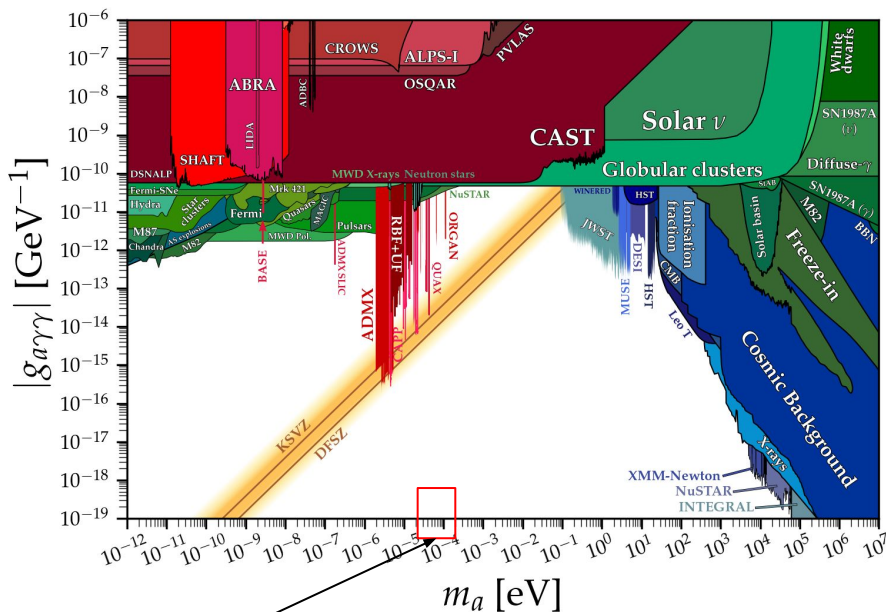


# **An efficient method of tuning plasma haloscopes with spirals and concentric rings**

Jianyang Qi  
Stockholm University for IDM 2026  
June 4th, 2026

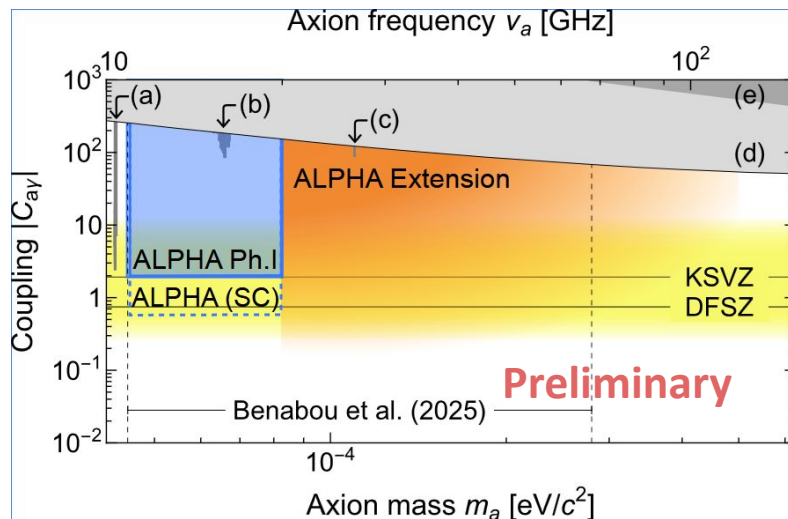
# Obligatory axion slide

Ciaran O'Hare limits

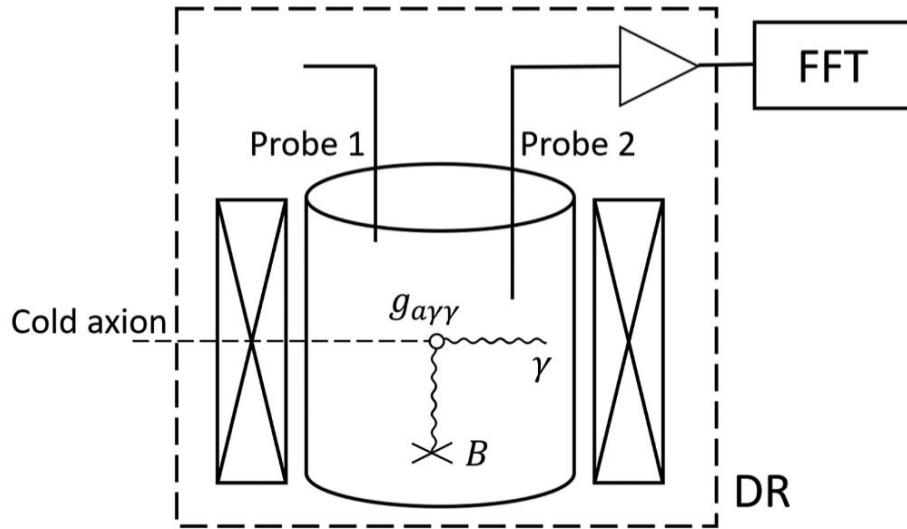


ALPHA target  
mass range

- ❖ Simulations place (post-inflationary) axion mass of several 10s of  $\mu\text{eV}$
- ❖ Few  $\mu\text{eV}$  region has good coverage, what about higher masses?



# The problem with higher masses



- ❖ Axion interacts with a magnetic field to convert to real photons
- ❖ Real photons excite a resonance inside of a cavity
- ❖ Mode of interest:  $TM_{010}$

*H. Chang. Rev. Sci. Instrum. 93, 084501 (2022)*

# The problem with higher masses

$$\frac{df}{dt} = g_{a\gamma\gamma}^4 \frac{\rho_a^2}{m_a^2} \frac{1}{\text{SNR}^2} \frac{B_0^4 \boxed{V^2 C^2}}{k_B^2 T_{\text{sys}}^2} \frac{\beta^2}{(1 + \beta)^2} Q_a \min(\boxed{Q_l}, Q_a)$$

