

Novel constraints on the primordial power spectrum from compact objects formation

IPPP/25/40

Updated constraints on the primordial power spectrum at sub-Mpc scales

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The primordial power spectrum of matter density perturbations contains highly valuable information about new fundamental physics, in particular cosmological inflation, but is only very weakly constrained observationally for small cosmological scales $k \gtrsim 3 \text{ Mpc}^{-1}$. We derive novel constraints, $\mathcal{P}_{\mathcal{R}}(k) \lesssim 5 \cdot 10^{-6}$ over a large range of such scales, from the formation of ultracompact minihalos in the early universe. Unlike most existing constraints of this type, our results do not rest on the assumption that dark matter can annihilate into ordinary matter.

I. INTRODUCTION

The power spectrum of primordial density fluctuations, $\mathcal{P}(k)$, is a key quantity in cosmology. It describes the initial conditions of the Universe right after the big bang, tiny perturbations in an almost homoge-

background [29, 30] provide complementary constraints directly on the linear, or only mildly non-linear, power spectrum at scales down to roughly one pc.

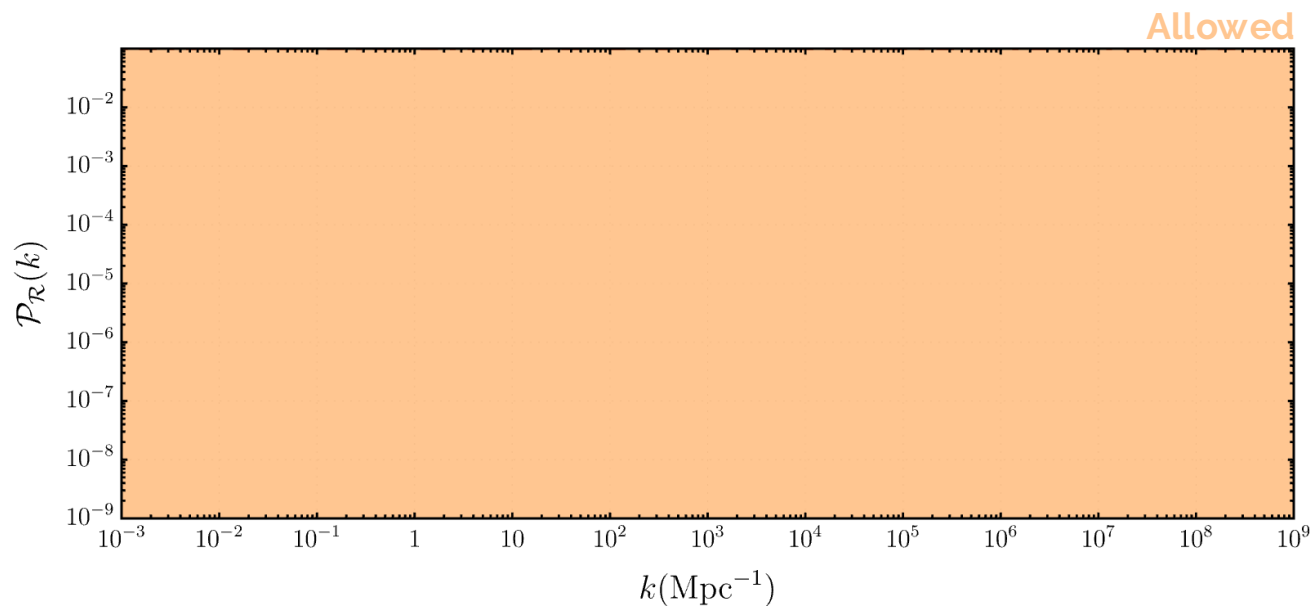
UCMHs are only one example of extended DM objects (EDOs) that may have existed since primordial times. Recently, in fact, the possibility of more general

Based on arXiv: 2506.20704

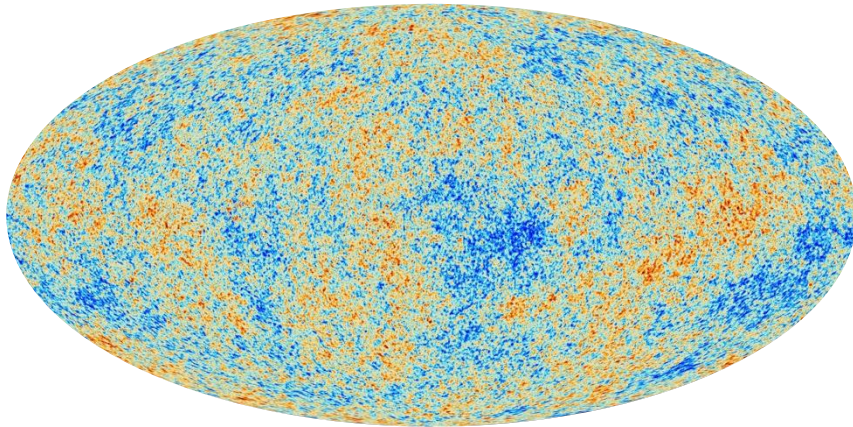
Sergio Sevillaño Muñoz

Our plan for today

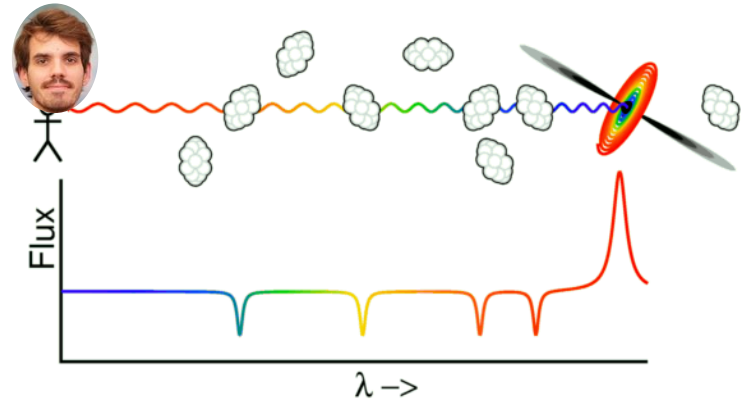
Update constraints from the primordial curvature power spectrum



Cosmic Microwave Background



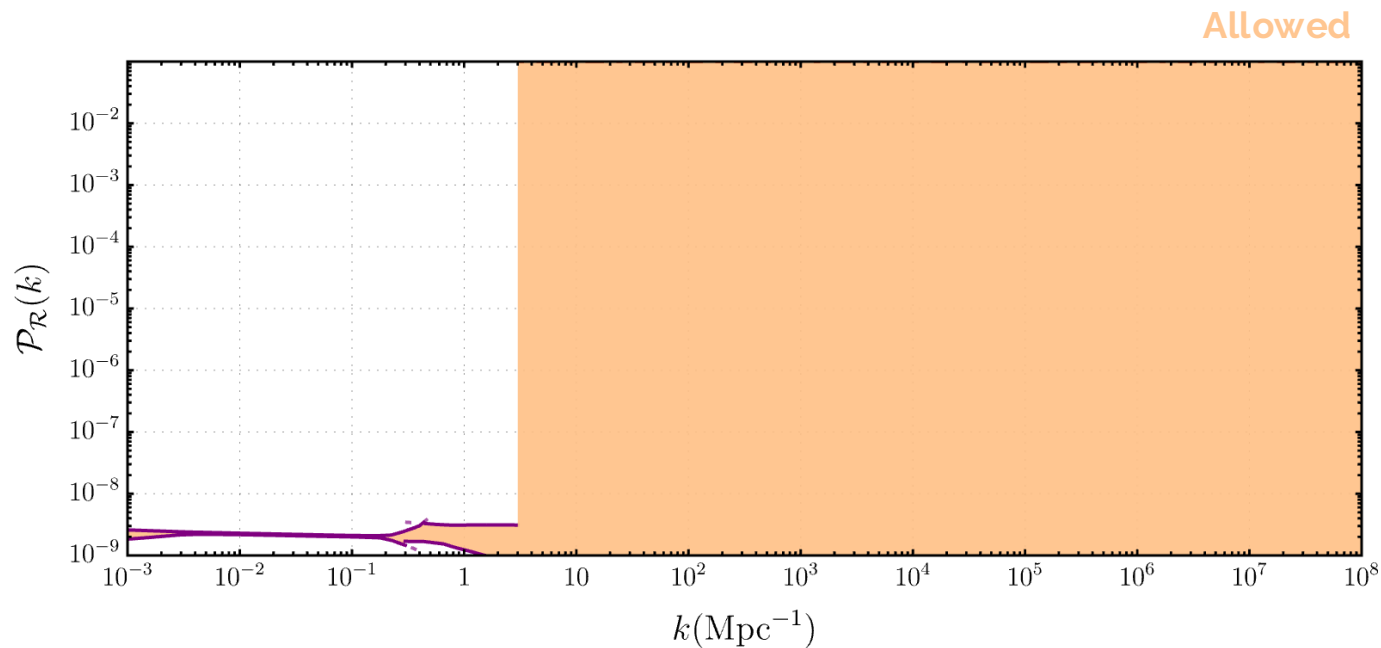
Lyman α Forest



S. Bird, H. V. Peiris, M. Viel, L. Verde (1010.1519v2)

Lyman α and Planck

Inferring smaller scales by looking at matter distribution in the late universe



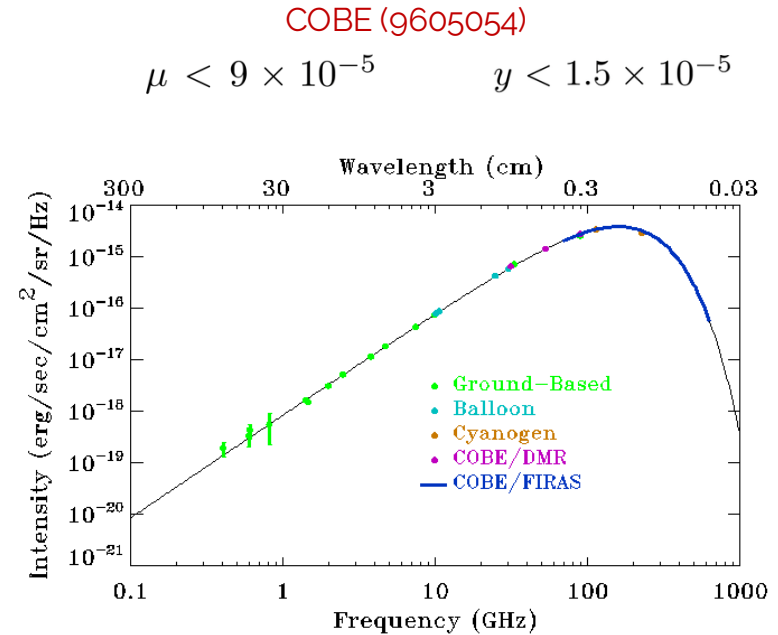
y and μ distortions

Primordial perturbations create injections of energy into the background that lead to a departure of the perfect CMB black body spectrum:

