The background features a complex, abstract pattern of overlapping, glowing orbits and particles in various colors (purple, blue, orange, yellow) against a dark, starry space. The orbits are represented by thin, curved lines, and the particles are small, bright spheres of different sizes and colors. The overall effect is reminiscent of a dynamic, multi-colored atomic model or a visualization of celestial mechanics.

Finsler spacetimes with (α, β) -metrics and their isometries

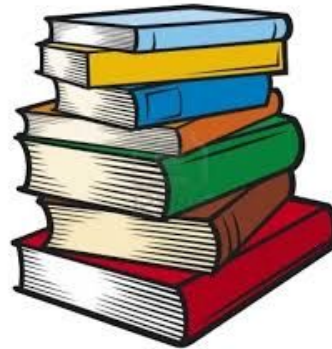
Nicoleta VOICU

Transilvania University of Brasov, Romania

joint work with: E. Popovici-Popescu, A. Friedl-Szasz, C. Pfeifer

Talk based on:

N. Voicu, A. Friedl-Szasz, E. Popovici-Popescu, C. Pfeifer, *The Finsler space-time condition for (α, β) -metrics and their isometries*, Universe 9, 198 (2023).



1 Main results

★ (α, β) -metrics:

- most commonly used class of Finsler metrics
- obtained by deforming a Riemannian metric (α) by a 1-form (β)

For general Finsler metrics of (α, β) -type:

- ✓ We find a set of (minimal) conditions such that these admit a *well defined causal structure*.
- ✓ We determine all Killing vector fields...

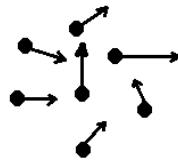


2 Motivation

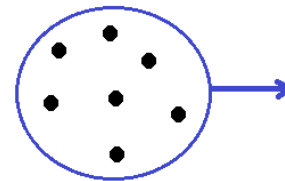
Modified dispersion relations \rightarrow Finslerian arc length

Also:

- ✓ SME models (Kostelecky&co.);
- ✓ minimal extension of Ehlers-Pirani-Schild axioms (Perlick&Laemmerzahl);
- ✓ more accurate description of gravity, taking into account *multiple sources, with different velocities* - **kinetic gas** (Hohmann, Pfeifer, NV, 2020)



kinetic gas
(individual velocities)



fluid
(averaged velocity)

3 Finsler gravity

General relativity $g_{\mu\nu}(x)$	Finsler gravity $L(x, \dot{x})$
$ds^2 = g_{\mu\nu}(x) dx^\mu dx^\nu$	$ds^2 = L(x, dx)$
$\frac{d^2 x^\mu}{ds^2} + \Gamma^\mu_{\nu\sigma}(x) \frac{d\dot{x}^\nu}{ds} \frac{d\dot{x}^\sigma}{ds} = 0$	$\frac{d^2 x^\mu}{ds^2} + N^\mu_\nu \left(x, \frac{d\dot{x}}{ds}\right) \frac{d\dot{x}^\nu}{ds} = 0$
$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \kappa T_{\mu\nu}$	$\frac{1}{2} g^{\mu\nu} \frac{\partial^2 Ric}{\partial \dot{x}^\mu \partial \dot{x}^\nu} - 3L^{-1} Ric - \text{div } P = \kappa \phi$

✓ Finslerian Ricci scalar Ric :

$$Ric := R_{\mu\nu} \dot{x}^\mu \dot{x}^\nu \quad (1)$$

Take: M - 4-dim manifold.

