Searching for new physics in B semileptonic decays using lattice QCD

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The motivation

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 $\bar{B} \rightarrow D^* \ell \bar{\nu}$ in LQCD

December 19th, 2024 2/

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- The Standard Model (SM) describes phenomena in a wide range of scales
- Yet, we expect it to fail at some point
 - Hierarchy problem, too many parameters, absence of gravity, dark matter/energy, neutrino mixing...
 - SM regarded as an Effective Field Theory (EFT)
- New physics searches more important than ever



Motivation: Searches for new physics

Energy frontier



Cosmology frontier





- High expectations with the LHC
- Intensity frontier becoming increasingly important

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Motivation: New physics in the flavor sector of the SM

$$\left(\begin{array}{ccc} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & \boldsymbol{V_{cb}} \\ V_{td} & V_{ts} & V_{tb} \end{array}\right)$$

The CKM matrix

- Matrix must be unitary (preserve the norm)
- Tensions have been there for a long time
- Evolution of the tensions according to PDG



Motivation: CKM matrix elements



Phys. Rev. Lett. 114, 011802 (2015)

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Motivation: Tensions in LFU ratios

$$R\left(D^{(*)}\right) = \frac{\mathcal{B}\left(B \to D^{(*)}\tau\nu_{\tau}\right)}{\mathcal{B}\left(B \to D^{(*)}\ell\nu_{\ell}\right)}$$



• Current $\approx 3.3\sigma$ combined tension with the SM (HFLAV)

• Tension in $R(D) \approx 1.6\sigma$ Tension in $R(D^*) \approx 2.5\sigma$

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The theory

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December 19th, 2024 8

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Semileptonic B decays on the lattice: Exclusive $|V_{cb}|$



• Decay described in the SM by 4 form factors: $g, f, \mathcal{F}_1, \mathcal{F}_2 \implies \mathcal{F}$

$$\underbrace{\frac{d\Gamma}{dw}\left(\bar{B}\to D^*\ell\bar{\nu}_\ell\right)}_{\text{Experiment}} = \underbrace{K_1(w,m_\ell\approx 0)}_{\text{Known factors}} \underbrace{\left|\mathcal{F}(w)\right|^2}_{\text{Theory}} \left|V_{cb}\right|^2, \quad w = v_{D^*} \cdot v_B$$

 $\bullet\,$ The amplitude ${\cal F}$ must be calculated in QCD, which is non-perturbative

- HQET provides $\mathcal{F}(1)$, but $K_1(w, m_\ell \approx 0) \propto (w^2 1)^{\frac{1}{2}} \rightarrow$ no data Q w = 1
- LQCD can calculate the form factors away from $\boldsymbol{w}=\boldsymbol{1}$

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- The amplitudes $\mathcal{F}_1\mathcal{F}_2$ must be calculated in the theory
- Since $K_2(w,0) = 0$, \mathcal{F}_2 only contributes significantly with the au
- Knowing these amplitudes, one can extract $\left|V_{cb}
 ight|$ from experiment
 - It is possible to extract R(D*) without experimental data!

$$R(D^{*}) = \frac{\int_{1}^{w_{\text{Max},\tau}} dw \left[K_{1}(w,m_{\tau}) \left| \mathcal{F}(w) \right|^{2} + K_{2}(w,m_{\tau}) \left| \mathcal{F}_{2}(w) \right|^{2} \right] \times \mathcal{V}_{cons}^{2}}{\int_{1}^{w_{\text{Max}}} dw \left[K_{1}(w,0) \left| \mathcal{F}(w) \right|^{2} \right] \times \mathcal{V}_{cons}^{2}}$$

• $|V_{cb}|$ cancels out

10 / 20

Semileptonic B decays on the lattice: Introduction to Lattice QCD

$$\mathcal{L}_{QCD} = \sum_{f} \bar{\psi}_{f} \left(\gamma^{\mu} D_{\mu} + m_{f} \right) \psi_{f} + \frac{1}{4} \text{tr} F_{\mu\nu} F^{\mu\nu}$$



- Discretize space-time in a computer
 - Finite lattice spacing a
 - Finite spatial volume L

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Finite time extent T

• Perform simulations in an unphysical setup and approach the physical limit

- Enlarge the volume and reduce a
- Quark masses \implies Pion masses (hadrons are matched)
- Number of sea quarks $n_f = 2 + 1, 2 + 1 + 1, 1 + 1 + 1 + 1 \dots$

Semileptonic B decays on the lattice: Introduction to Lattice QCD

The systematic error analysis is based on **EFT** descriptions of QCD The EFT description:

- provides functional form for different extrapolations (or interpolations)
- can be used to construct improved actions
- can estimate the size of the systematic errors



In order to keep the systematic errors under control we must repeat the calculation for several lattice spacings, volumes, light quark masses... and use the EFT to extrapolate to the physical theory

12/20

- Heavy quark treatment in Lattice QCD
 - For heavy quarks $(m_Q > \Lambda_{QCD})$, discretization errors typically grow as $\sim a^2 m_Q^2$
 - Two treatments
 - EFTs: Physical heavy masses, but requires matching and renormalization

Semileptonic B decays on the lattice: Heavy quarks

- Same action as the light quarks: Unphysical heavy masses, requires extrapolation
- Not all actions perform equally well
 - Typical action $\sim a^2 m_Q^2$ VS HISQ action $\sim \alpha_s a^2 m_Q^2$





Image: A matrix

The mess

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December 19th, 2024 14

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The mess: Lattice results



The mess: Lattice results





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The mess: Lattice results





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The mess: Combined lattice fits



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Image: A math a math

The mess: Experimental data

Total Events = $\int dt$ Luminosity \times Cross Section



• LHCb integrated luminosity (IL) \approx 7 fb⁻¹

- Target 300-350 fb⁻¹
- Cross-section $\approx 10^5$ larger than in Belle, but spread around many particles

• Belle II IL \approx 360 ${\rm fb}^{-1}$

Image: A matching of the second se

- Target 50 $ab^{-1} = 50000 \text{ fb}^{-1}$
- Belle I 711 fb⁻¹ (×70!!)

Conclusions

- $\bullet\,$ Great progress in LQCD calculations of $B\to D^*\ell\nu$ form factors
- Good agreement between different LQCD results
 - Not so good between LQCD and experiment
- New calculations are needed to clarify the situation

Thank you for your attention

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BACKUP SLIDES

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Semileptonic B decays on the lattice: Parametrizations

Most parametrizations perform an expansion in the z parameter



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22 / 20

Semileptonic B decays on the lattice: Parametrizations

• Boyd-Grinstein-Lebed (BGL)

 $f_X(w) = \frac{1}{B_{f_X}(z)\phi_{f_X}(z)} \sum_{n=1}^{\infty} a_n z^n \frac{\frac{P_{\text{hys. Rev. Lett. 74 (1995) 4603-4606}}{N_{\text{ucl. Phys. B461 (1996) 493-511}}}{\frac{N_{\text{ucl. Phys. B461 (1996) 493-511}}{N_{\text{ucl. Phys. B461 (1996) 493-511}}}$

- B_{f_X} Blaschke factors, includes contributions from the poles
- ϕ_{f_X} is called outer function and must be computed for each form factor
- Weak unitarity constraints $\sum_n |a_n|^2 \leq 1$
- Caprini-Lellouch-Neubert (CLN)

Nucl. Phys. B530 (1998) 153-181

 $F(w) \propto 1 - \rho^2 z + c z^2 - d z^3$, with $c = f_c(\rho), d = f_d(\rho)$

- Relies strongly on HQET, spin symmetry and (old) inputs
- Tightly constrains F(w): four independent parameters, one relevant at w = 1
- Current consensus: abandon CLN
 - Spiritual sucessors of CLN Bernlochner et al. Phys.Rev.D 95 (2017) 115008, Phys.Rev.D 97 (2018) 059902

Bordone, Gubernari, Jung, Straub, Van Dyk... Eur.Phys.J.C 80 (2020) 74, Eur.Phys.J.C 80 (2020) 347, JHEP 01 (2019) 009

23 / 20

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Semileptonic B decays on the lattice: Chiral-continuum extrapolation



$B \to D^* \ell \nu$: Chiral-continuum extrapolation



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December 19th, 2024 25 / 20

The mess: Combined lattice fits



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December 19th, 2024

26 / 20

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27 / 20

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Future projects: HISQ + Fermilab

- Fermilab/MILC calculation
- Using 7 $N_f = 2 + 1 + 1$ ensembles of sea HISQ quarks
- The heavy quarks use the Fermilab effective action
 - $\bullet\,$ Correlated with a $B\to L\ell\nu$ analysis using the same data
 - Four channels $B_{(s)} \to D_{(s)}^{(*)} \ell \nu$ in a single correlated analysis



Future projects: HISQ²

- Fermilab/MILC calculation
- Planning to use 9 $N_f = 2 + 1 + 1$ ensembles of sea HISQ quarks
- The heavy quarks use the HISQ action
 - Physical bottom mass reachable with the finest ensembles
- m_{π} physical in several ensembles



29 / 20