$$\mu \approx 2.2 \int_{k_{\min}}^{\infty} \mathcal{P}_{\mathcal{R}}(k) \left[\exp\left(-\frac{\hat{k}}{5400}\right) - \exp\left(-\left[\frac{\hat{k}}{31.6}\right]^2\right) \right] d \ln k$$
$$y \approx 0.4 \int_{k_{\min}}^{\infty} \mathcal{P}_{\mathcal{R}}(k) \exp\left(-\left[\frac{\hat{k}}{31.6}\right]^2\right) d \ln k,$$

Where we assume
a locally scale-invariant power spectrum

$$\mathcal{P}_{\mathcal{R}}(k) = \begin{cases} \mathcal{P}_{\mathcal{R}}(k_{\mathcal{R}}) & \text{for } 1/3 < k/k_{\mathcal{R}} < 3 \\ 0 & \text{otherwise} \end{cases}$$

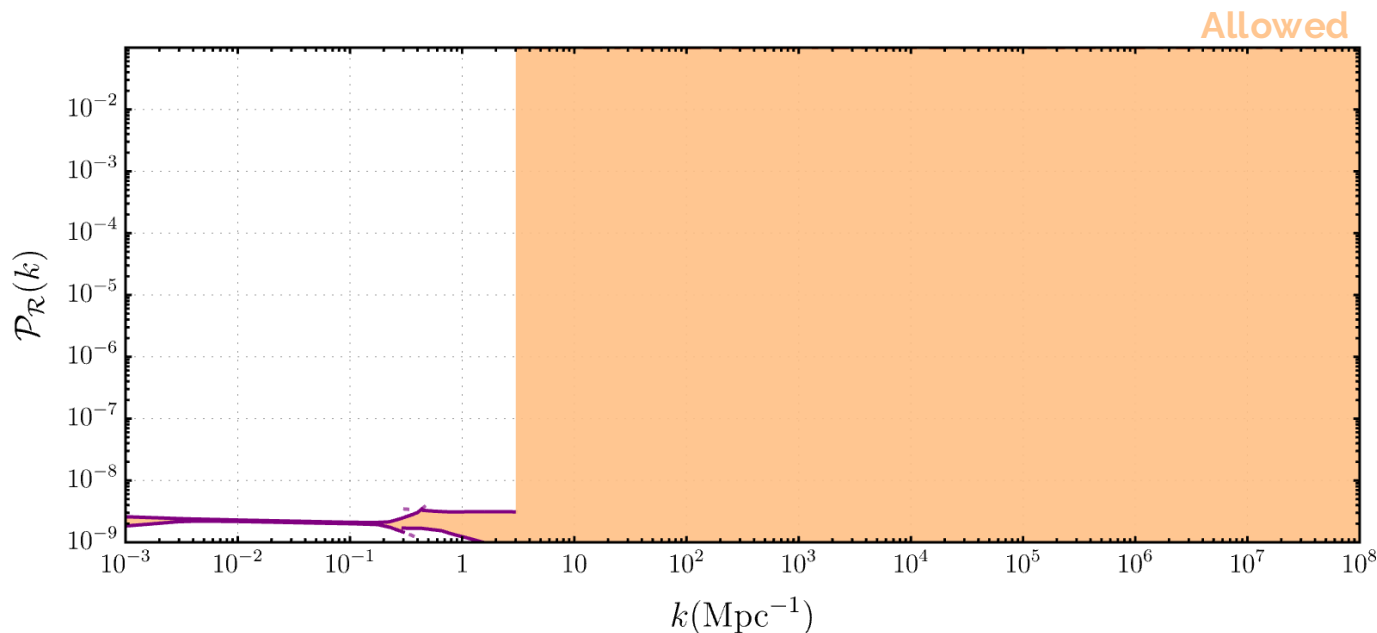


J. Chluba, A. L. Erickcek, I. Ben-Dayan (1203.2681)

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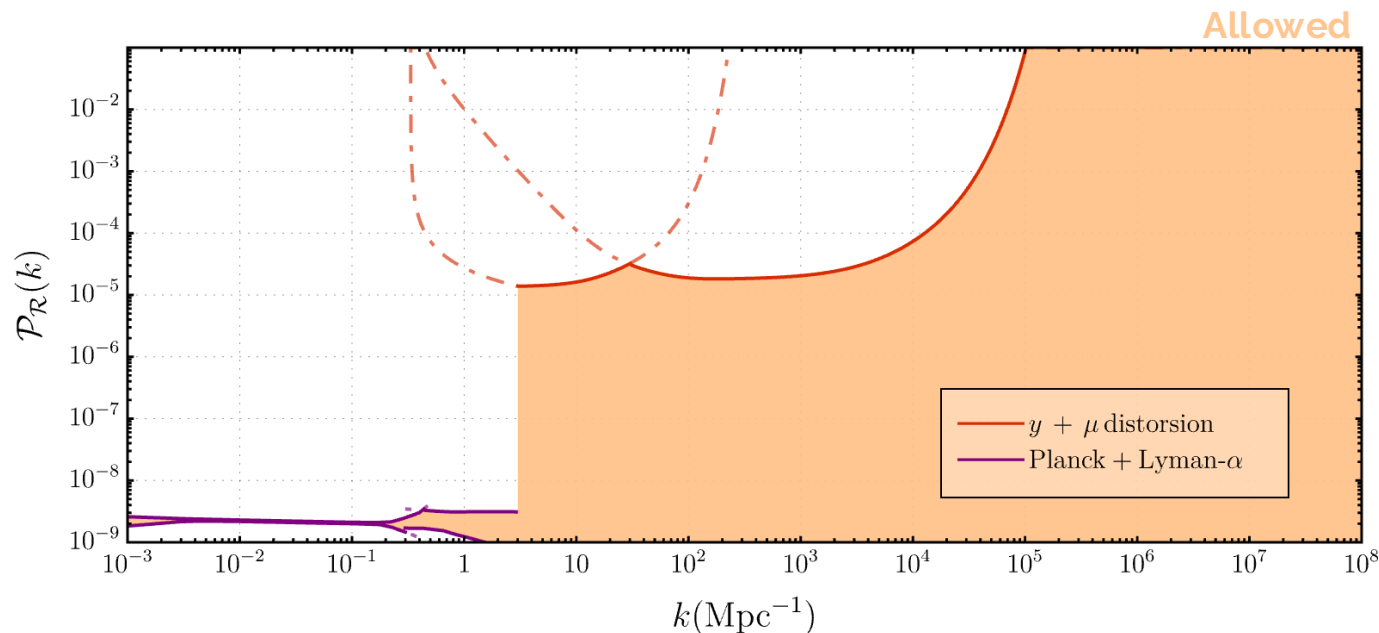
COBE (9605054) J. Chluba, A. L. Erickcek, I. Ben-Dayan (1203.2681)



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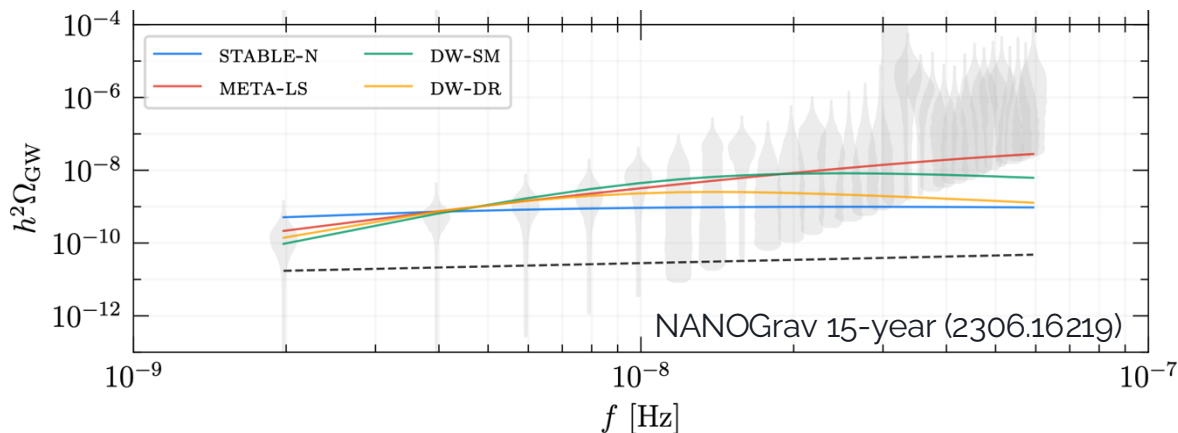