Finsler spacetime structure on M : a smooth scalar function

$$L = L(x, \dot{x}) \quad (2)$$

on some subset of TM , with:

(i) 2-homogeneity: $L(x, \alpha\dot{x}) = \alpha^2 L(x, \dot{x}), \forall \alpha > 0$;

(ii) Future timelike cones \mathcal{T}_x

at each $x \in M$, \exists a connected conic set $\mathcal{T}_x \neq \emptyset$ such that:

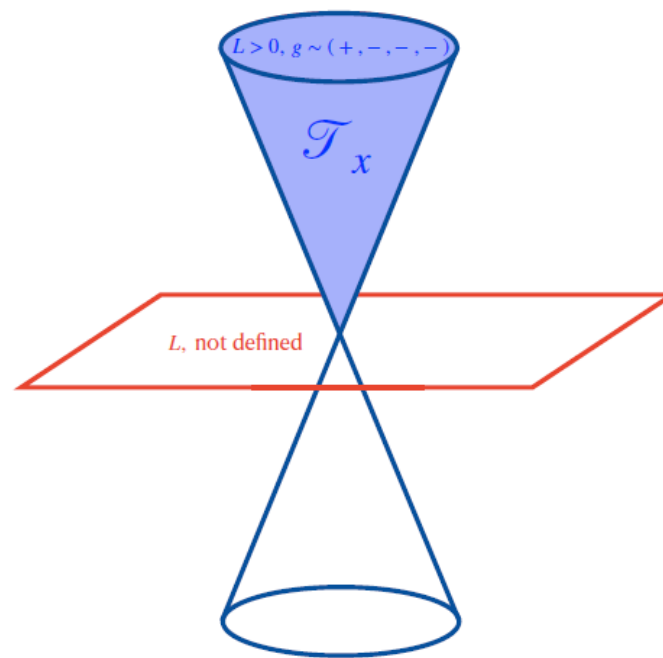
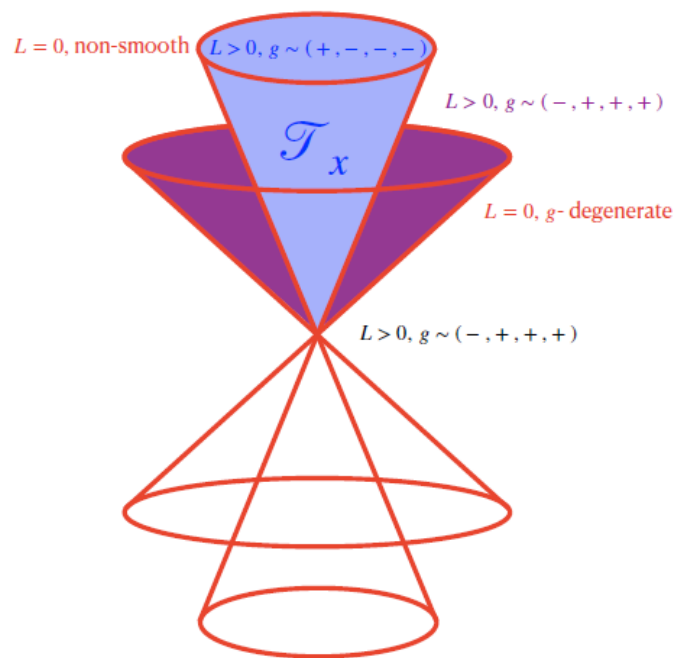
✓ $L > 0$ on \mathcal{T}_x , $L|_{\partial\mathcal{T}_x} = 0$

✓ The matrix

$$g_{\mu\nu}(x, \dot{x}) = \frac{1}{2} \frac{\partial^2 L}{\partial \dot{x}^\mu \partial \dot{x}^\nu}$$

has Lorentzian signature $(+, -, -, -)$.

Finslerian cones:



4 (α, β) -metric spacetimes

Setting: $(M, a_{\mu\nu}(x))$ - Lorentzian (GR) spacetime, $b := b_\mu(x) dx^\mu$.

