a_{LIV}K + \beta \,,$$

$$K \equiv \frac{1}{1+z} \int_0^z \frac{(1+\tilde{z}) d\tilde{z}}{h(\tilde{z})}.$$

$$a_{LIV} \equiv \Delta E / (H_0 E_{QG})$$

$$\Delta t_{\rm obs} = \Delta t_{\rm int} + \Delta t_{\rm QG} + \Delta t_{\rm spec} + \Delta t_{\rm DM} + \Delta t_{\rm grad}$$

To investigate cosmology, we combine combine GRB TD data with other astrophysical sources

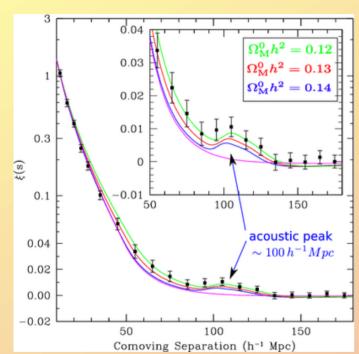


 $H(z)/H_0 = E(z)$ $E(z)^2 = \Omega_r(1+z)^4 + \Omega_m(1+z)^3 + \Omega_k(1+z)^2 + \Omega_{DE}(z)$,

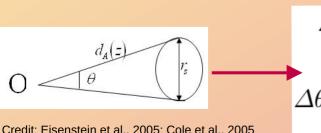
The DDE models we consider

Model	$\Omega_{DE}(z) = \Omega_{\Lambda} \times$	w(z)
CPL	$\exp\left[\int_0^z \frac{3(1+w(z'))dz'}{1+z'}\right]$	$w_0 + w_a \frac{z}{z+1}$
BA	$(1+z)^{3(1+w_0)}(1+z^2)^{\frac{3w_1}{2}}$	$w_0 + z \frac{1+z}{1+z^2} w_1$
LC	$(1+z)^{(3(1-2w_0+3wa))}e^{\frac{9(w_0-wa)z}{(1+z)}}$	$\frac{(-z+z_c)w_0+z(1+z_c)w_c}{(1+z)z_c}$
JPB	$(1+z)^{3(1+w_0)}e^{\frac{3w_1z^2}{2(1+z)^2}}$	$w_0 + w_1 \frac{z}{(1+z)^2}$
FSLLI	$(1+z)^{3(1+w_0)}e^{\frac{3w_1}{2}\arctan(z)}(1+z^2)^{\frac{3w_1}{4}}(1+z)^{-\frac{3}{2}w_1}$	$w_0 + w_1 \frac{z}{1+z^2}$
FSLLII	$(1+z)^{3(1+w_0)}e^{-\frac{3w_1}{2}\arctan(z)}(1+z^2)^{\frac{3w_1}{4}}(1+z)^{+\frac{3}{2}w_1}$	$w_0 + w_1 \frac{z^2}{1+z^2}$
PEDE	$\frac{1-\tanh(\bar{\Delta}\log_{10}(\frac{1+z}{1+z_t}))}{1+\tanh(\bar{\Delta}\log_{10}(1+z_t)}$	$-\tfrac{(1+\tanh[\log_{10}{(1+z)]})}{3\ln 10}-1$

BAO – "standard ruler" in cosmology



- Baryonic acoustic oscilations are periodic fluctuations in the density of the visible baryonic matter of the universe.
- Created by the intrerplay of gravity, radiative pressure and the expansion of the universe
- The distance at which plasma waves induced by radiation pressure froze at recombination the sound horizon, r_d (Planck 2018: r_d =147.5 Mpc, z_d =1059, z_* =1100)
- Measured by looking at the large scale structure of matter