Pulsar Timing Arrays (PTAs)

Any stochastic background of gravitational waves can be constrained via $h^2\Omega_{\text{GW}}$

$$\Omega_{\text{GW}}(k, \eta) = \frac{1}{24} \left(\frac{k}{aH} \right)^2 \mathcal{P}_h(k, \eta) \quad \mathcal{P}_{\mathcal{R}}(k) = \begin{cases} \mathcal{P}_{\mathcal{R}}(k_{\mathcal{R}}) & \text{for } 1/3 < k/k_{\mathcal{R}} < 3 \\ 0 & \text{otherwise} \end{cases}$$

$$\mathcal{P}_h(k, \eta) = 4 \int_0^\infty dv \int_{|1-v|}^{1+v} du \left(\frac{4v^2 - (1 + v^2 - u^2)^2}{4vu} \right)^2 \mathcal{P}_{\mathcal{R}}(ku) \mathcal{P}_{\mathcal{R}}(kv) I^2(u, v, k\eta),$$

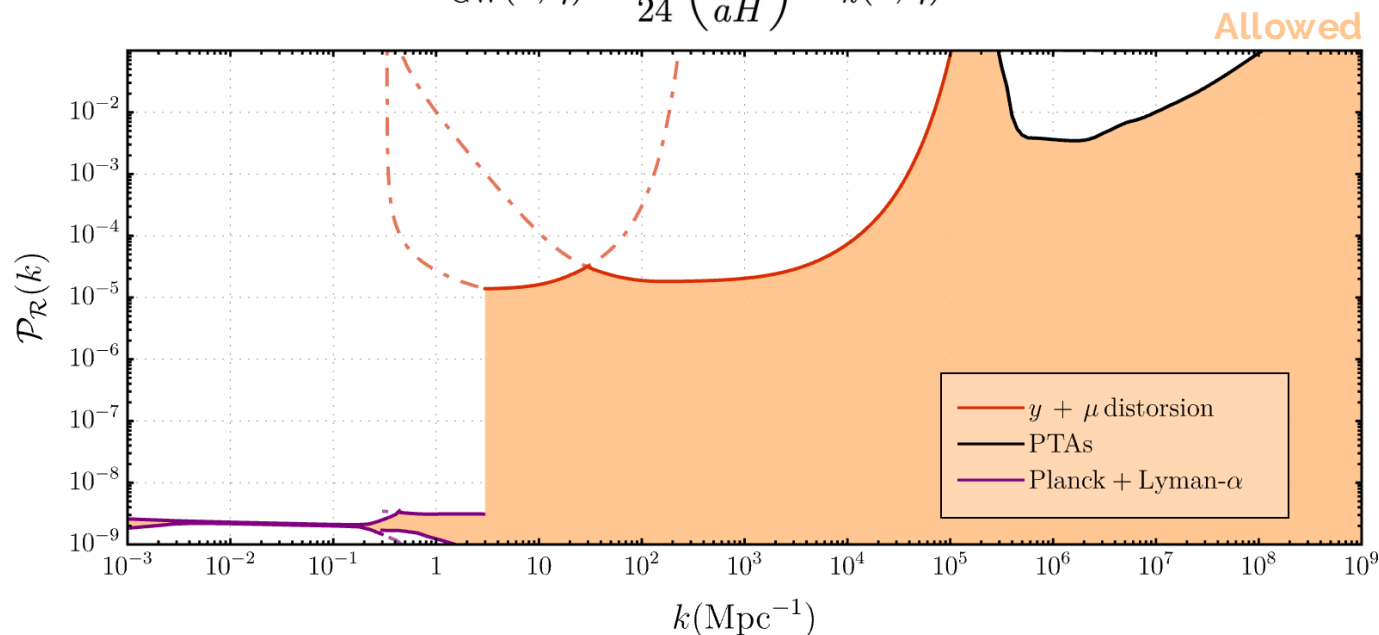


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NANOGrav 15-year (2306.16219)

$$\Omega_{\text{GW}}(k, \eta) = \frac{1}{24} \left(\frac{k}{aH} \right)^2 \mathcal{P}_h(k, \eta)$$



Primordial Black Holes (PBHs)

Large enough overdensities can lead to the collapse of a region into a black hole!

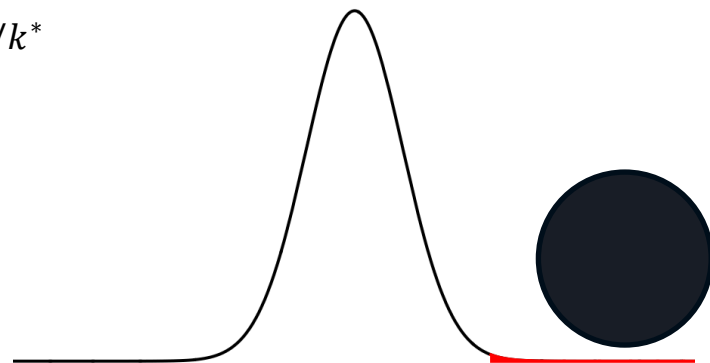
$$\beta(M_{\text{H}}) = \sqrt{\frac{2}{\pi}} \frac{\alpha}{\sigma(R)} \int_{\delta_c}^{\infty} \exp\left(-\frac{\delta^2(R)}{2\sigma^2(R)}\right) d\delta(R)$$

$R = 1/k^*$

$$\delta \equiv (\rho - \bar{\rho})/\bar{\rho}$$

$$\sigma^2(R) = \frac{16}{81} \int_0^{\infty} (kR)^4 W^2(kR) \mathcal{P}_{\mathcal{R}}(k) T^2(kR/\sqrt{3}) \frac{dk}{k}$$

Fraction of black holes

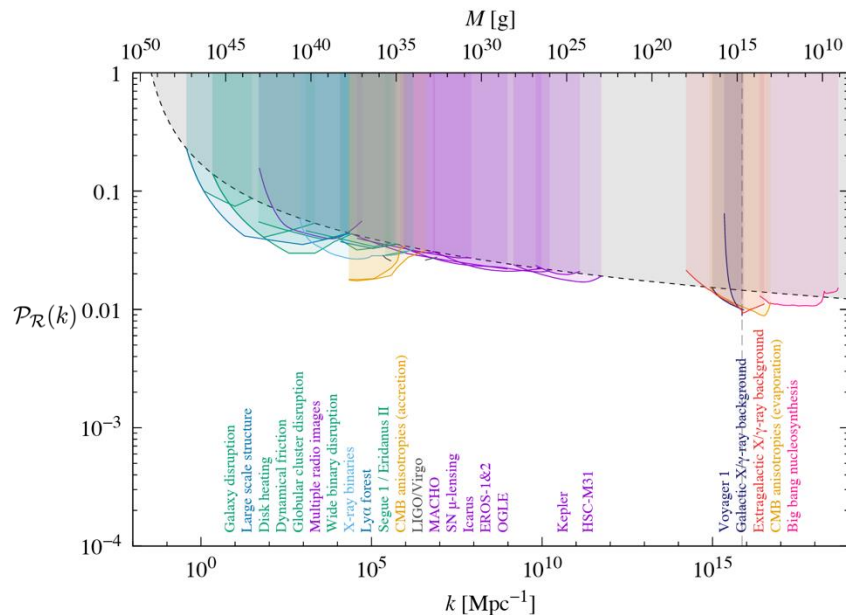
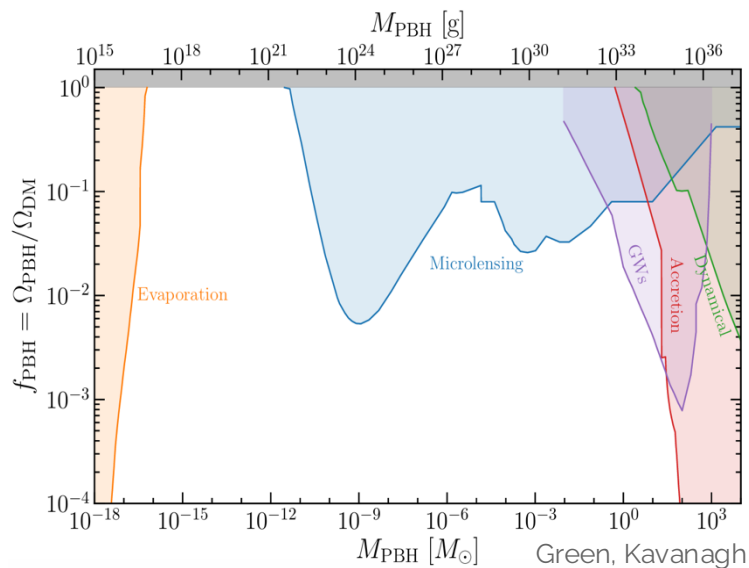


We can therefore constrain the power spectrum using limits on PBHs!