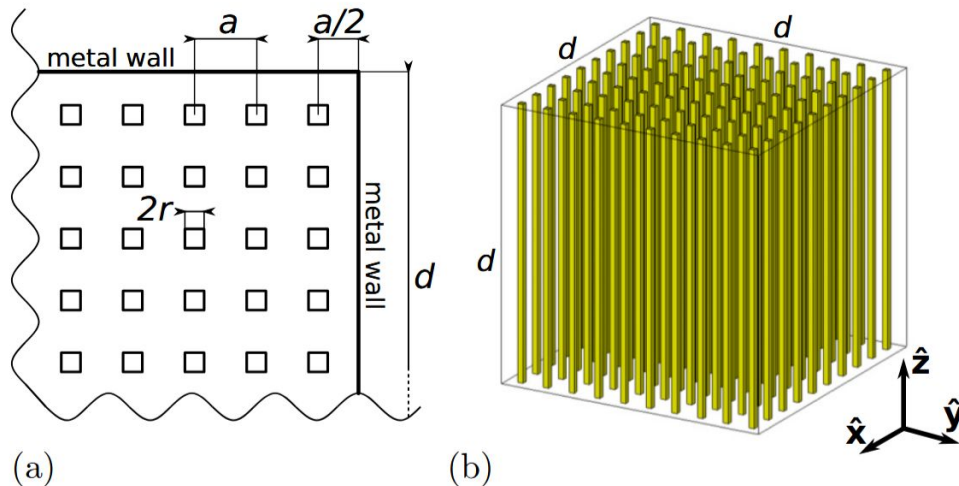
Controllable by cavity design

$\left\{ \begin{array}{l} V = \text{volume} \\ C = \text{form factor (coupling of E to external B)} \\ Q = \text{quality factor (proxy for losses)} \end{array} \right.$

$$\boxed{f_{010} = \frac{c}{2\pi} \frac{j_{01}}{R}} \quad (\text{for an empty cavity})$$

Higher mass  $\Rightarrow$  higher frequency  $\Rightarrow$  smaller  
volume  $\Rightarrow$  **slower scan speed**

# The ALPHA experiment and plasma haloscopes



- ❖ Wire grid decouples frequency from volume by using an array of wires
- ❖ Large volume, high frequency!
- ❖ Trade off: lower quality factor ( $Q$ )
  - Scan rate is only linear with  $Q$
  - Can be improved with superconducting wires
- ❖ ALPHA: **A**xion **L**ongitudinal **P**lasma **H**aloscope

*R. Balafendiev et al. Phys. Rev. B* **106**, 075106

$$k_p^2 = \frac{2\pi/a^2}{\ln \frac{a}{2\pi r} + F(1)}$$

(for the cavity shown above)

# ALPHA Collaboration



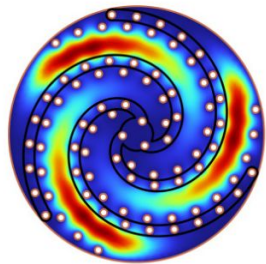
**2025 ALPHA Collaboration Meeting**



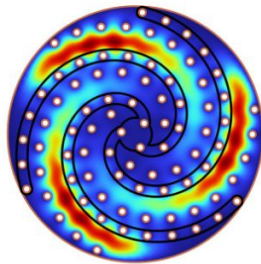
**Supported by:**



# An example of a plasma haloscope: The spiral cavity



(a)  $30^\circ$



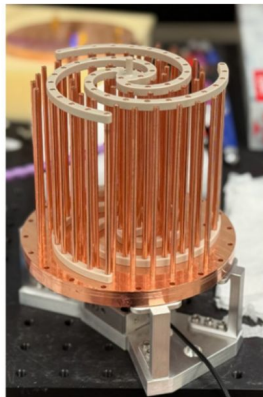
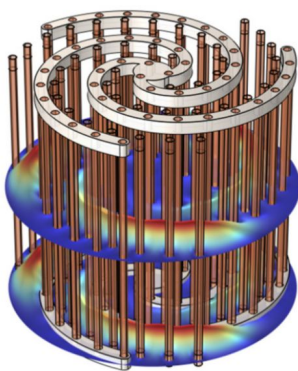
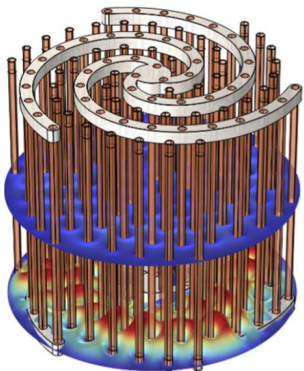
(b)  $15^\circ$



(c)  $0^\circ$

Low frequency

High frequency



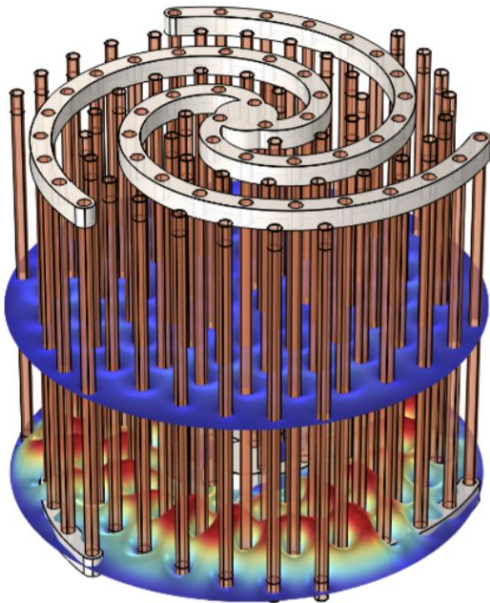
- ❖ Wires on spiral arms
- ❖ Turn one set of spirals relative to another
- ❖ Example shown: tunes from 9.7 to 12.7 GHz
- ❖ Currently being used for a dark photon run at Yale!
- ❖ Published in PR Applied: *J. Lindahl et. al. Phys. Rev. Applied* **25**, 024073



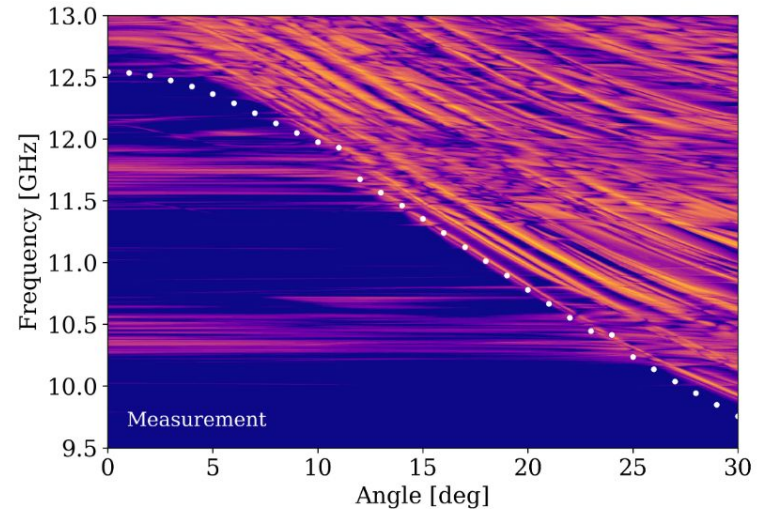
# An issue with the existing spiral cavity