$$A(x, \dot{x}) : = a_{\mu\nu}(x) \dot{x}^\mu \dot{x}^\nu, \quad (3)$$

$$B(x, \dot{x}) : = b_\mu(x) \dot{x}^\mu. \quad (4)$$

Definition: $L = L(x, \dot{x})$ is an (α, β) -metric if:

$$L = A\Psi(s), \quad s = \frac{B^2}{A}. \quad (5)$$

Traditional notations:

$$|A| = \alpha^2, \quad B = \beta \quad (6)$$

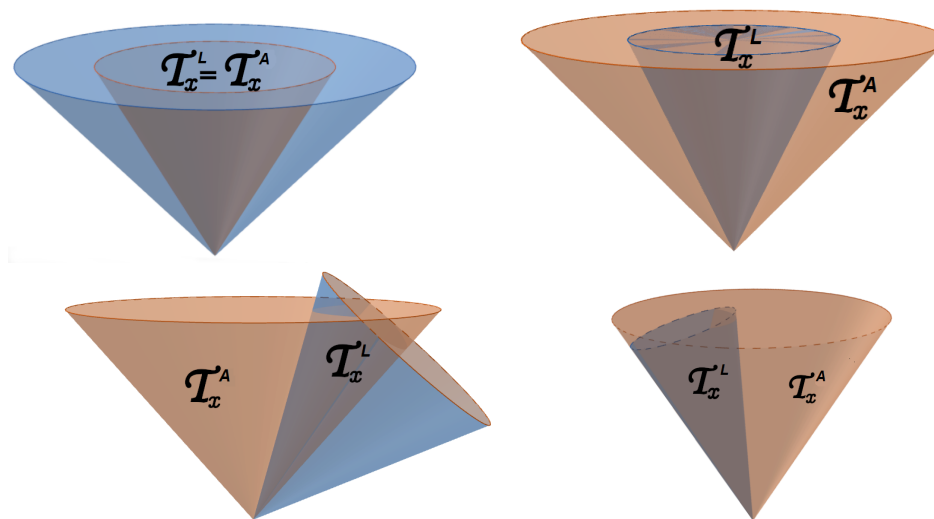
Assumption:

$$\mathcal{T}_x^A \cap \mathcal{T}_x^L \neq \emptyset. \quad (7)$$

Then:

$$\mathcal{T}_x^L : A > 0, \quad \psi > 0. \quad (8)$$

$$\partial\mathcal{T}_x^L : A = 0 \quad \text{or} \quad \psi = 0. \quad (9)$$



Determinant of $g_{\mu\nu}(x, \dot{x})$:

$$\det(g_{\mu\nu}) = \Psi^2(\Psi - s\Psi') \det(a) \frac{d}{ds} [(s - \langle b, b \rangle)\sigma], \quad (10)$$

where:

$$\sigma := \frac{(\Psi - s\Psi')^2}{\Psi} \quad (11)$$

Theorem 1: $L = A\Psi(s)$ - Finsler spacetime structure \Leftrightarrow at each $x \in M$, \exists a connected conic set \mathcal{T}_x such that:

(i) $A > 0$, $\Psi > 0$ on \mathcal{T}_x and $\lim_{\dot{x} \rightarrow \partial\mathcal{T}_x} (A\Psi) = 0$.

(ii) at all $\dot{x} \in \mathcal{T}_x$:

$$\Psi - s\Psi' > 0, \quad (s - \langle b, b \rangle) \frac{d}{ds} \ln \sigma > -1. \quad (12)$$

5 Examples

1. Randers metrics

$$L = \epsilon(\sqrt{|A|} + B)^2 = A(1 - \sqrt{s})^2. \quad (13)$$

(where $\epsilon = \text{sgn}(A)$)