Primordial Black Holes (PBHs)

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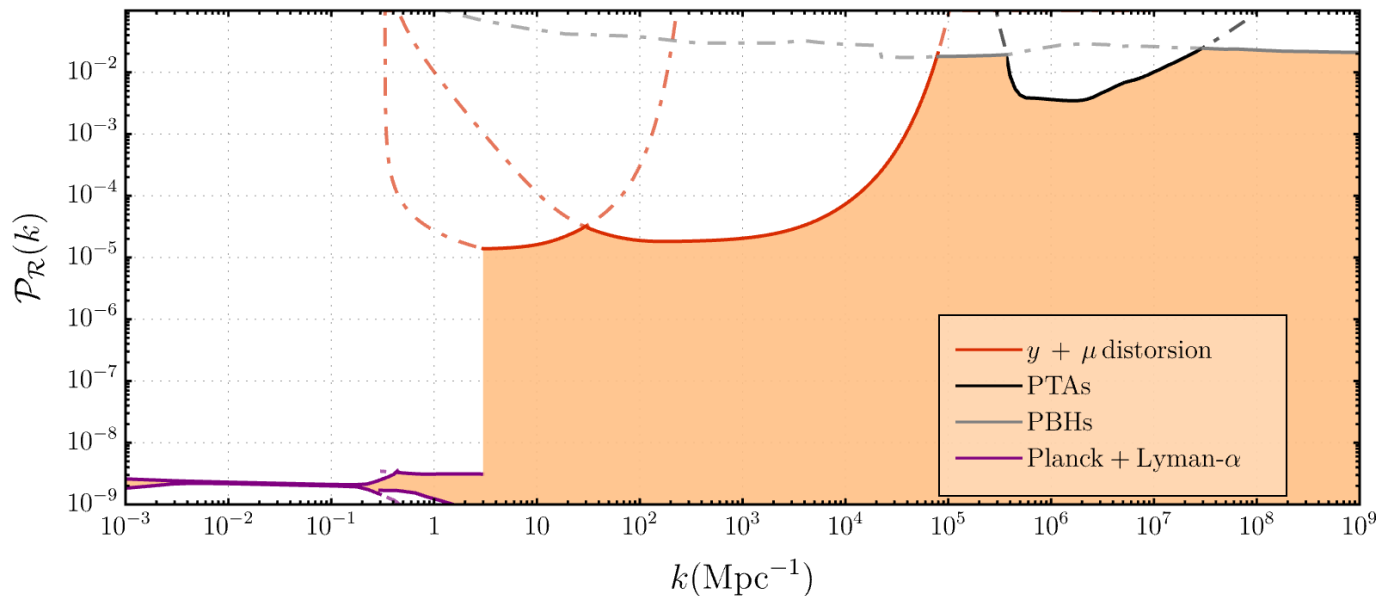
Carr, Kohri, Sndouda and Yokoyama (2002.12778)



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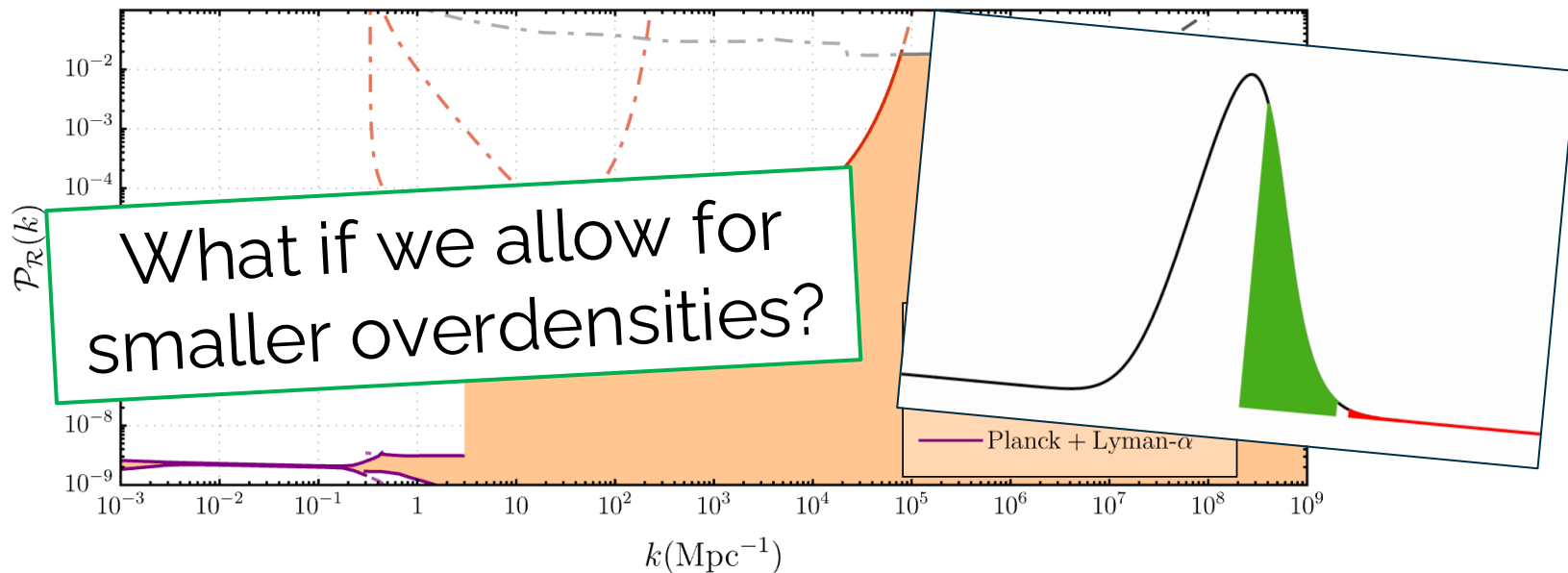
$$\beta(M_{\text{H}}) = \sqrt{\frac{2}{\pi}} \frac{\alpha}{\sigma(R)} \int_{\delta_c}^{\infty} \exp\left(-\frac{\delta^2(R)}{2\sigma^2(R)}\right) d\delta(R)$$



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Ultra-compact minihaloes (UCMHs)

That was the idea from this paper, assuming SM-DM interactions

**Improved constraints on the primordial power spectrum at small scales from
ultracompact minihalos**

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Luruper Chausse 149, DE-22761 Hamburg, Germany*

Pat Scott[†]

Department of Physics, McGill University, 3600 rue University, Montréal, QC, H3A 2T8, Canada

Yashar Akrami[‡]

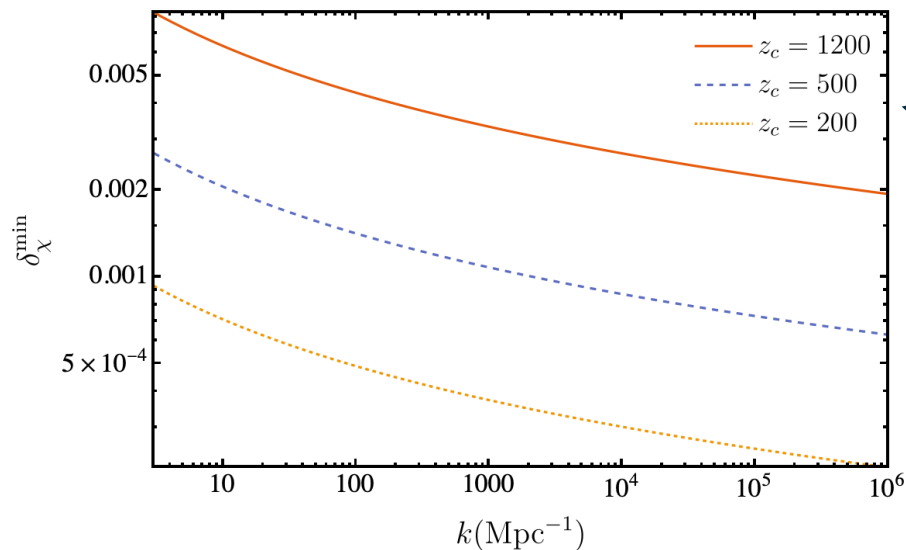
*The Oskar Klein Centre for Cosmoparticle Physics,
Department of Physics, Stockholm University,
AlbaNova, SE-106 91 Stockholm, Sweden*

Can we do this only using gravitational interactions?

Ultra-compact minihaloes (UCMHs)

To know how many of these are created we need two things:

- 1) How large does an overdensity need to be to collapse into an object?



Depends on *when* it collapses

Ultra-compact minihaloes (UCMHs)

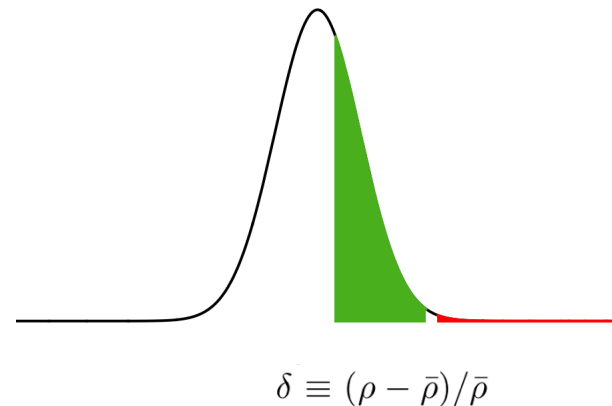
To know how many of these are created we need two things:

2) What is the distribution of overdensity amplitudes given a scale?