Fixing structure:

- ❖ Complicated geometry
- ❖ Causes mode mixing

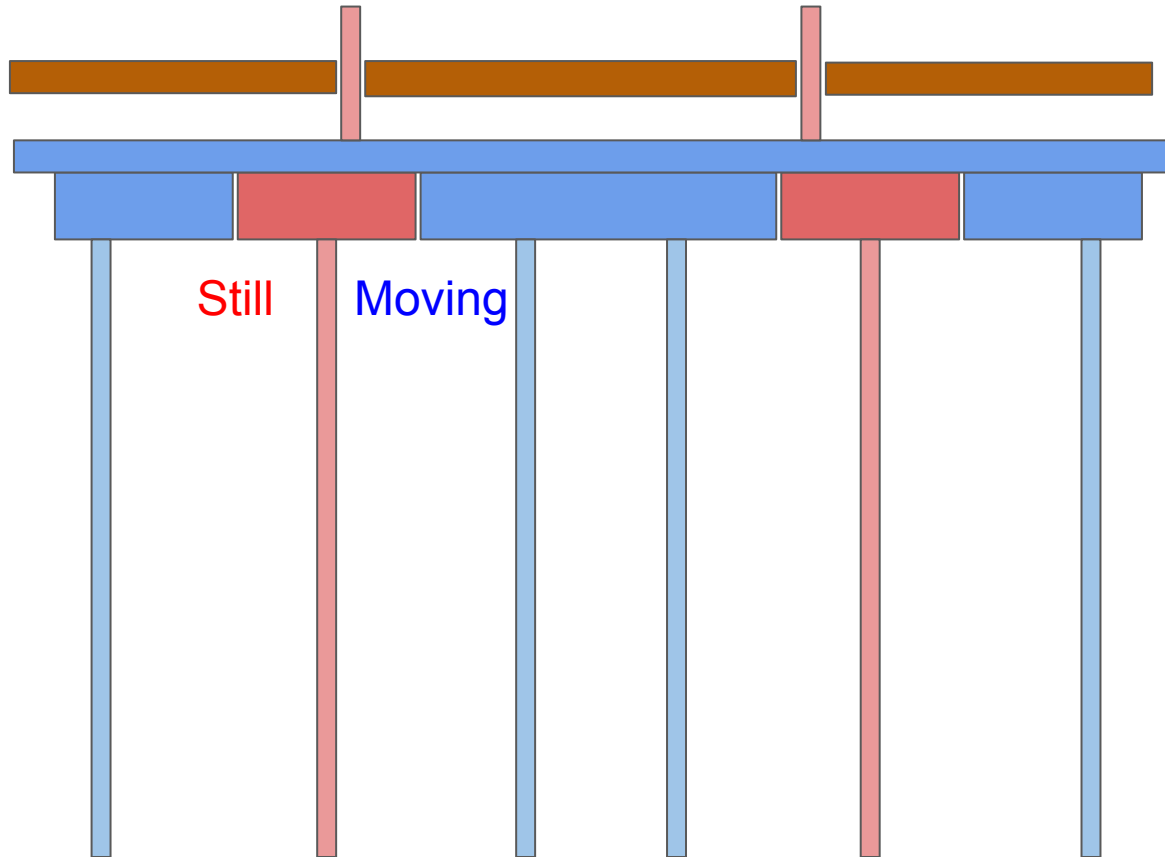


Longitudinal symmetry is broken with the spiral arm fixtures



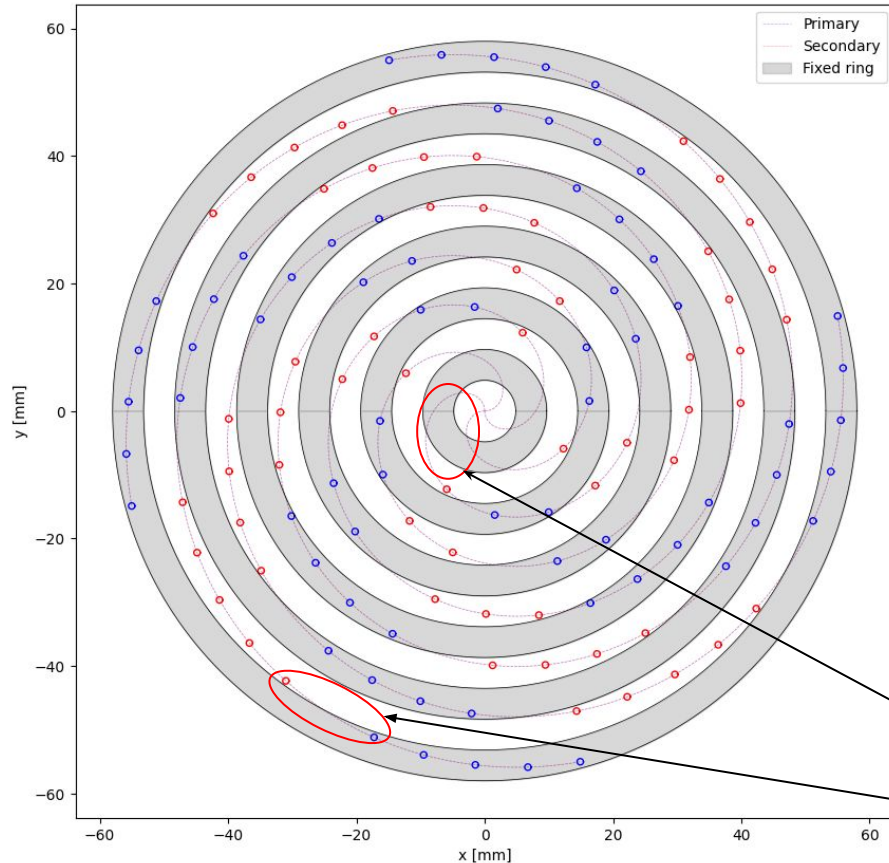


# A flush inner surface



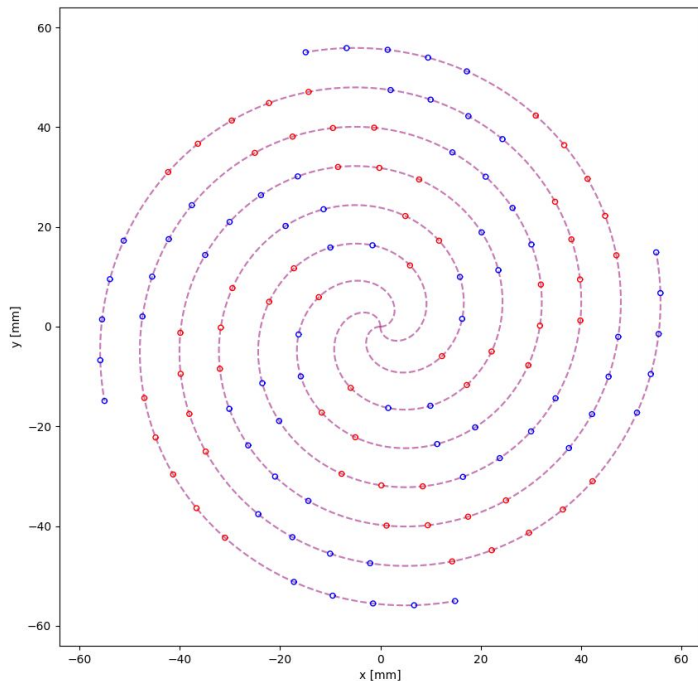
- ❖ Put still and moving rods on different concentric rings
- ❖ Mount all moving rods into one plate and move them together
- ❖ Inner surface is flush
- ❖ **One single rotation axis!**

# Rinnegan cavity: Ring Nested Group Rotation

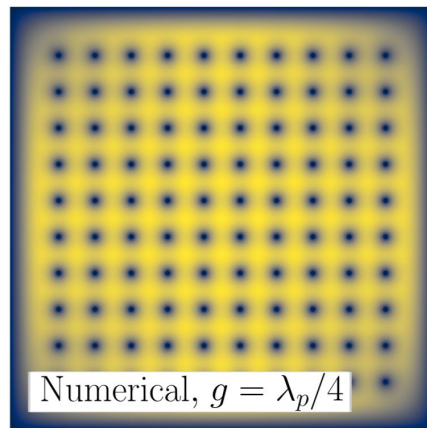


- ❖ Spirals are broken up along concentric rings
- ❖ **Sections** of a spiral move relative to each other