Proposition 1: L defines a Finsler spacetime structure \Leftrightarrow

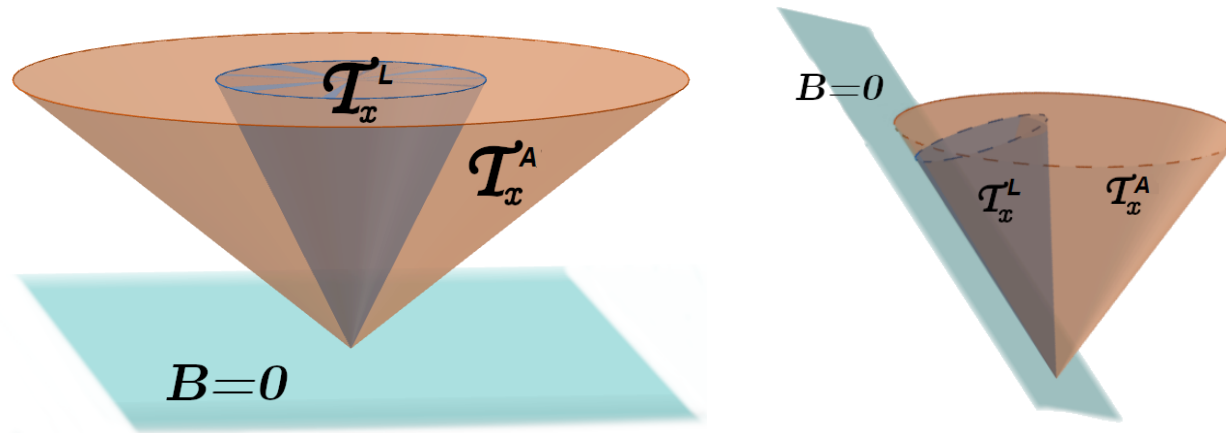
$$\langle b, b \rangle \in [0, 1). \quad (14)$$

In this case:

$$\mathcal{T}_x^L = \{\dot{x} \mid (a_{\mu\nu} - b_\mu b_\nu)\dot{x}^\mu \dot{x}^\nu > 0\} \cap \mathcal{T}_x^A. \quad (15)$$

Properties:

1. \mathcal{T}_x^L - "sharper" cone than \mathcal{T}_x^A .
2. on \mathcal{T}_x^L : $B < 0$.



Variant: Modified Randers metrics (Heefer&Fuster 2023):

$$L = \epsilon(\sqrt{|A|} - |B|)^2 = A(1 - \sqrt{s})^2 \quad (16)$$

→ allows for

$$\langle b, b \rangle \in (-1, 1) \quad (17)$$

(and has well defined *past timelike cones*).

2. Bogoslovsky-Kropina (VSR/VGR, m-Kropina) metrics

$$L = A^{1-q} B^{2q} = A s^q. \quad (18)$$

Proposition 2: L - Finsler spacetime structure in precisely one of the following cases:

(i) $\langle b, b \rangle > 0$, $q \in [-1, 1)$. In this case:

$$\mathcal{T}_x^L = \mathcal{T}_x^A. \quad (19)$$

(ii) $\langle b, b \rangle = 0$, $q \in (-1, 1)$. In this case:

$$\mathcal{T}_x^L = \mathcal{T}_x^A. \quad (20)$$

(iii) $\langle b, b \rangle < 0$, $q \in (0, 1)$. In this case:

$$\mathcal{T}_x^L = \mathcal{T}_x^A \cap \{B > 0\}. \quad (21)$$

3. Generalized m-Kropina metrics:

$$L(x, \dot{x}) = As^{-p} (k + ms)^{p+1}, \quad (22)$$

where $p, k, m \in \mathbb{R}$ - const.

Proposition 3: L - Finsler spacetime structure if and only if

$$k > 0, \quad m < 0 \quad (23)$$

and one of the following occurs:

(i) $\langle b, b \rangle \geq 0$, $p \in (-1, 1]$. In this case:

$$\mathcal{T}_x^L = \mathcal{T}_x^A \cap \{kA + mB^2 > 0\}. \quad (24)$$

(ii) $\langle b, b \rangle < 0$, $p \in (-1, 0)$. In this case:

$$\mathcal{T}_x^L = \mathcal{T}_x^A \cap \{kA + mB^2 > 0\} \cap \{B > 0\}. \quad (25)$$