Similar to primordial black holes

$$\beta(k) = \frac{1}{\sqrt{2\pi}\sigma_{\chi,H}(k)} \int_{\delta_{\chi}^{\min}(k)}^{\delta_{\chi}^{\max}} \exp\left[-\frac{\delta_{\chi}^2}{2\sigma_{\chi,H}^2(k)}\right] d\delta_{\chi}$$

$$\sigma_{\chi,H}^2(k) = \frac{1}{9} \int_0^{\infty} dx x^3 W_{\text{TH}}^2(x) \mathcal{P}_{\mathcal{R}}(xk) T_{\chi}^2(\theta = x/\sqrt{3})$$



Integrating for smaller overdensities gives density fraction!

$$\Omega_{\text{UCMH}}(k) = \Omega_{\text{DM}}\beta(k)$$

Ultra-compact minihaloes (UCMHs)

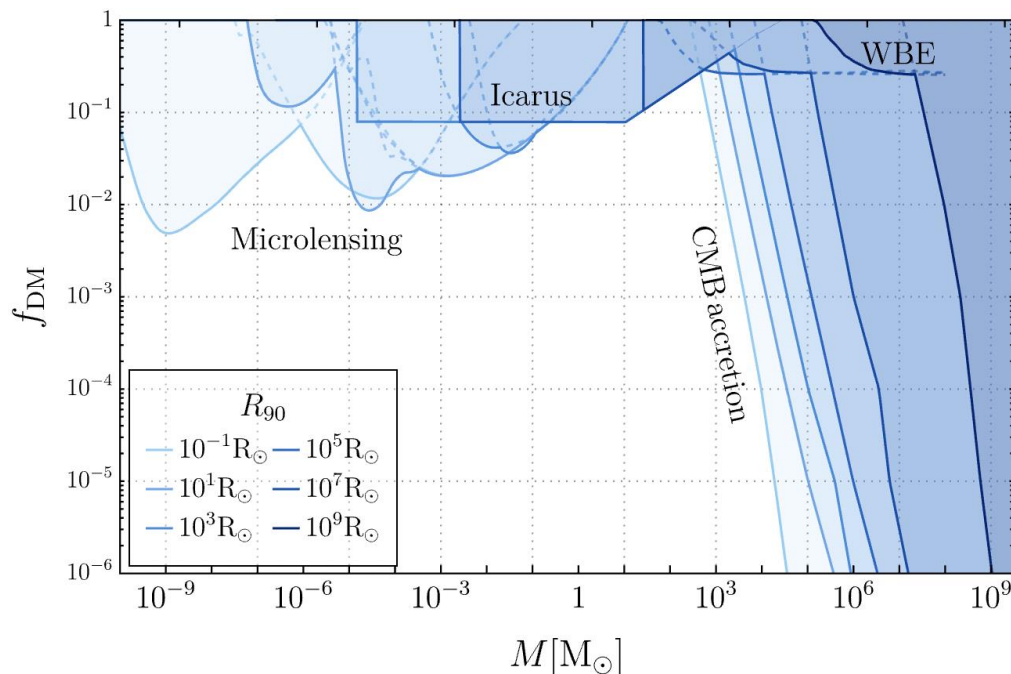
EDObounds:

Choose:

- Shape of the object
- Radius
- Bounds to plot
- Mass distributions!!

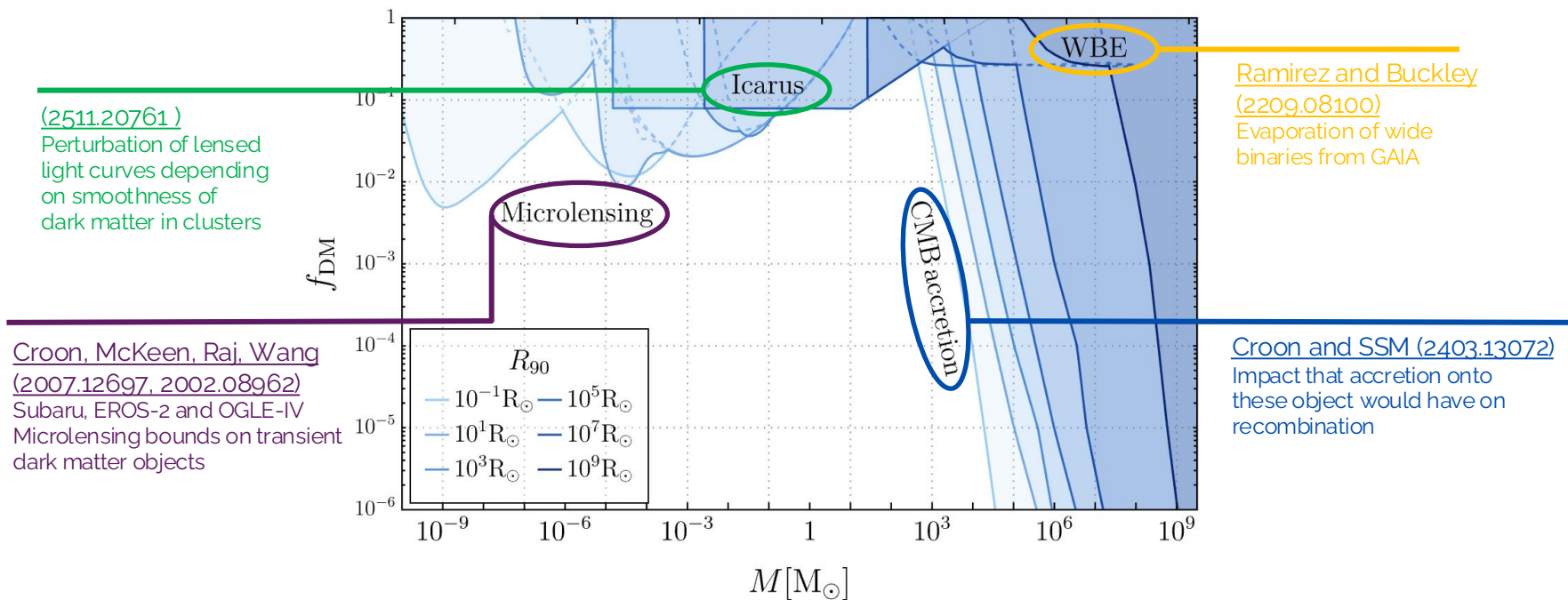
[arXiv:2407.02573](https://arxiv.org/abs/2407.02573)

There has been recent interest on constraining these type of objects:



Ultra-compact minihaloes (UCMHs)

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Ultra-compact minihaloes (UCMHs)

The formed objects have a given mass, radius and density for a given scale $k \sim 1/R$:

$$M \approx \left[\frac{4\pi}{3} \rho_\chi H^{-3} \right]_{aH=1/R} = \frac{H_0^2}{2G} \Omega_\chi R^3 \quad \rho_\chi(r, z_c) = \frac{3f_\chi M(z_c)}{16\pi R(z_c)^{3/2} r^{3/2}}$$

$$\frac{R(z_c)}{\text{pc}} = 0.019 \left(\frac{1000}{z_c + 1} \right) \left(\frac{M(z_c)}{M_\odot} \right)^{1/3}$$