- Intersection points feature gaps (or off-spiral wires)

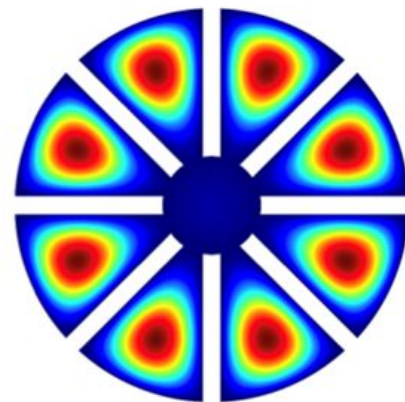
# A Broken Spiral on Concentric Rings



- ❖ Blue wires are on fixed rings
- ❖ Red wires are on moving rings
- ❖ **One single motor rotation!**
- ❖ Transitions from “uniform” to “pizza cavity”

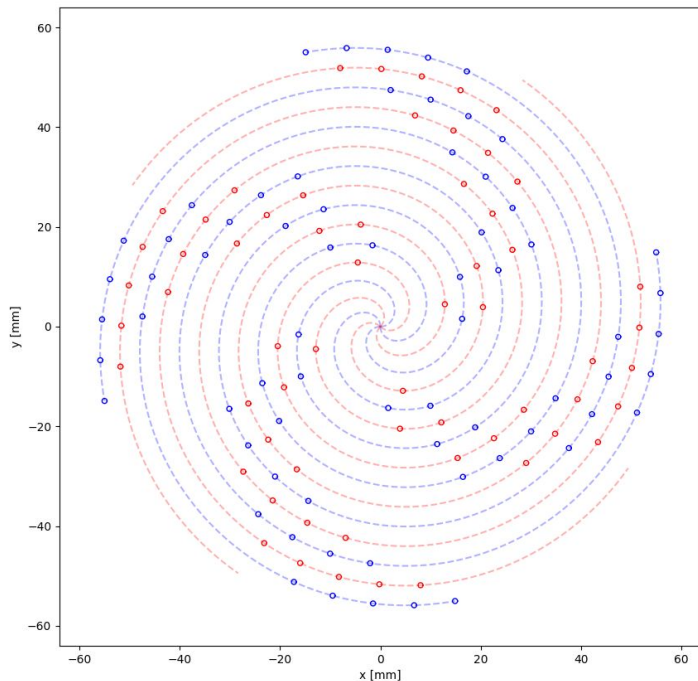


*Phys. Rev. Applied* **23**,  
054053 (2025)

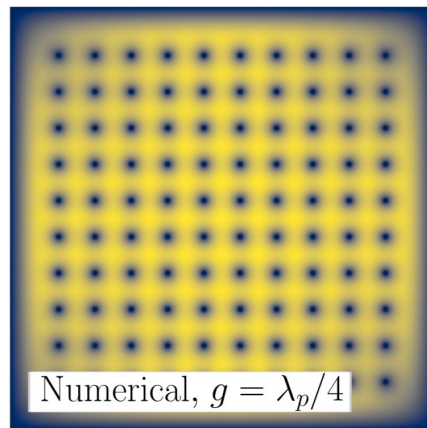


*Phys. Rev. Lett.* **125**,  
221302 (2020)

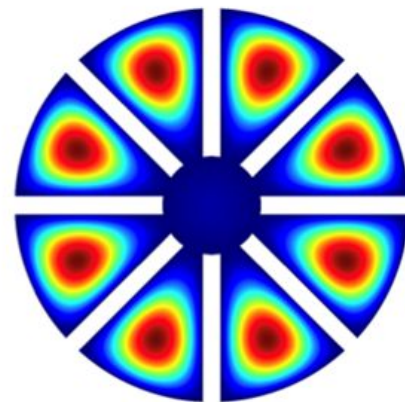
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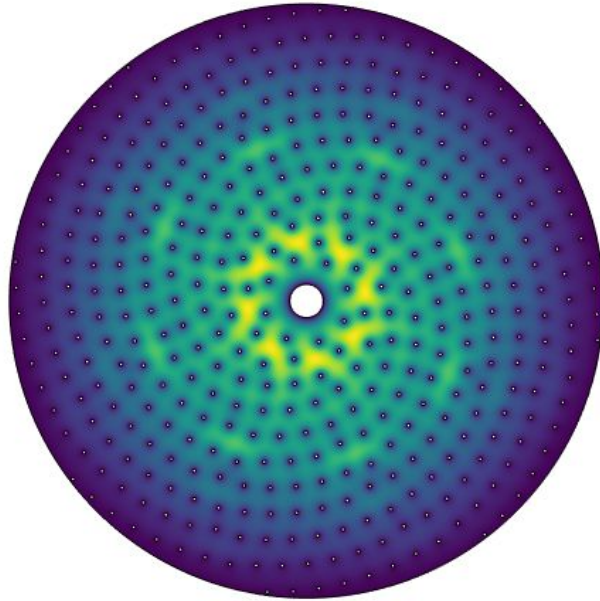
*Phys. Rev. Applied* **23**,  
054053 (2025)