4. Exponential metrics:

$$L = Ae^{P(s)}. \quad (26)$$

Proposition 4: L - spacetime structure iff :

$$\lim_{s \rightarrow \infty} \frac{e^{P(s)}}{s} = 0, \quad (27)$$

$$1 - sP' > 0, \quad 1 - sP' > (s - \langle b, b \rangle) \left(s (P')^2 + 2sP'' + P' \right), \quad \forall s \in I.$$

✓ **Concrete example:**

$$A = \eta_{\mu\nu} \dot{x}^\mu \dot{x}^\nu, B = \langle b, b \rangle \dot{x}^0$$

$$\Psi : [\langle b, b \rangle, \infty) \rightarrow \mathbb{R}_+, \quad \Psi(s) = e^{k - \frac{\langle b, b \rangle^2}{2s^2}}.$$

6 Isometries of (α, β) -metrics

Isometry (symmetry):

$$\varphi : M \rightarrow M, \quad d\varphi(L) = 0. \quad (28)$$

Generators: $\xi \in \mathcal{X}(M)$ - **Killing vector field (symmetry)** for L :

$$\xi^C(L) = 0, \quad (29)$$

where $\xi^C = \xi^\mu \partial_\mu + \xi^\mu_{,\nu} \dot{x}^\nu \dot{\partial}_\mu$ - complete lift of ξ to TM .

For (α, β) -metrics

$$L = A\Psi(s), \quad s = \frac{B^2}{A}, \quad (30)$$

\Rightarrow the Killing vector condition:

$$\xi^C(A) (\Psi - s\Psi') + 2\Psi' B \xi^C(B) = 0, \quad (31)$$

where

$$\xi^C(A) = (\mathfrak{L}_\xi a_{ij}) \dot{x}^i \dot{x}^j, \quad \xi^C(B) = (\mathfrak{L}_\xi b_i) \dot{x}^i.$$

Classes of symmetries:

1. **Trivial symmetries:** ξ - Killing vector field for A :

$$\xi^C(A) = 0 \quad (\Rightarrow \xi^C(B) = 0). \quad (32)$$

2. **Nontrivial symmetries:** $\xi^C(A) \neq 0$.

Theorem 2: $L = A\Psi(s)$ admits nontrivial symmetries $\xi \Leftrightarrow \exists \kappa, m_1, m_2, n_1, n_2$ - smooth functions of x only (with κ not identically zero) such that:

$$\Psi(s) = \begin{cases} cs^{\frac{m_1}{m_1-2n_2}} |m_2s + m_1 - 2n_2|^{\frac{-2n_2}{m_1-2n_2}}, & \text{if } m_1 \neq 2n_2 \\ cse^{\frac{-2n_2}{m_2s}}, & \text{if } m_1 = 2n_2, \quad m_2 \neq 0 \end{cases} \quad (33)$$

and:

$$\xi^C(B) = \kappa n_2 B, \quad \xi^C(A) = \kappa (m_1 A + m_2 B^2). \quad (34)$$

✓ **Concrete example: (curved) Bogoslovsky-Kropina metric on \mathbb{R}^4 :**

$$A = e^{2qx^0} \eta_{\mu\nu} \dot{x}^\mu \dot{x}^\nu, \quad B = e^{(q-1)x^0} (\dot{x}^0 + \dot{x}^1),$$

$$L = A \left(\frac{B^2}{A} \right)^q, \quad q \in (0, 1).$$

That is: $n_2 = q - 1$, $m_1 = 2q$, $m_2 = 0$, $\kappa = c = 1$. Then

$$\xi := (1, 0, 0, 0). \tag{35}$$

obeys:

$$\xi^C(A) = 2qA \neq 0, \quad \xi^C(B) = (q - 1)B$$

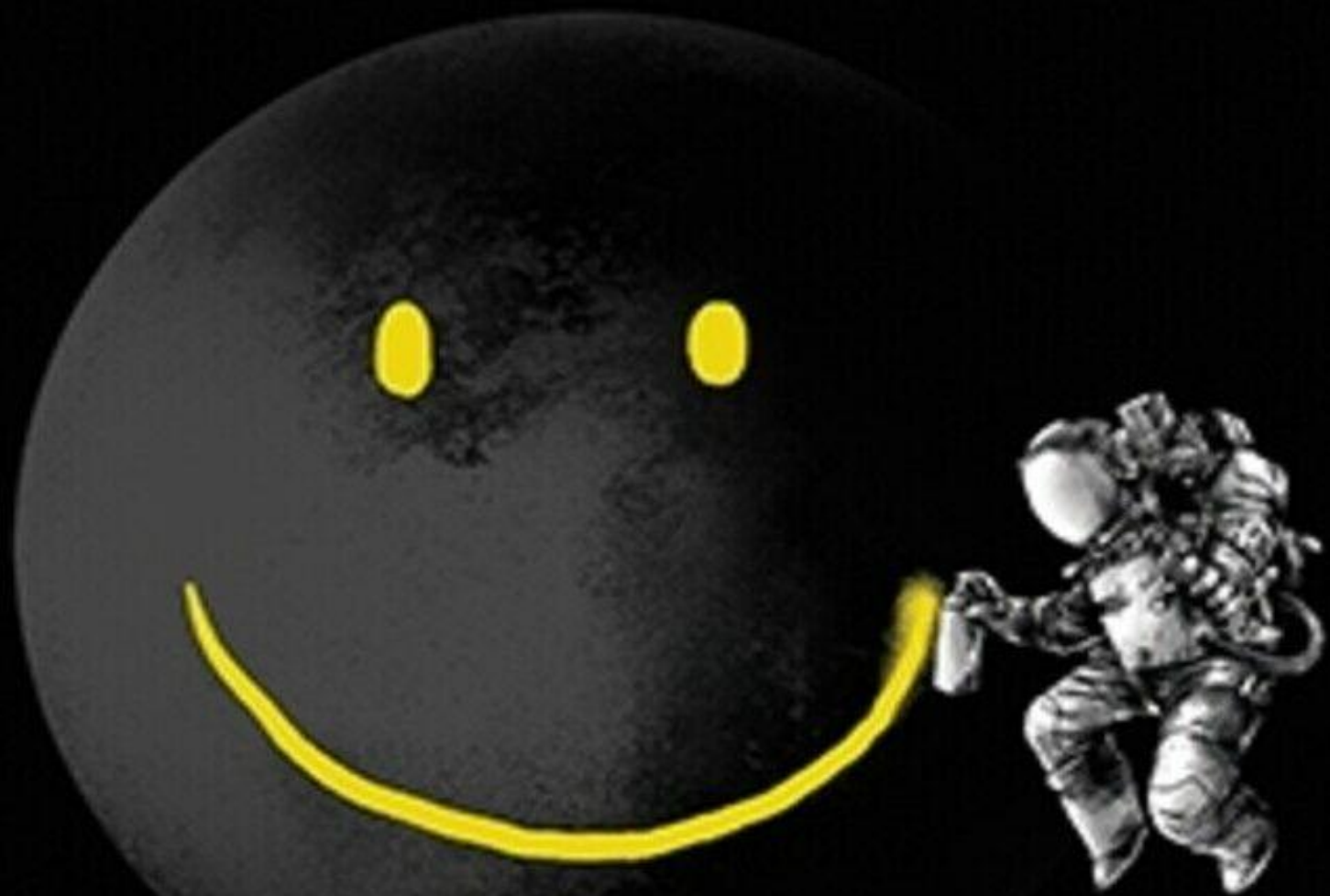
$\Rightarrow \xi$ is a nontrivial Killing vector field for L .

7 Conclusion

- For (α, β) -metric deformations L of Lorentzian metrics, we have determined the conditions such that they admit a *smooth distribution of convex cones* \mathcal{T}_x^L , with null boundary.
- We proved that there exist (curved) (α, β) -metric spacetimes admitting symmetries that are not necessarily symmetries of the underlying Lorentzian metric.

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Thank you!