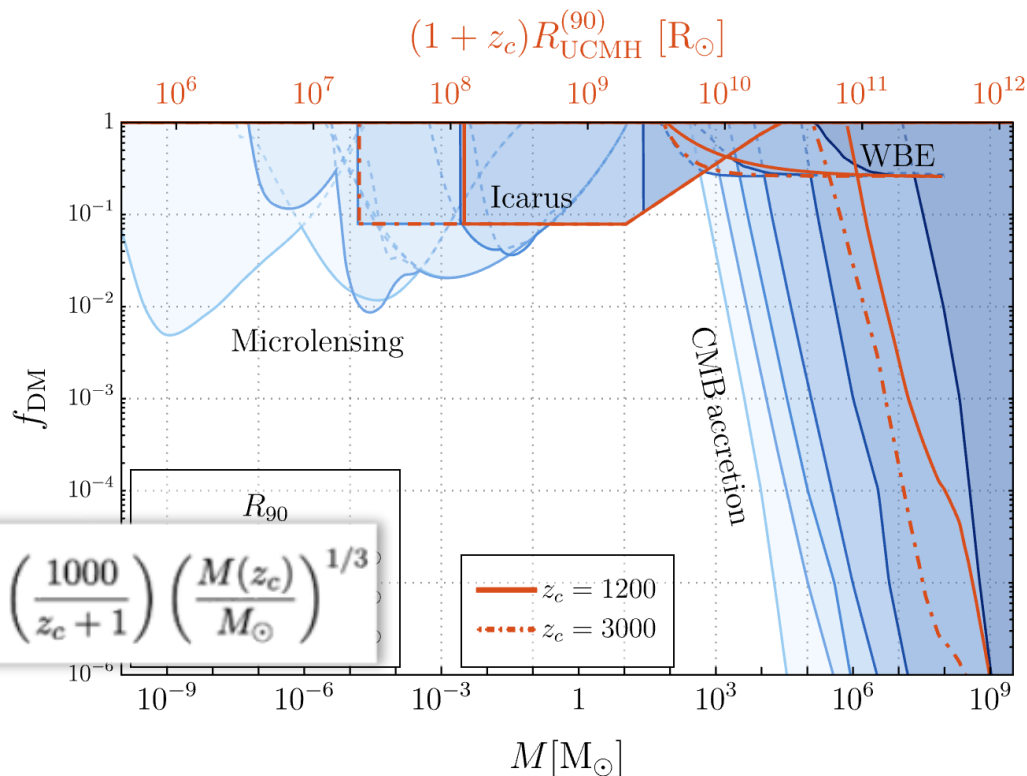
where z_c is the collapse redshift



Different density functions have been used, but this is the most conservative

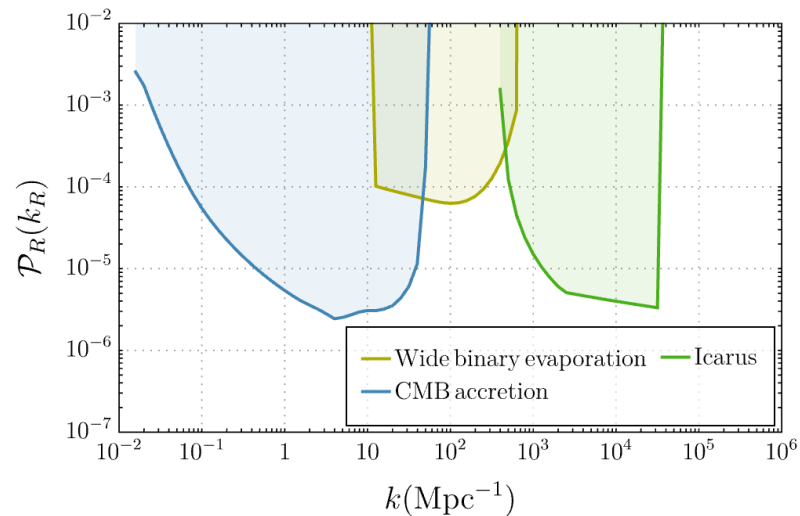
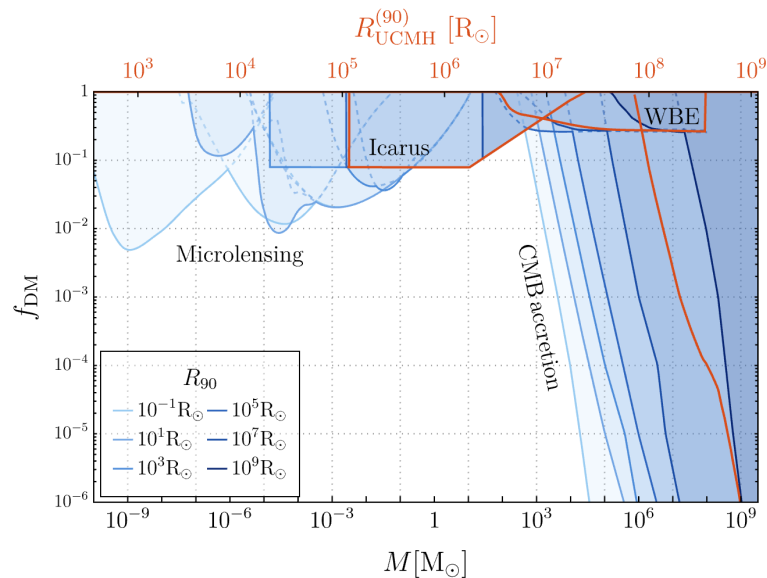
Ultra-compact minihaloes (UCMHs)

Only specific objects that have the mass-radius relation can be used to constrain the power spectrum



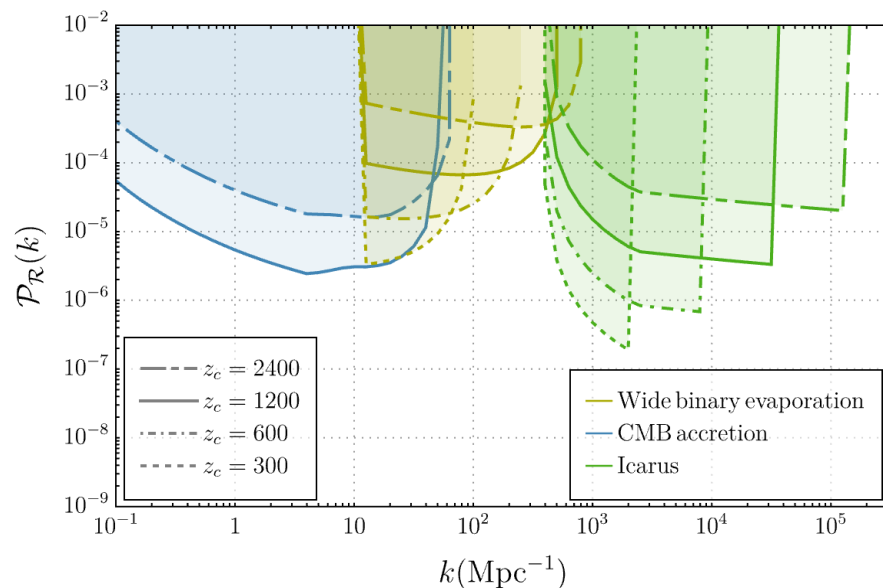
Ultra-compact minihaloes (UCMHs)

We can now translate these bounds (in red) to power spectrum using $\Omega_{\text{UCMH}}(k) = \Omega_{\text{DM}}\beta(k)$



Ultra-compact minihaloes (UCMHs)

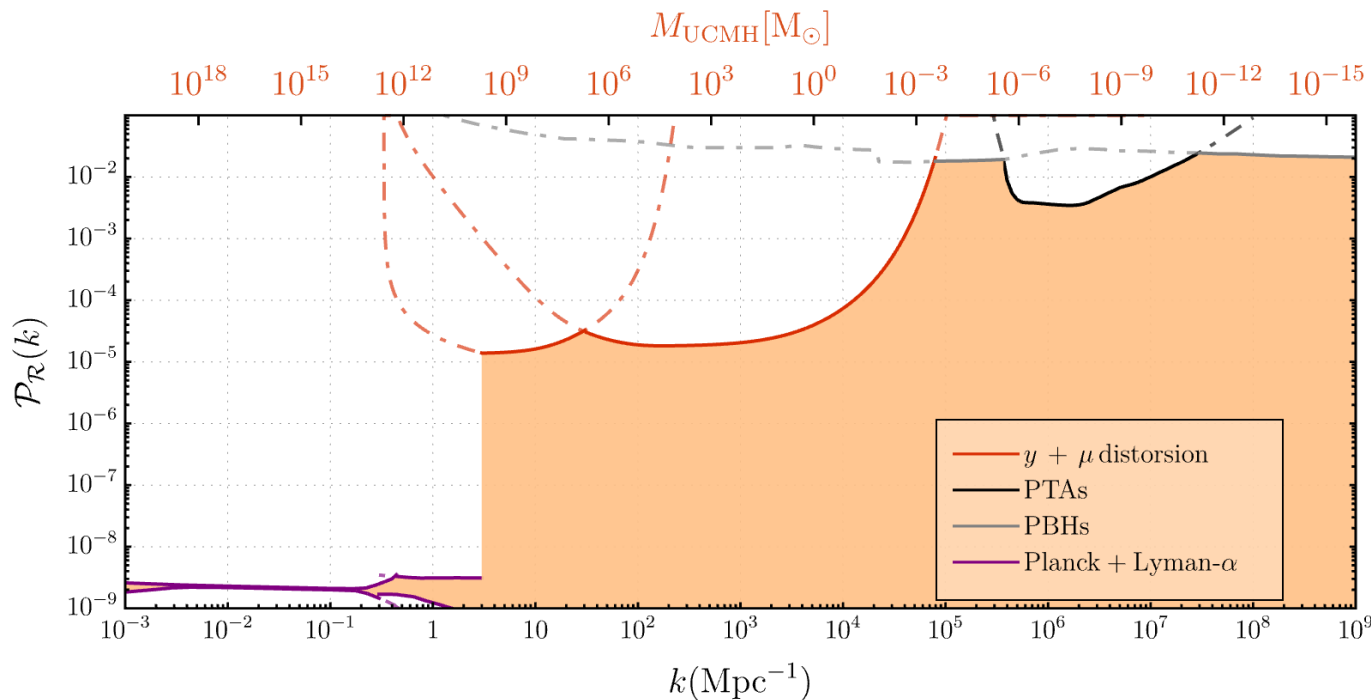
Different collapse times lead to different constraints



We will take the conservative choice $z_c > 1200$

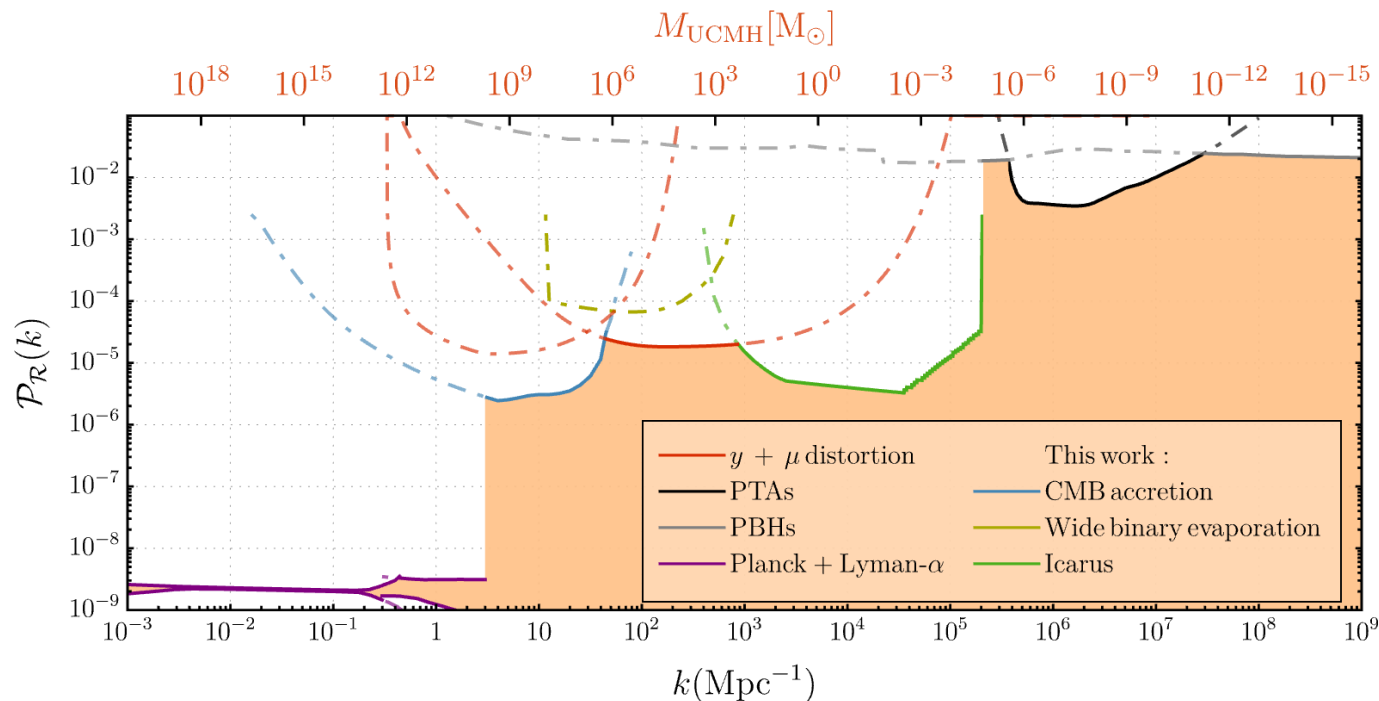
Ultra-compact minihaloes (UCMHs)