*Phys. Rev. Lett.* **125**,  
221302 (2020)

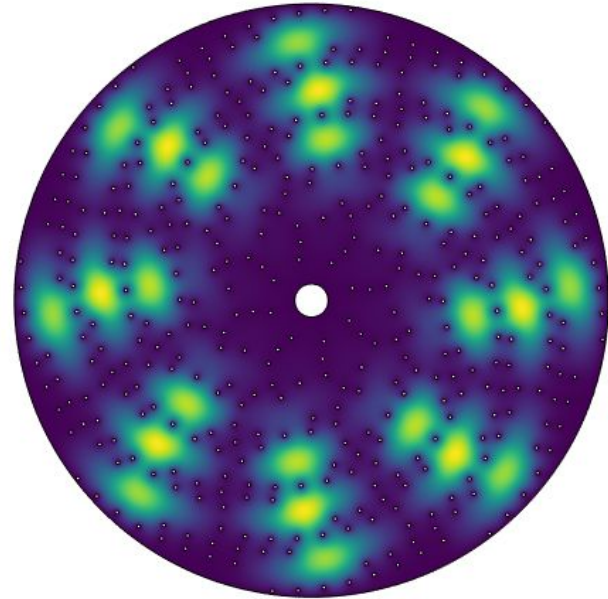
# Field simulations

- ❖ Uniform
- ❖ High  $f_0$



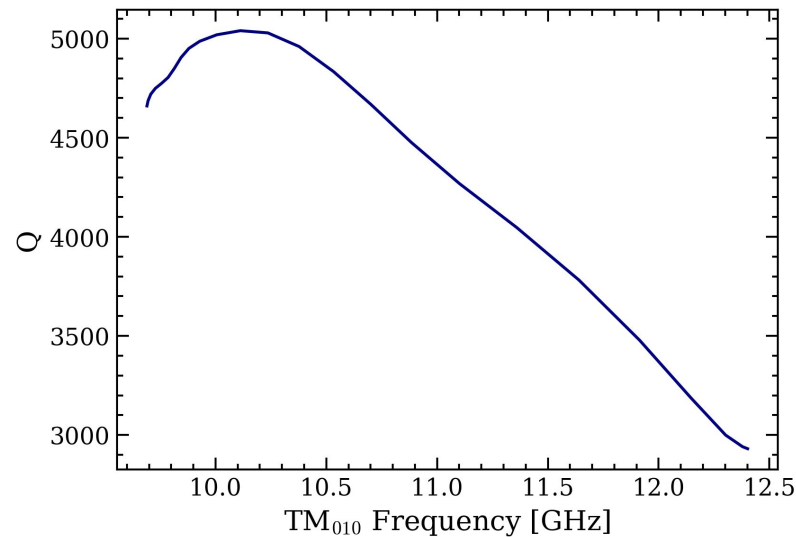
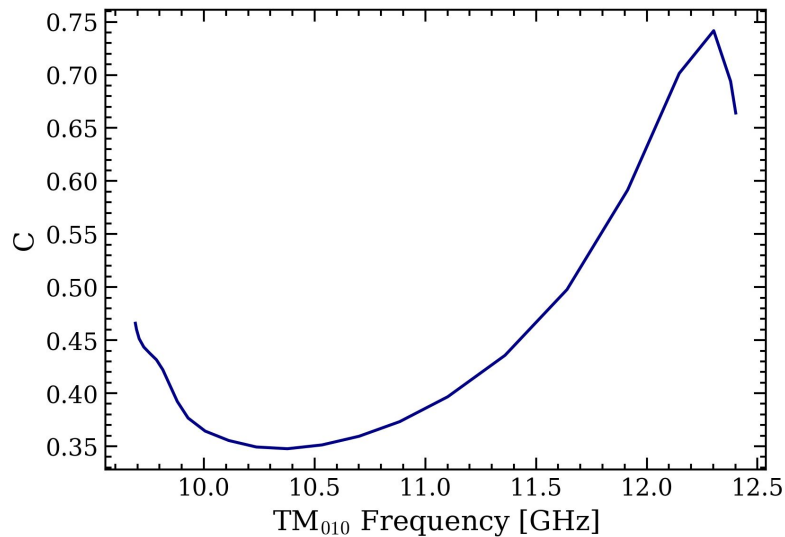
$\theta = 0$