 $\Delta z = r_d H(z)/c$

 $D_M = \frac{c}{H_0} S_k \left(\int_0^z \frac{dz'}{E(z')} \right)$

Credit: Eisenstein et al., 2005; Cole et al., 2005

The quantities we use

SN/GRB

$$\mu_B(z) - M_B = 5 \log_{10} [d_L(z)] + 25,$$

CMB distance priors

$$l_{\rm A} = (1 + z_*) \frac{\pi D_{\rm A}(z_*)}{r_s(z_*)},$$

$$R \equiv (1 + z_*) \frac{D_{\rm A}(z_*) \sqrt{\Omega_m} H_0}{c},$$

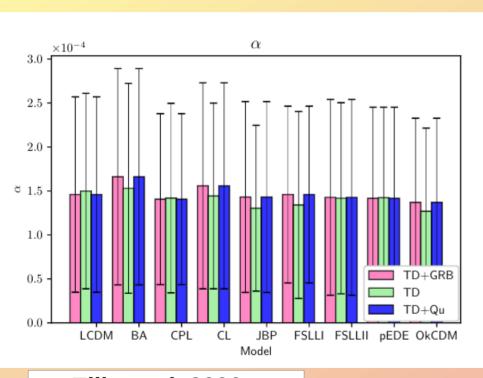
BAO

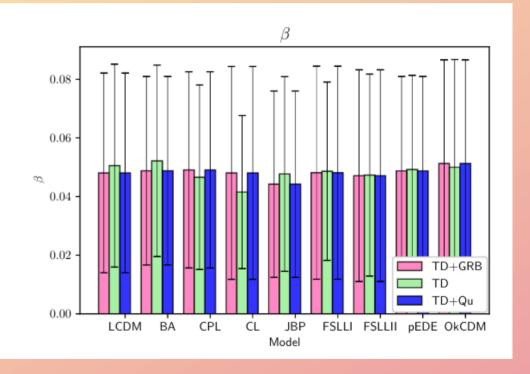
$$D_{A} = \frac{c}{(1+z)H_{0}\sqrt{|\Omega_{k}|}} \operatorname{sinn}\left[|\Omega_{k}|^{1/2} \int_{0}^{z} \frac{dz'}{E(z')}\right]$$

where $S_k(x) = \begin{cases} \frac{1}{\sqrt{\Omega_k}} \sinh\left(\sqrt{\Omega_k}x\right) & \text{if } \Omega_k > 0\\ x & \text{if } \Omega_k = 0\\ \frac{1}{\sqrt{-\Omega_k}} \sin\left(\sqrt{-\Omega_k}x\right) & \text{if } \Omega_k < 0 \end{cases}$

All depend on c/H₀r_d so we take it as 1 factor!

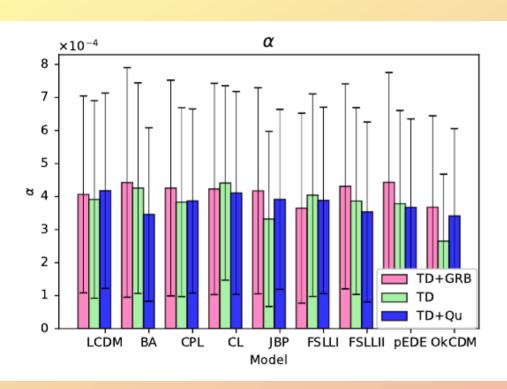
The results - TD1

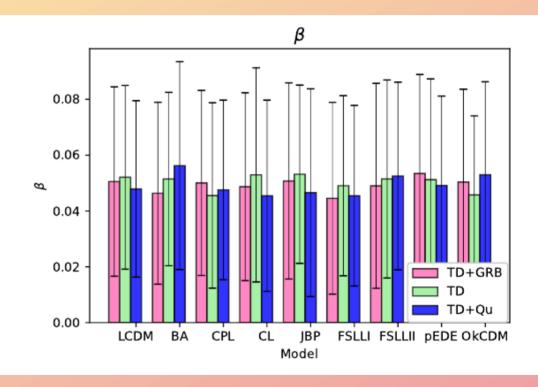




Ellis et al. 2006, 35 GRBs, wavelet method

The results - TD2





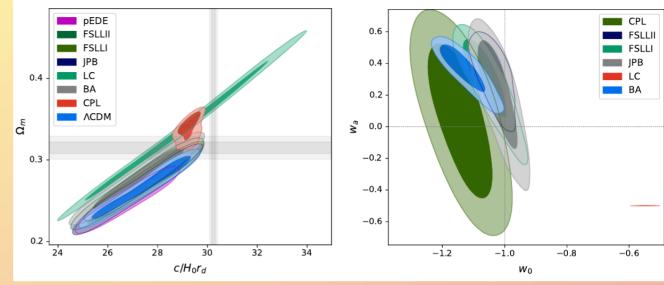
Vardanyan et al. 2022, 49 GRBs, descrete CCF