We can now translate these bounds to power spectrum using $\Omega_{\text{UCMH}}(k) = \Omega_{\text{DM}}\beta(k)$

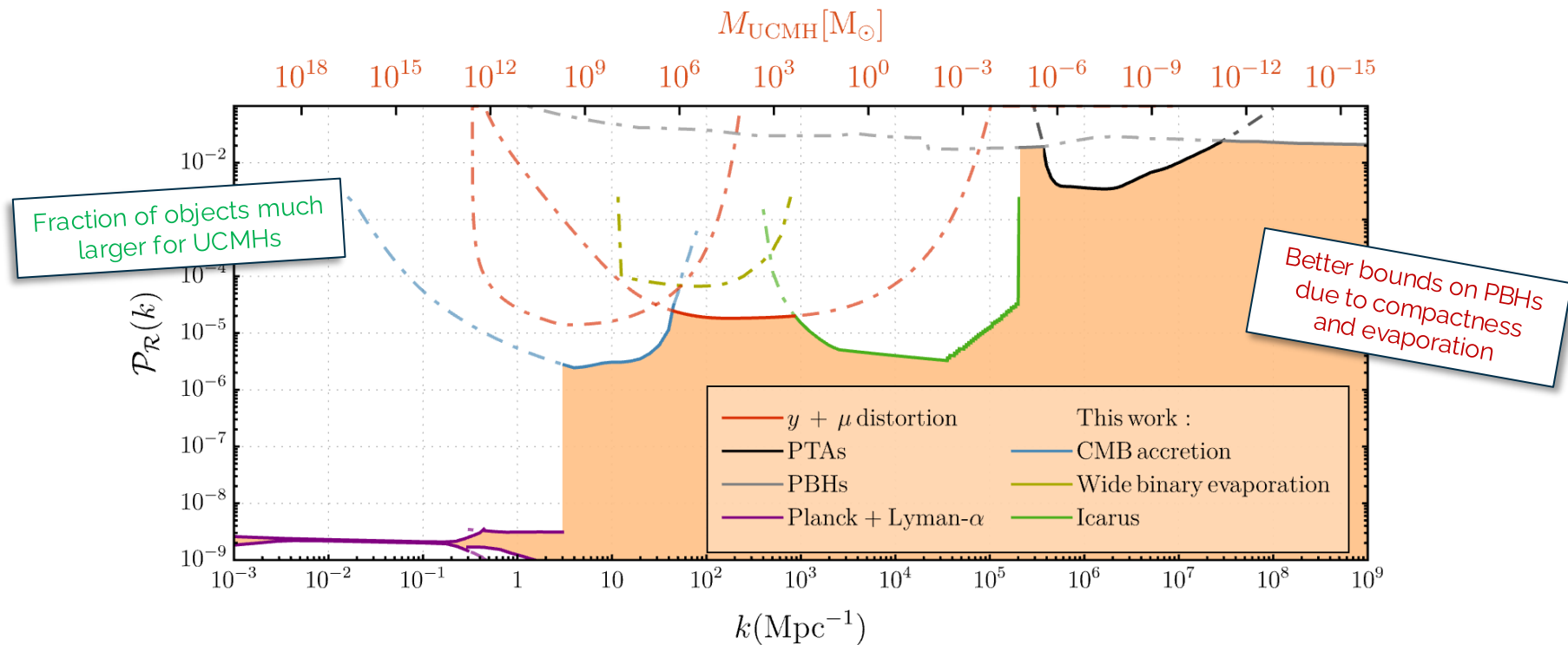


Ultra-compact minihaloes (UCMHs)

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PBHs vs UCMH



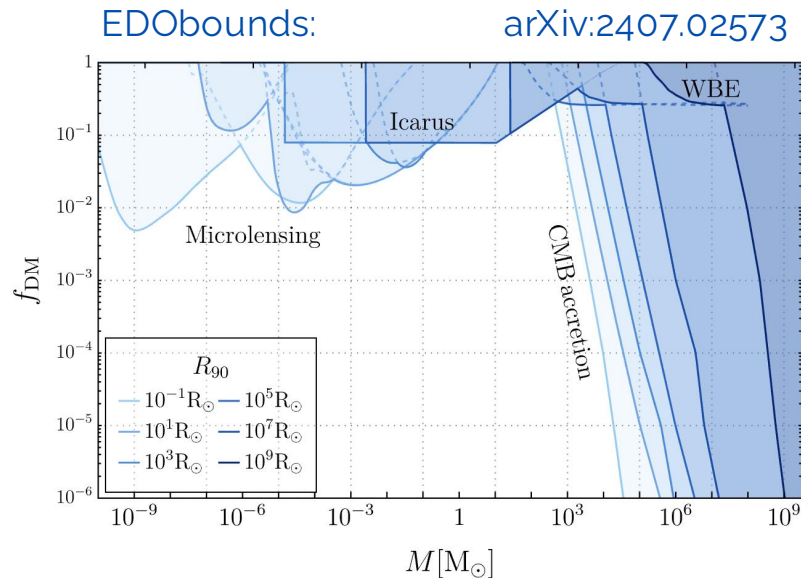
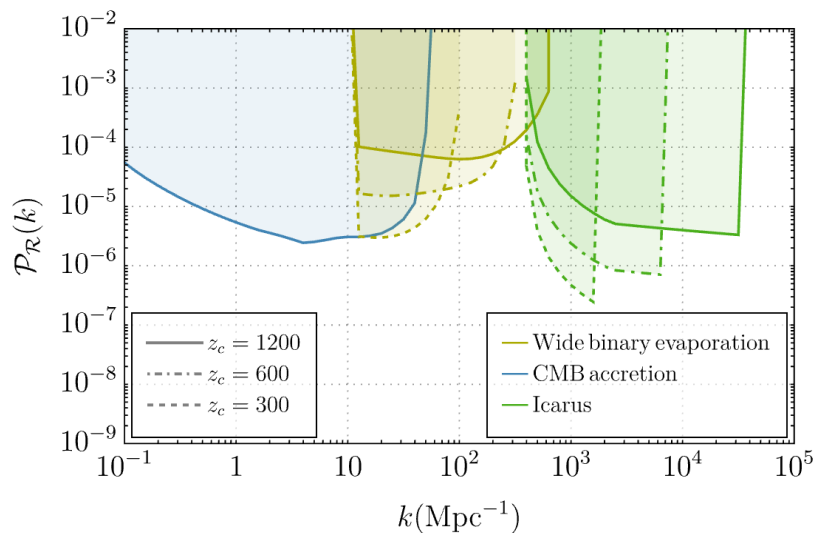
No <100% of PBHs with mass $> 100 M_{\odot}$

Soon no masses $> 10^{-1} M_{\odot}$

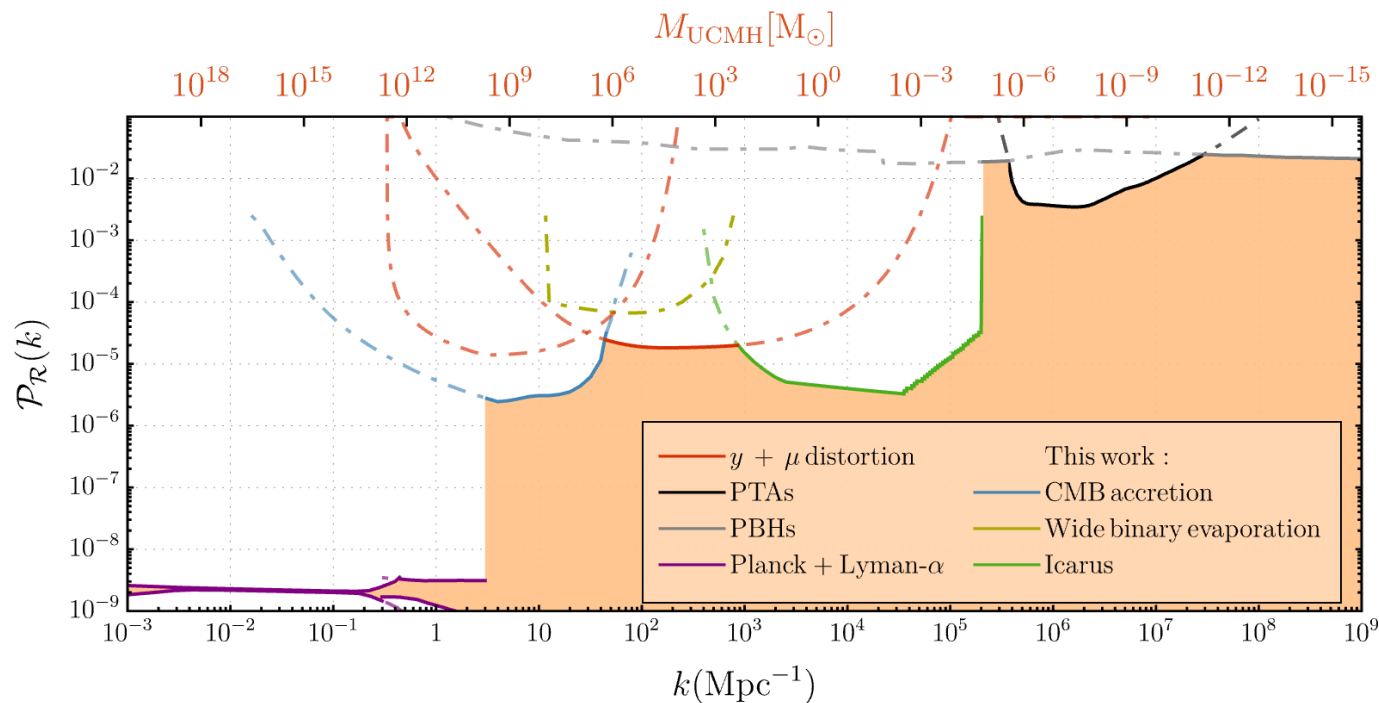
Conclusion

We can obtain improved constraints on the Power Spectrum using UCMHs!