- ❖ Pizza
- ❖ Low  $f_0$



$\theta = \pi/8$

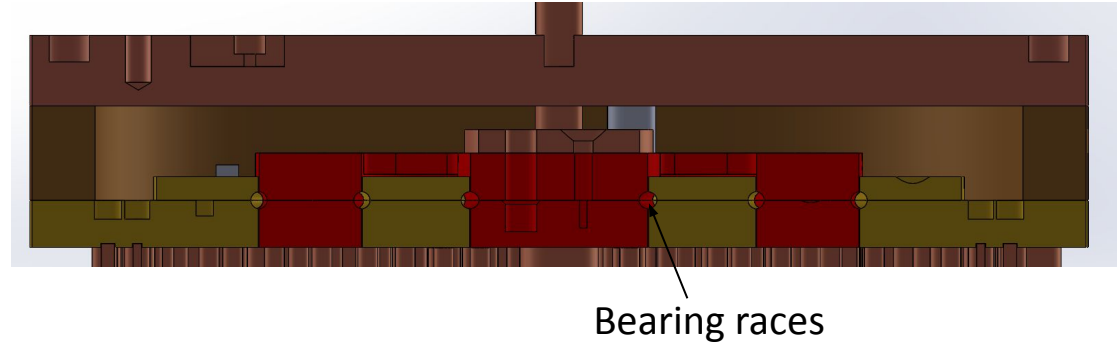
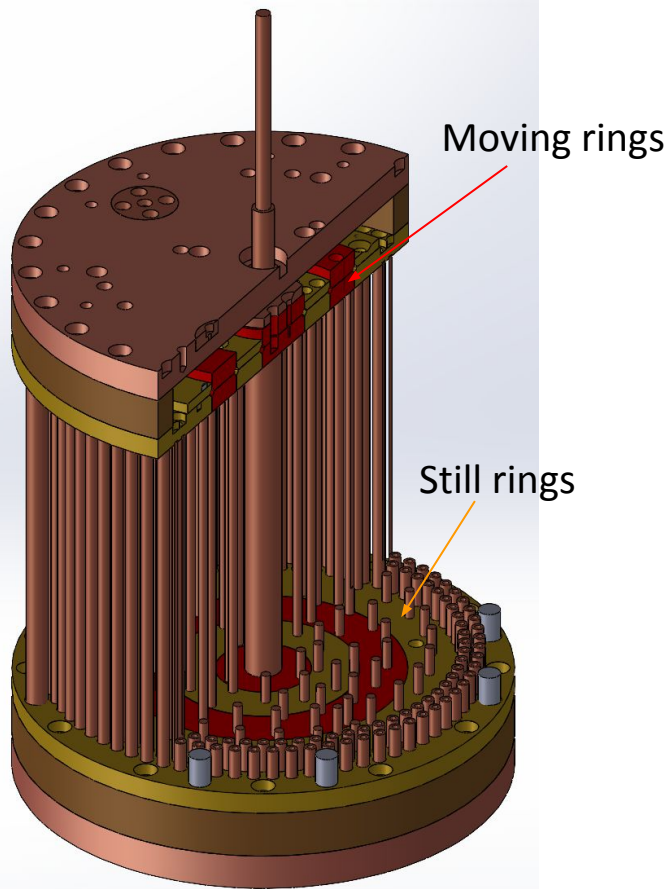
# Frequency Range and Form Factor



- ❖ Tuning range from 9.7 GHz to 12.4 GHz
- ❖ Maintains a reasonable form factor for 1mm wires
- ❖ Quality factor decreases for thin wires

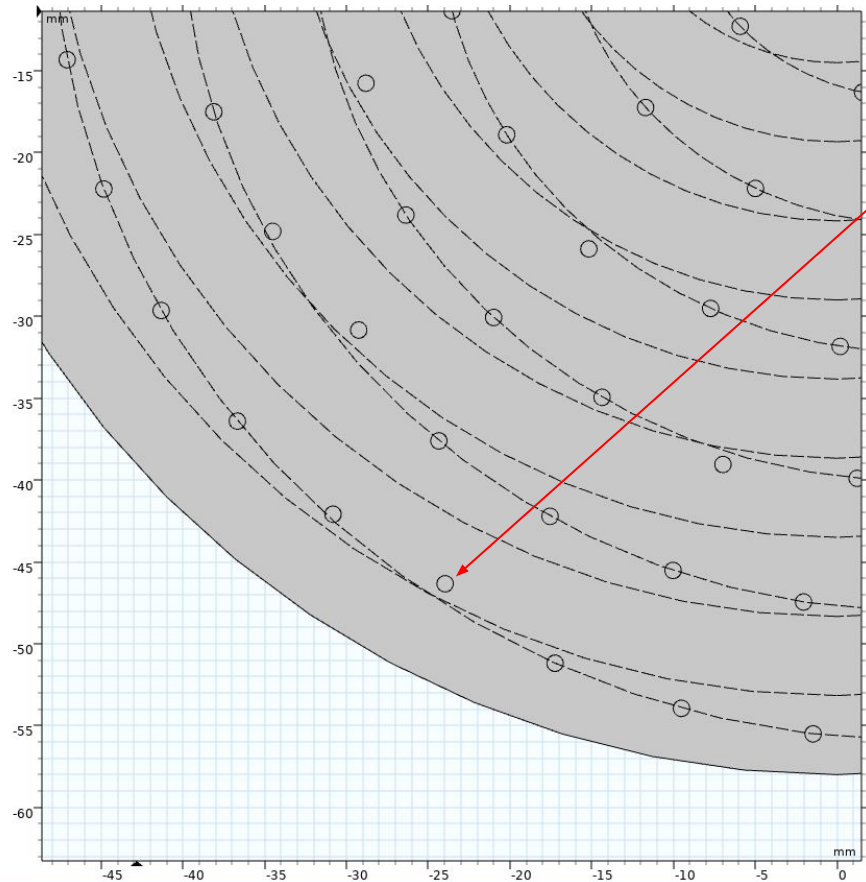


# Ongoing prototyping for the Rinnegan cavity



- ❖ 4 rings
- ❖ 3 spirals
- ❖ 25 wires per ring
- ❖ 2.5 mm diameter wires
- ❖ 100  $\mu\text{m}$  gaps between rings

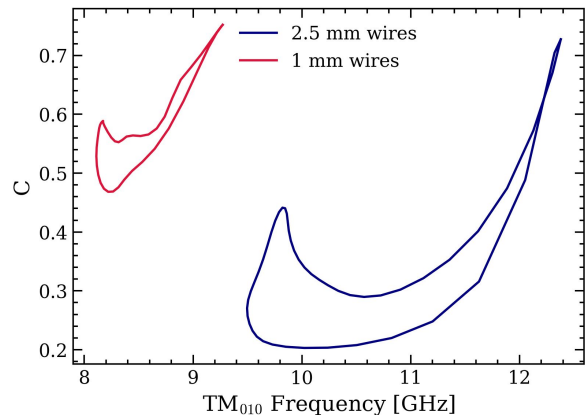
# Potential drawback #1: Wires that intersect the rings