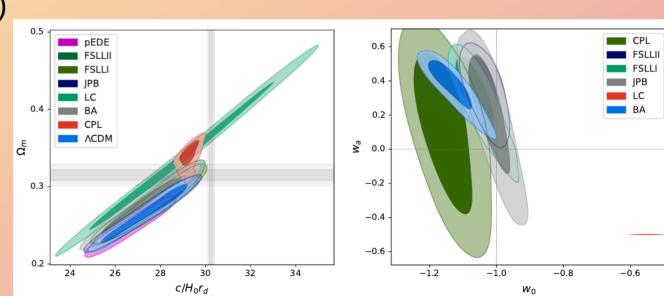
In both cases - we see deviations between DE models TD1

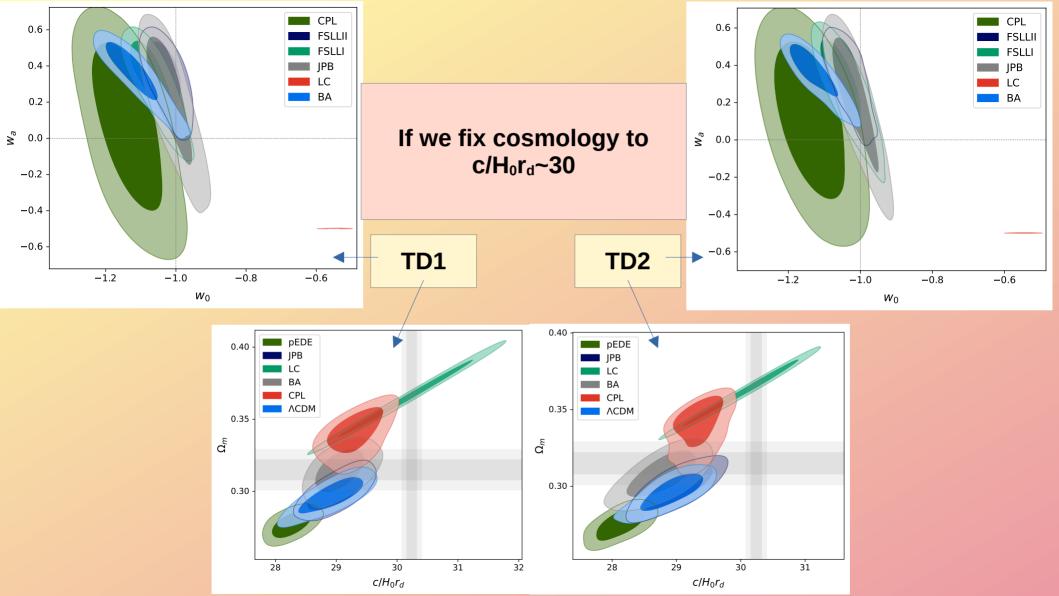
The cosmology TD1 and TD2 look very similar)

In both cases c/H₀r๙~28!!!