There is room for improvement :)



Thank you for your attention!

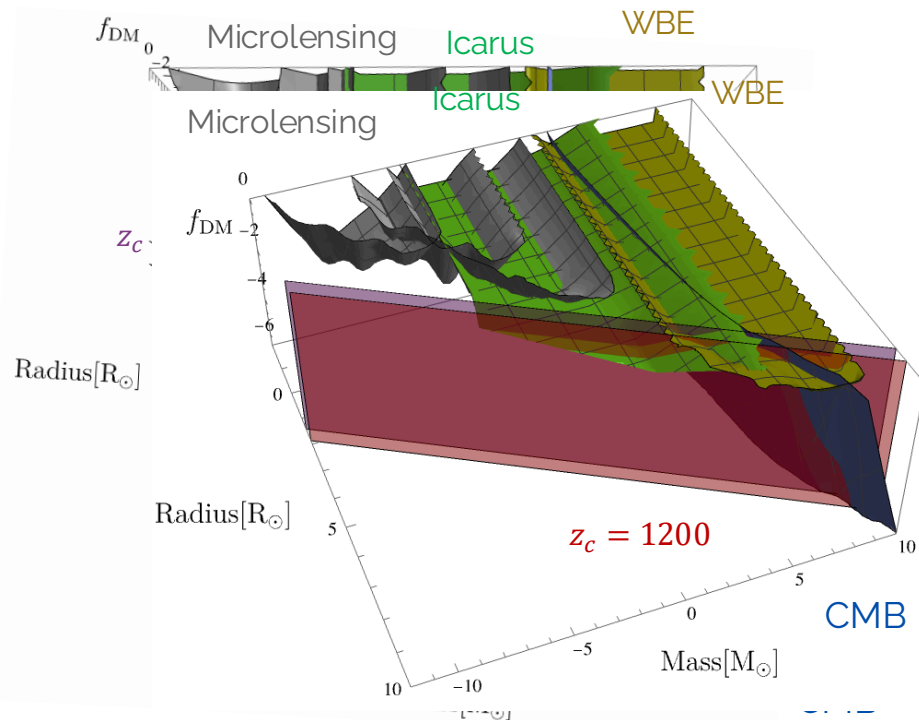
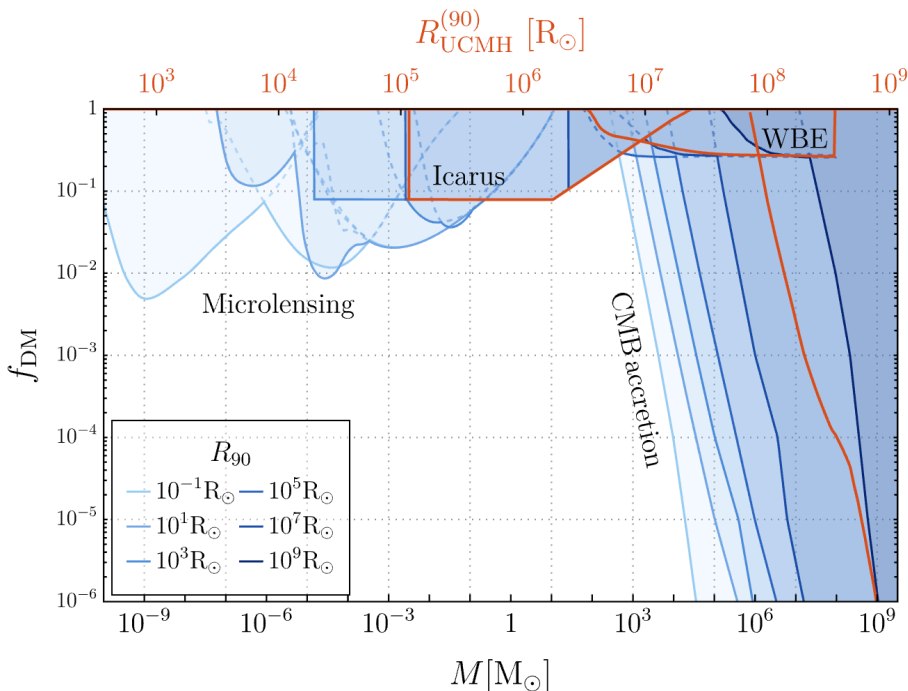


Based on arXiv: 2506.20704

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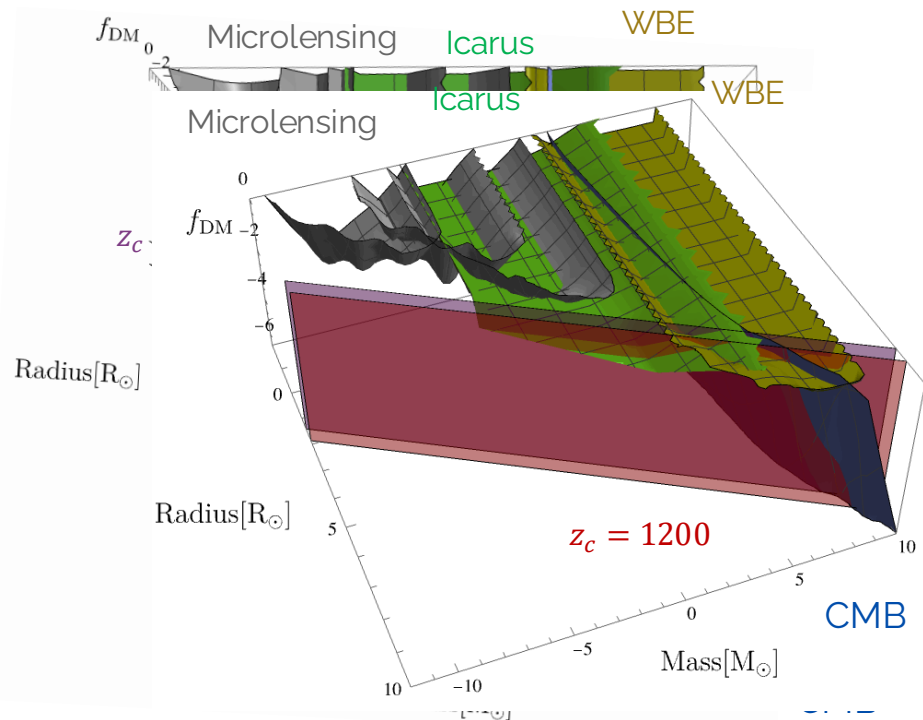
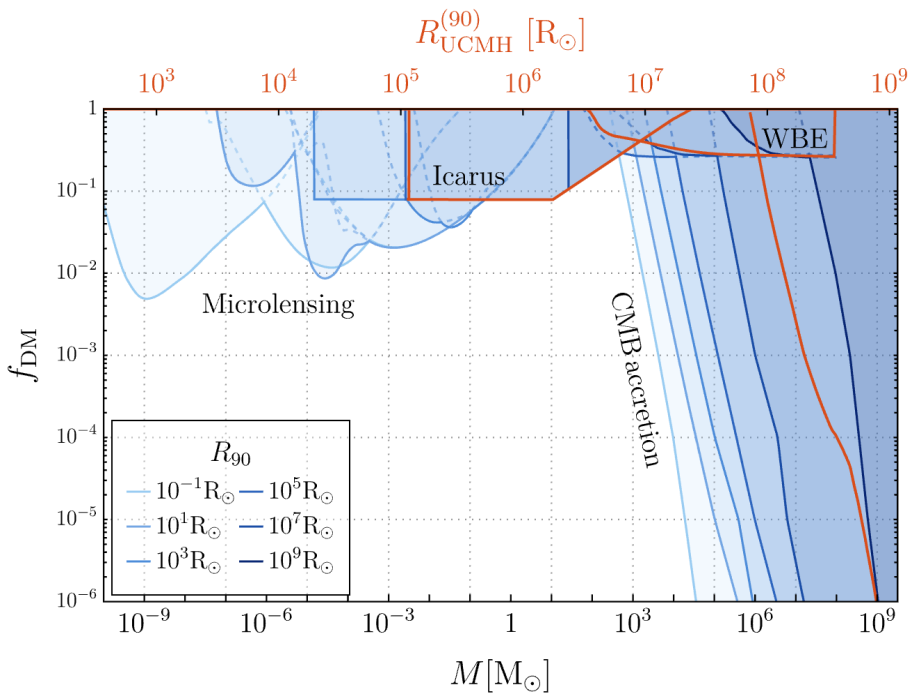
Only specific objects that have the mass-radius relation can be used to constrain the power spectrum



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