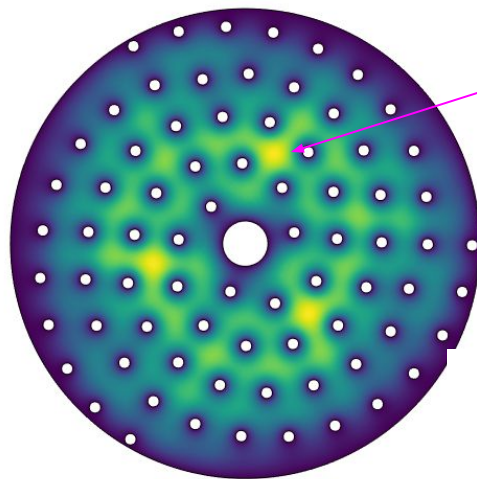
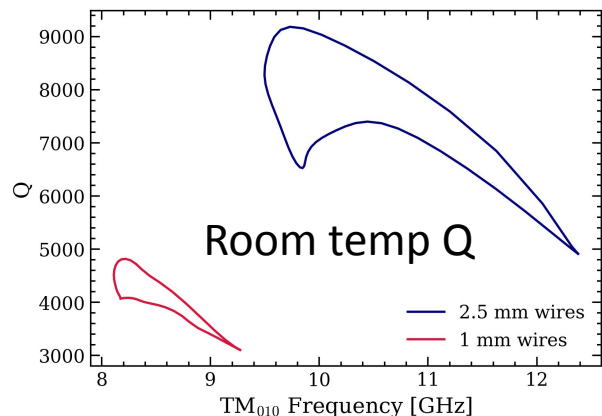
Removed wires create a gap for the field to localize

- ◆ Degrades form factor
- ◆ Must be reintroduced with a slight offset

# Potential drawback #1: Wires that intersect the rings

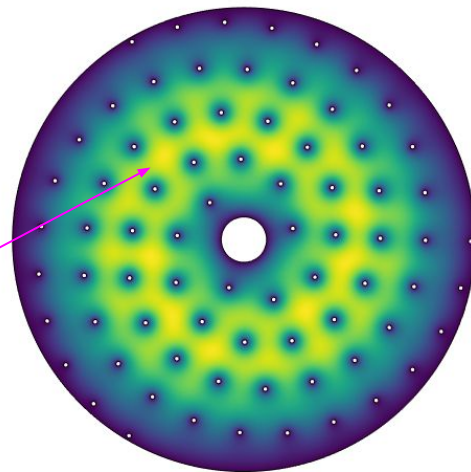


**Not an apples to apples comparison!**



Thicker wires has more localization

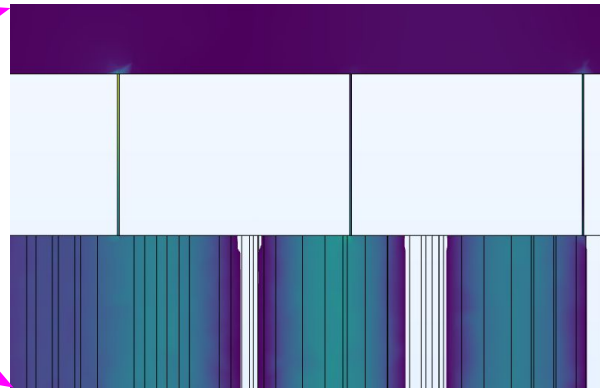
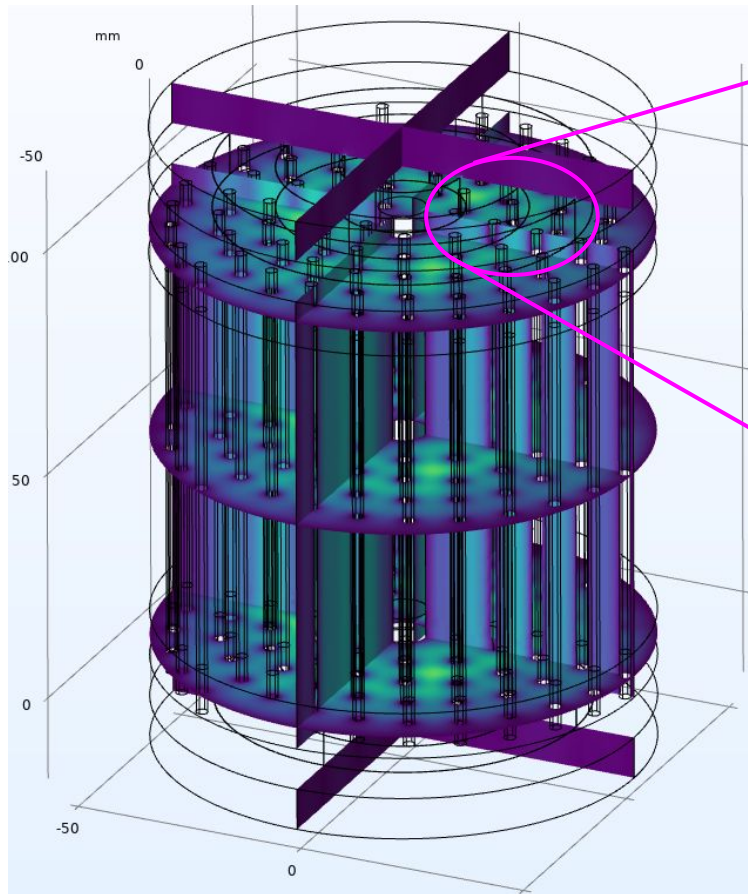
Offset can be as big as the wire radius