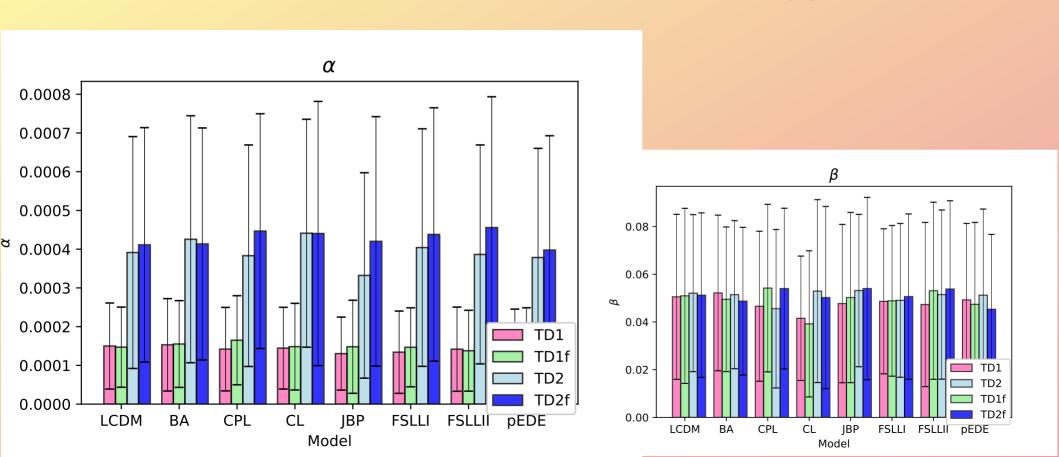
TD2



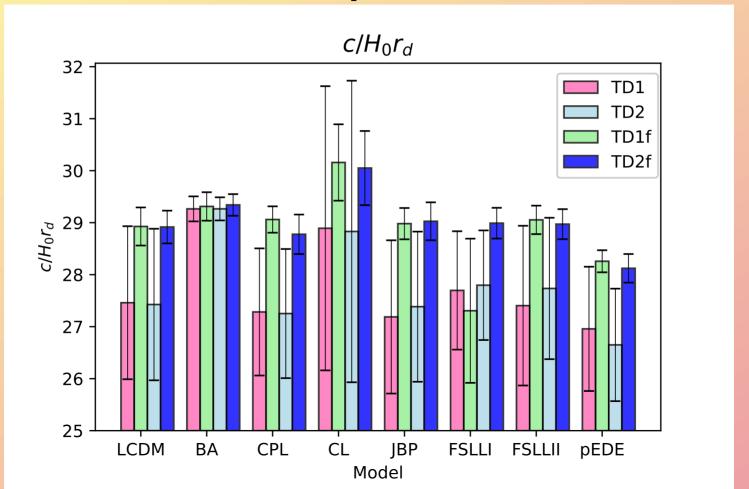




Or if we compare fixed cosmology with non-fixed cosmology...



The c/H₀r_d parameter



Conclusions:

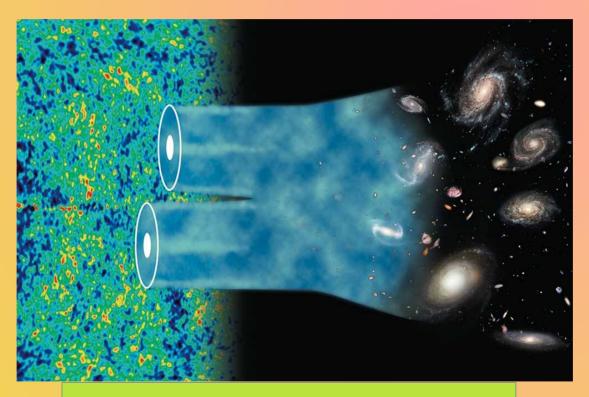
Variation within the model larger than variation between models! Cosmology can account for at least 10% of any constraint on LIV effects! We need better understanding of intrinsic effects!



- $c/H_0r_d\sim 27$ lower than pure DDE $c/H_0r_d\sim 30$
- Lower than Planck's matter density
- Similar DDE parameters
- Some limited preference for CPL, BA, **FSLLII**
- TD1 $E_{QG} > 5 \times 10^{17} \; \mathrm{GeV}$ TD2 $E_{QG} > 1.1 \times 10^{17} \; \mathrm{GeV}$
- 20%-60% deviation due to cosmology
- Consistent with previous results

"Effect of the cosmological model on LIV constraints from GRB Time-Delays datasets,, DS, arXiv:2305.06504

Thank you for your attention!



Thanks to COST CA18108 for the financial support and the **inspiration!**

Datasets/Methods

- Two GRB TD datasets
 (TD1 35 pts, TD2 49 pts)
- CMB distance prior (2 pts)
- BAO (15 pts)
- SN (40 binned pts)
- GRB/Quasars (162/24 pts)

- MCMC with Polychord
- Marginalized
 - SN and GRB
- DDE models
- We remove

$$F = \sum_{i} C_{cov}^{-1},$$

$$\chi_{BAO}^{2} = \sum_{i} \frac{\left(\vec{v}_{obs} - \vec{v}_{model}\right)^{2}}{\sigma^{2}},$$

$$\chi_{TD}^{2} = \sum_{i} \frac{\left(\vec{\Delta t}_{obs} - \vec{\Delta t}_{model}\right)^{2}}{\sigma^{2}},$$

 $\tilde{\chi}_{SN,GRB}^2 = D - \frac{E^2}{F} + \ln \frac{F}{2\pi}$

 $D = \sum_{i} \left(\Delta \mu \, C_{cov}^{-1} \, \Delta \mu^{T} \right)^{2},$

 $E = \sum_{i} \left(\Delta \mu \, C_{cov}^{-1} \, E \right),\,$

 $\chi^2 = \chi^2_{BAO} + \chi^2_{CMB} + \chi^2_{SN} + \chi^2_{TD}(+\chi^2_{GRB})(+\chi^2_{Qua}).$