Thinner wires have less localization

But lower Q

## Potential drawback #2: Losses due to seams



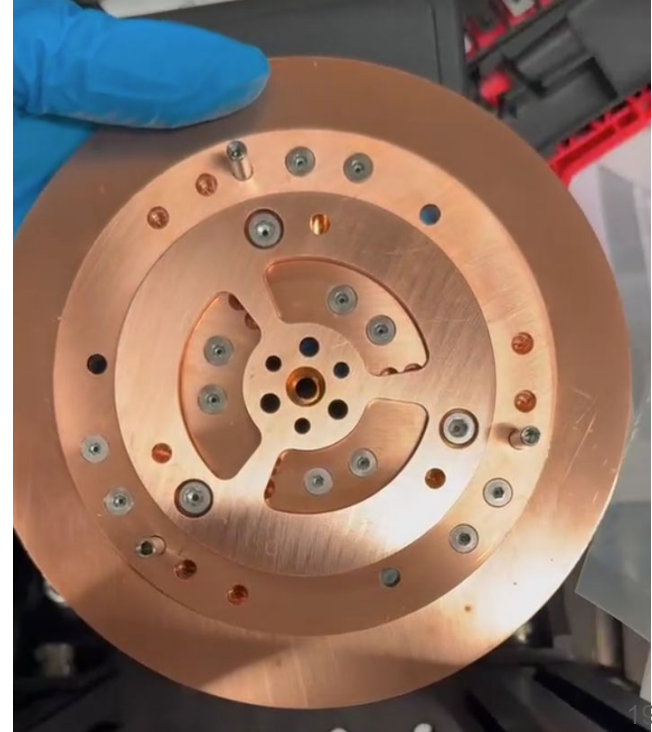
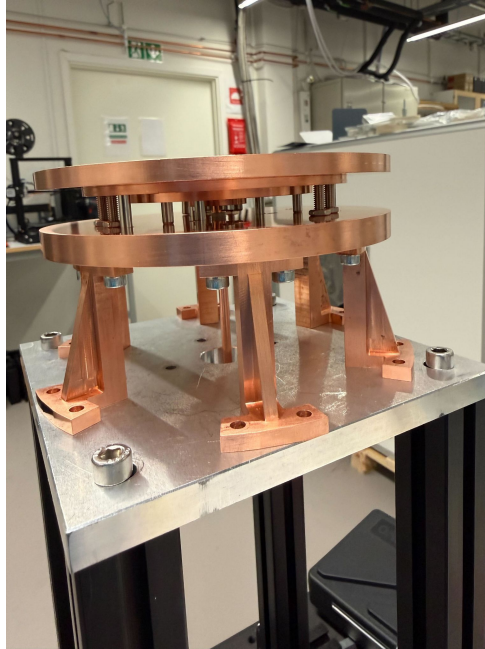
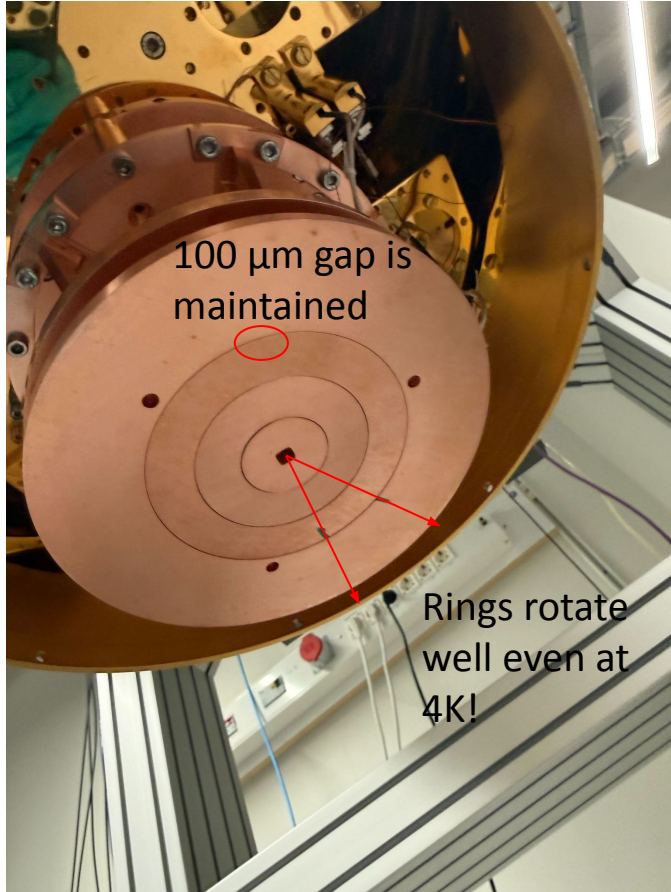
Some field is localized in the seams

**Q and C drop, but slightly at 100  $\mu\text{m}$  gap**

- ❖ Without seams (highest frequency  $\text{TM}_{010}$ ):
  - $Q = 4780, C = 0.73$
- ❖ With seams:
  - $Q = 4577, C = 0.61$
- ❖ **Effect of ball bearings to be seen**



# Cryogenic rotation tests



# Thanks for listening!

- ❖ Several GHz tunability to look for high mass axions
- ❖ Single rotating axis
- ❖ To be seen:
  - The effect of ball bearings and seams on  $Q$
  - The optimal way of packing wires

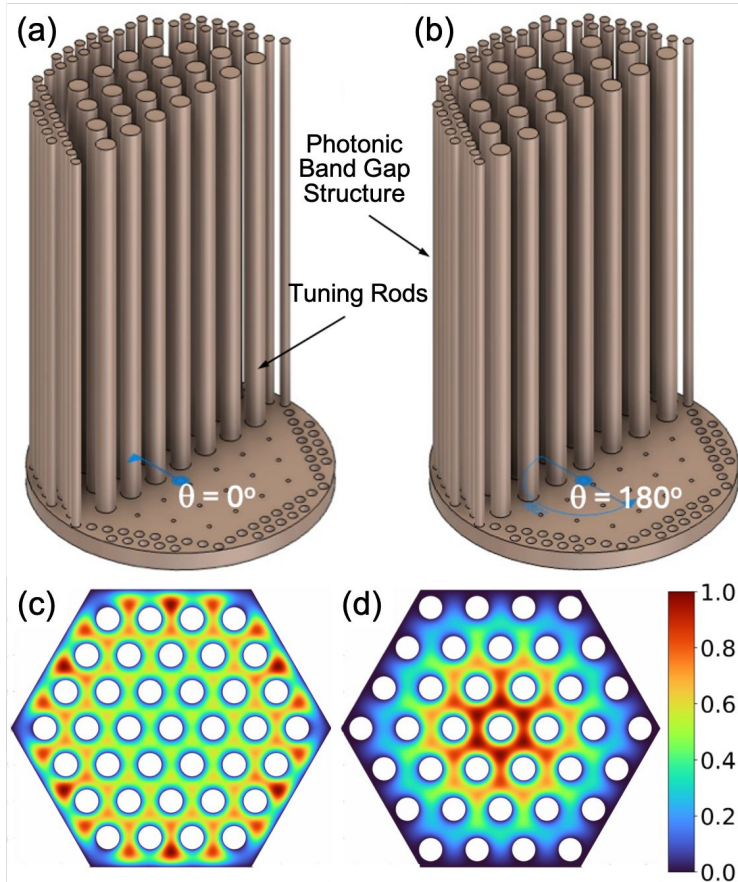


**Backup**

# ALPHA Timeline

- ❖ Currently commissioning the dark photon run with a prototype spiral cavity
- ❖ Resonator design is being finalized and production will begin soon
- ❖ Magnet delivery and full commissioning in 2027
- ❖ Physics data taking to begin in 2028

# Berkeley cavity



- ❖ Preserves symmetry
- ❖ Thicker rods for high Q
- ❖ Off axis axle drives several rod axes